



# Novel miniaturized arrow linked figure of eight square loop FSS-based multiband linear-circular and linear-cross reflective microwave polarizer

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## Research Paper

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### Abstract

This work demonstrates a novel, reflective-type, dual-polarized multiband polarization transformer for various C to Ku band applications. The design constitutes a top FSS printed on a thin FR-4 grounded dielectric substrate. A  $y/x$ -polarized incident wave experiences linear-circular conversion from 4.61 to 4.71 GHz (right-hand circular polarized (RHCP)/left-hand circular polarized (LHCP)), 5.44 to 7.68 GHz (LHCP/RHCP), 10.13 to 14.98 GHz (RHCP/LHCP), and 16.68 to 17.19 GHz (LHCP/RHCP) with 3.41, 34.15, 38.63, and 3.01% FBW, respectively. In addition, it performs linear-cross conversions with a minimum 90% polarization conversion ratio (PCR) from 4.92 to 5.18 GHz, 8.32 to 9.41 GHz, and 15.56 to 16.25 GHz with 5.15, 12.30, and 4.34% FBW. Multiple surface plasmonic resonances are responsible for these polarization conversions in different bands. The device performance is angular stable for both transverse electric (TE) and transverse magnetic (TM) incidences. The design novelty lies in compacted periodicity  $0.112\lambda_L$ , low-profile thickness  $0.037\lambda_L$  with a multi-conversion (switching the orthogonal circular rotation in consecutive bands), and stable angular response, which can potentially be used in real-time applications.

## Introduction

Polarization of an electromagnetic (EM) wave gives information about the electric field oscillating direction in the plane perpendicular to the direction of propagation [1]. Manipulating wave polarization has a significant role in applications like RCS reduction, sensing, and imaging from microwave [2] to THz regime [3]. Generally, the Birefringence effect [4], circular dichroism [5], and molecular level chirality [6] occur in naturally occurring materials and are responsible for polarization modulation. But these crystals suffer from huge power losses and greater thickness, limiting their usage in real-time sub-miniaturized applications [7].

Frequency selective surfaces (FSSs) are the 2D-spatial filters that possess advantages like planar nature, low profile, and fabrication ease [8]. The advancement of FSS technology has drawn researchers' attention in modulating the EM wave features like amplitude, phase, and polarization states using FSS [9]. Controlling the propagation states of EM waves has been reported in reflection and transmission modes over the EM spectrum [10]. A compact, high-efficiency microwave linear-cross polarizer based on square split rings is demonstrated [9], with unity cross-pol reflectance from 8.2 to 23 GHz. Huang *et al.* proposed an ultrathin dual-band reflective polarizer that performs linear-cross conversion in frequency ranges 4.40–5.30 GHz and 9.45–13.60 GHz [11]. A 45° oriented H-shaped linear-cross reflective polarizer is demonstrated for K- and Ka-bands (17.97–40.23 GHz) with minimal bandwidth reduction at higher oblique angles [12]. This design can also maintain handedness to the circularly polarized wave. Asymmetric meanderline structures are also utilized for broadband linear-cross conversion [13].

For circular polarization conversion, Jerusalem cross-based reflective linear-circular polarizer with an axial ratio (AR)  $\leq 1$  dB bandwidth of 15% is proposed in X-band [14]. Multiband reflective type multi-polarization conversion is demonstrated by a double L-type resonator of different dimensions [15]. From 11.3 to 20.2 GHz, a flexible broadband linear-cross reflective polarizer with an average polarization conversion ratio (PCR) of 85% is designed [16]. Active devices like p-i-n, varactor diodes, and their integration with metasurfaces have also been illustrated for multifunctional characteristics [17]. The linear-circular invertibility is also achieved using an ultrathin broadband single-layer polarization converter in a near-infrared regime [18]. However, few studies have focused on converting a linear to a cross and circular polarization simultaneously across wideband.

This work proposes a novel single-layered dual-polarized multiband reflective type metasurface enabled linear-circular and linear-cross polarizer. The design of the unit cell and its simulation results are detailed in section "Design of polarizer's unit cell". Surface current

distributions in section “Analysis and discussion” elaborate the analysis for the multi-conversion. The measured results obtained from the free-space measurement technique are compared with the simulated one in section “Fabrication and measurement setup”. Good matching is observed between those two, and finally, the conclusion has been drawn in “Conclusion” section.

### Design of polarizer’s unit cell

The proposed reflective polarizer consists of an eight-lobbed square loop as top FSS and orthogonally linked arrow, as shown in Fig. 1(a). The FSS is printed on the top of a commercially available FR-4 substrate with relative permittivity  $\epsilon_r = 4.4$ , and loss-tangent  $\tan \delta = 0.02$ , which is metal-backed on the other side. Copper film with conductivity  $\sigma = 5.8 \times 10^7$  S/m is used as metallic layer. The unit cell dimensions are optimized in a finite element method (FEM) based EM solver ANSYS HFSS with appropriate periodic boundary conditions. The final parameters are:  $p = 7.3$  mm,  $l_0 = 3.52$  mm,  $l_1 = 1.584$  mm,  $l_2 = 3.96$  mm,  $w = 0.176$  mm,  $g_0 = 0.176$  mm, and  $g_1 = 0.13$  mm,  $h = 2.4$  mm, and  $t = 0.035$  mm, as depicted in Fig. 1.

For a  $y$ -polarized wave incidence on the reflective polarizer, the co-pol ( $r_{yy} = |E_{yr}|/|E_{yi}|$ ), and cross-pol ( $r_{xy} = |E_{xr}|/|E_{yi}|$ ) reflectances are defined due to the asymmetry that exists. Here,  $E_x$  and  $E_y$  represent electric field oscillations along the  $x$ - and  $y$ -axes, respectively, with indices  $i$  and  $r$  denoting incident and reflected components. The simulated magnitudes and relative phase of the reflectances are shown in Fig. 2, and it seems clear that  $r_{yy}$  approaches zero at three frequencies, 5.05, 8.85, and 15.95 GHz, while  $r_{xy}$  reaches near-unity (see Fig. 2(a)). The term PCR determines the efficiency of any reflective linear-cross polarizer, expressed as

$$PCR = r_{xy}^2 / (r_{xy}^2 + r_{yy}^2) \quad (1)$$

The PCR is above 90% from 4.92 to 5.18 GHz (C-band), 8.32 to 9.41 GHz (X-band), and 15.56 to 16.25 GHz (Ku-band), as shown in Fig. 3(a). In addition, the amount of energy getting dissipated in the substrate during the multiple reflections between the FSS and ground is estimated by energy conversion ratio

(ECR),  $ECR = r_{xy}^2 + r_{yy}^2$ . The remaining energy is wasted via substrate losses, with almost 90% of incident energy reflected (see Fig. 3(a)). Besides linear-cross conversion, it also demonstrates linear-circular conversion. The circular polarization efficiency is calculated from the axial ratio (AR), calculated as in [18],

$$R = \frac{r_{xy}}{r_{yy}} \quad (2)$$

$$Schi, \chi = 0.5 \sin^{-1} \left( \frac{2R \sin(\Delta\phi)}{1 + R^2} \right) \quad (3)$$

$$\text{Axial ratio, } AR = |1/\tan(\chi)| \quad (4)$$

Here, the term *schi* ( $\chi$ ) represents the shape of the polarization ellipse. The simulated AR is  $\leq 3$  dB over the ranges 4.61–4.71 GHz (C-band), 5.44–7.68 GHz (C-band), 10.13–14.98 GHz (X, Ku-band), and 16.68–17.19 GHz (Ku-band), as shown in Fig. 3(b).

Since polarizer designs work based on the multiple plasmonic resonances, and three plasmonic resonances are dominant for the proposed design. The first resonance depends on the outer periphery of the figure of eight loops in the  $v$ -axis direction. The second resonance mainly depends on the substrate parameters like relative permittivity and thickness concerning ground, and the third resonance depends on the dimensions of the figure-of-eight loop along the  $u$ -axis direction (see Fig. 6). The PCR, ECR, and AR (dB) performance with and without arrow dipole is given in Fig. 3. For the design optimization for better performance, it is necessary to shift the higher resonance to the lower side by coupling along the  $u$ -axis direction to achieve a good AR bandwidth  $\leq 3$  dB. The arrow-shaped dipole is included in the proposed polarizer with square loops. The term ellipticity ( $e$ ) can be used to describe the switching of circular rotation in successive bands, expressed as

$$\text{Ellipticity, } e = \frac{2|r_{xy}||r_{yy}| \sin(\Delta\phi)}{|r_{xy}|^2 + |r_{yy}|^2} \quad (5)$$

The ellipticity range is varied from  $-1$  (right-hand circular polarized, RHCP) to  $+1$  (left-hand circular polarized, LHCP), as

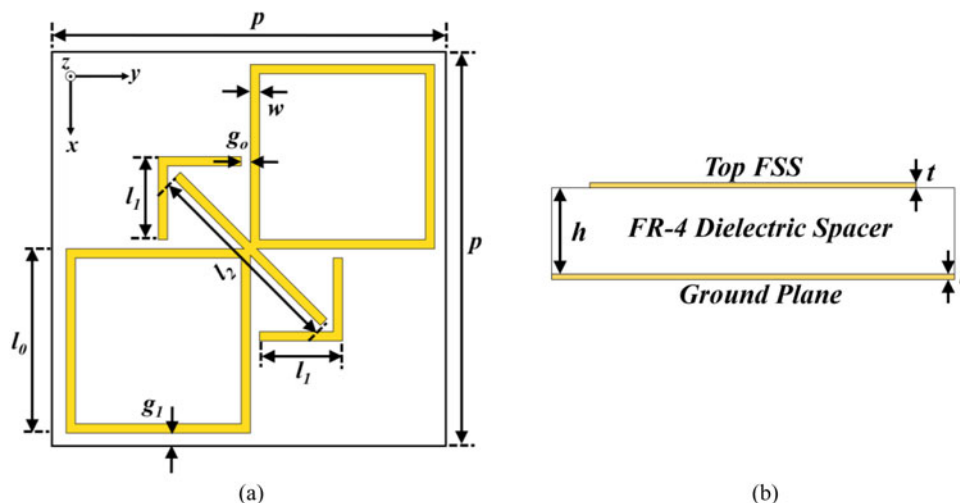


Fig. 1. (a) Top view, and (b) side view of the proposed reflective polarizer.

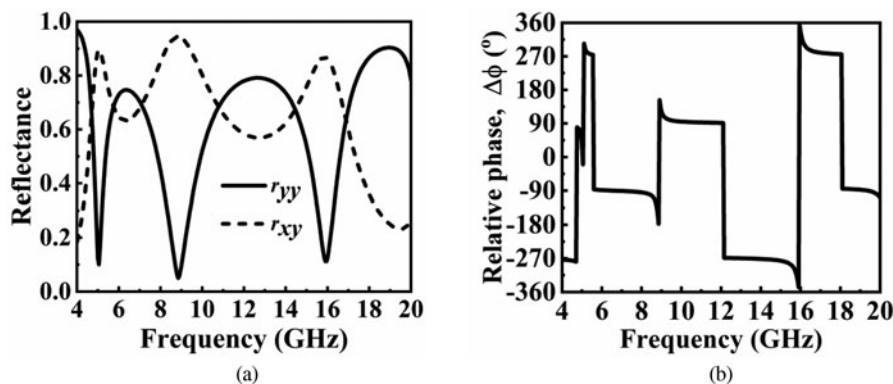


Fig. 2. Simulated (a) magnitudes and (b) phases of the co-pol and cross-pol reflectances of the proposed polarizer.

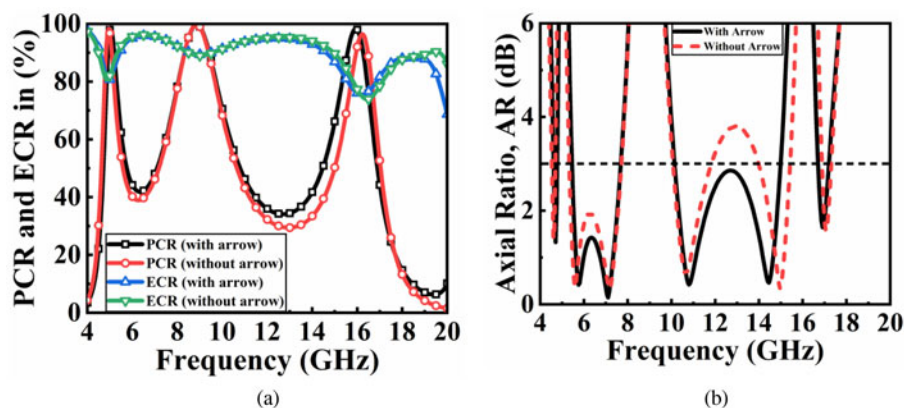


Fig. 3. Simulated (a) PCR, ECR, and (b) AR (dB) of the proposed polarizer with and without arrow dipole.

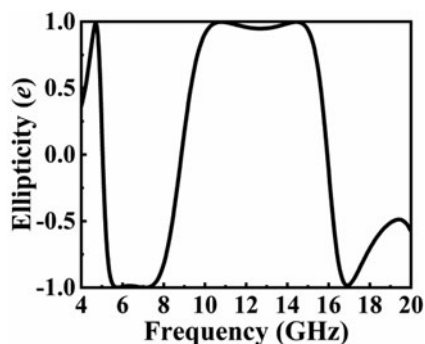


Fig. 4. Simulated ellipticity ( $e$ ) of the proposed reflective polarizer.

shown in Fig. 4. The handedness of circular rotation is sensitive to the relative phase ( $\Delta\phi$ ) between co- and cross-pol reflectances, as in Fig. 2(b). This design is dual-polarized; the circular polarization’s handedness will change if the  $y$ -polarized incidence is replaced with the  $x$ -polarized.

Since the incoming wave on the metasurface is not always normal in practical scenarios, it is necessary to investigate the performance variation with the angle of incidence for both the transverse electric (TE) and transverse magnetic (TM) modes. The PCR and AR variations for different angles ( $0^\circ$  to  $60^\circ$  in steps of  $15^\circ$ ) are depicted in Fig. 5. If the incident angle varies, it is trivial that the input surface impedance seen by the EM wave is different for different incident angles and modes. For the proposed design, it is observed that the response is almost stable up to  $45^\circ$  (see Fig. 5), and

performance degrades due to grating lobes at higher oblique angles and weak field couplings. Although PCR solely considers reflectance magnitudes, the behavior of PCR at various oblique angles is well explained by reflectance phases. PCR bandwidth is reduced as the angle of incidence is moved away from the normal. The cross-pol component ( $r_{xy}$ ) has nearly unchanged amplitude and phase, whereas the co-pol reflectance has a significant angle of incidence dependence. The additional path between the metasurface and the ground influences the phase of co-pol reflectance significantly, creating a phase difference compared to normal incidence calculated by Bragg’s diffraction [12].

$$\Delta\phi = \phi_{oblique} - \phi_{normal} = 2\sqrt{\epsilon_r}kd \left( \frac{1}{\sqrt{1 - \frac{\sin^2\theta}{\epsilon_r}}} - 1 \right) \quad (6)$$

Here,  $\theta$  is angle made by the incident ray with reference to polarizer normal,  $k$  is the wavevector in the dielectric medium. The reason behind the multiband multi-polarization conversion is analyzed in later sections.

### Analysis and discussion

#### Analysis using transfer matrix method (TMM)

To acquire a true understanding of the physical mechanism involved in the polarization conversion, two orthogonal axes,  $u$ -

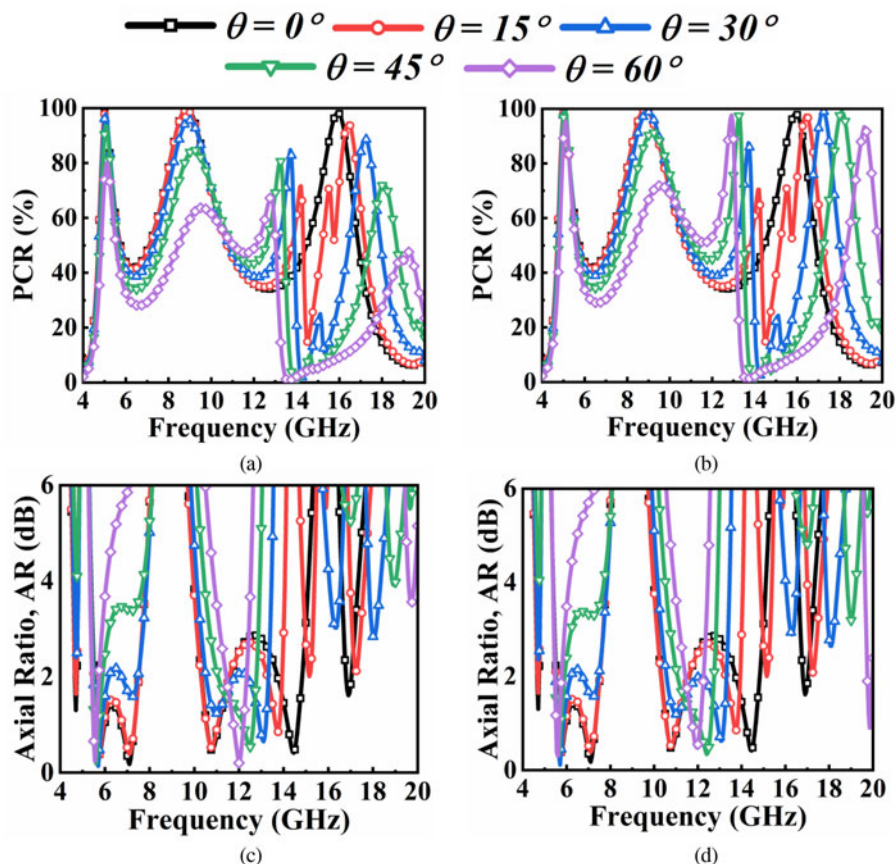


Fig. 5. Simulated PCR for the (a) TE, (b) TM, and AR (dB) for the (c) TE and (d) TM modes for different elevation angles ( $\theta$ ).

and  $v$ -, are specified at  $45^\circ$  counterclockwise to the  $x$ - $y$  axes (see Fig. 6). The incident electric field for a wave traveling along the negative  $z$ -axis is expressed as  $\vec{E}_i = (E_i^u \hat{u} + E_i^v \hat{v})e^{jkz}$ , where  $E_i^u$  and  $E_i^v$  are the components along  $u$ - and  $v$ -axes, respectively.

From the TMM technique [12], the reflected field components can be written as

$$\begin{bmatrix} E_r^u \\ E_r^v \end{bmatrix} = \begin{bmatrix} r_{uu} & r_{uv} \\ r_{vu} & r_{vv} \end{bmatrix} \begin{bmatrix} E_i^u \\ E_i^v \end{bmatrix} \tag{7}$$

The influence of cross-pol reflectance is very low and ignored since the unit cell design is symmetric along  $uv$ -axes. Then the reflected field component can be expressed as

$$\vec{E}_r = (r_{uu}E_i^u \hat{u} + r_{vv}E_i^v \hat{v})e^{-jkz} \tag{8}$$

The simulated magnitudes and phases are depicted in Figs 7(a) and 7(b), respectively. The magnitudes of  $r_{uu}$  and  $r_{vv}$  are almost unity over the band, with a tiny dip at frequencies 4.95, 9.05, and 16.15 GHz (see Fig. 7(a)). The eigenmodes at 4.95 and 16.15 GHz are stronger (have a higher Q-factor) than the mode at 9.05 GHz. If the substrate loss is ignored,  $r_{uu} = e^{j\varphi_u}$ ,  $r_{vv} = e^{j\varphi_v}$ , and  $\Delta\varphi = \varphi_u - \varphi_v$  are obtained [19]. Different combinations of magnitudes and phases with the conditions for possible polarization conversions are summarized in Table 1.

Surface current distribution analysis

Surface current patterns of the polarizer’s metallic portion are investigated to better understand the physical phenomena behind polarization conversion. At resonant frequencies, 5.05, 8.85, and 15.95 GHz, the current profiles for the top FSS and ground are depicted in Fig. 8. The symmetric and asymmetric linkage of

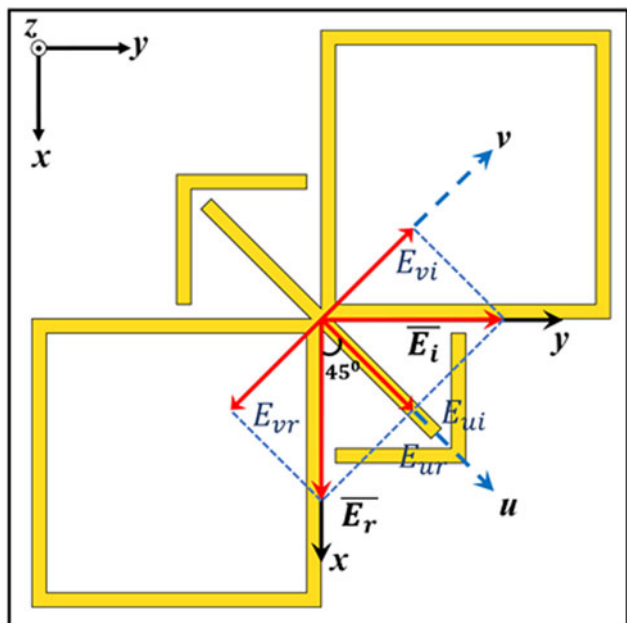


Fig. 6. Schematic view of the proposed reflective polarizer unit cell for  $uv$ -analysis.

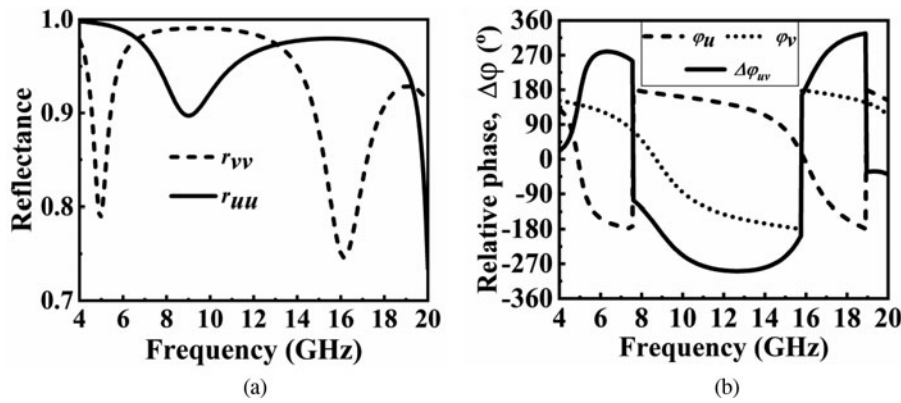


Fig. 7. Simulated (a) magnitudes of  $r_{uu}$ ,  $r_{vv}$ , and (b) relative phase ( $\Delta\phi_{uv}$ ).

Table 1. Summary of polarization conversion for different relative phases ( $\Delta\phi$ ) and magnitudes ( $r_{uu}$  and  $r_{vv}$ )

Sl no	$\phi_u$	$\phi_v$	$\Delta\phi = \phi_u - \phi_v$	$r_{uu}$	$r_{vv}$	Reflected field component ( $\vec{E}_r$ )	Polarization conversion
1	$\pi$	0	$+\pi$	-1	+1	$\vec{E}_r = (-E_x^i \hat{u} + E_y^i \hat{v})e^{-jkz}$	Linear-cross (negative x-direction)
2	0	$\pi$	$-\pi$	+1	-1	$\vec{E}_r = (+E_x^i \hat{u} - E_y^i \hat{v})e^{-jkz}$	Linear-cross (positive x-direction)
3	$-\pi/2$	$+\pi$	$+\pi/2$	-j	-1	$\vec{E}_r = (-jE_x^i \hat{u} - E_y^i \hat{v})e^{-jkz}$	Linear-circular (LHCP)
4	$+\pi/2$	$-\pi$	$-\pi/2$	+j	-1	$\vec{E}_r = (+jE_x^i \hat{u} - E_y^i \hat{v})e^{-jkz}$	Linear-circular (RHCP)

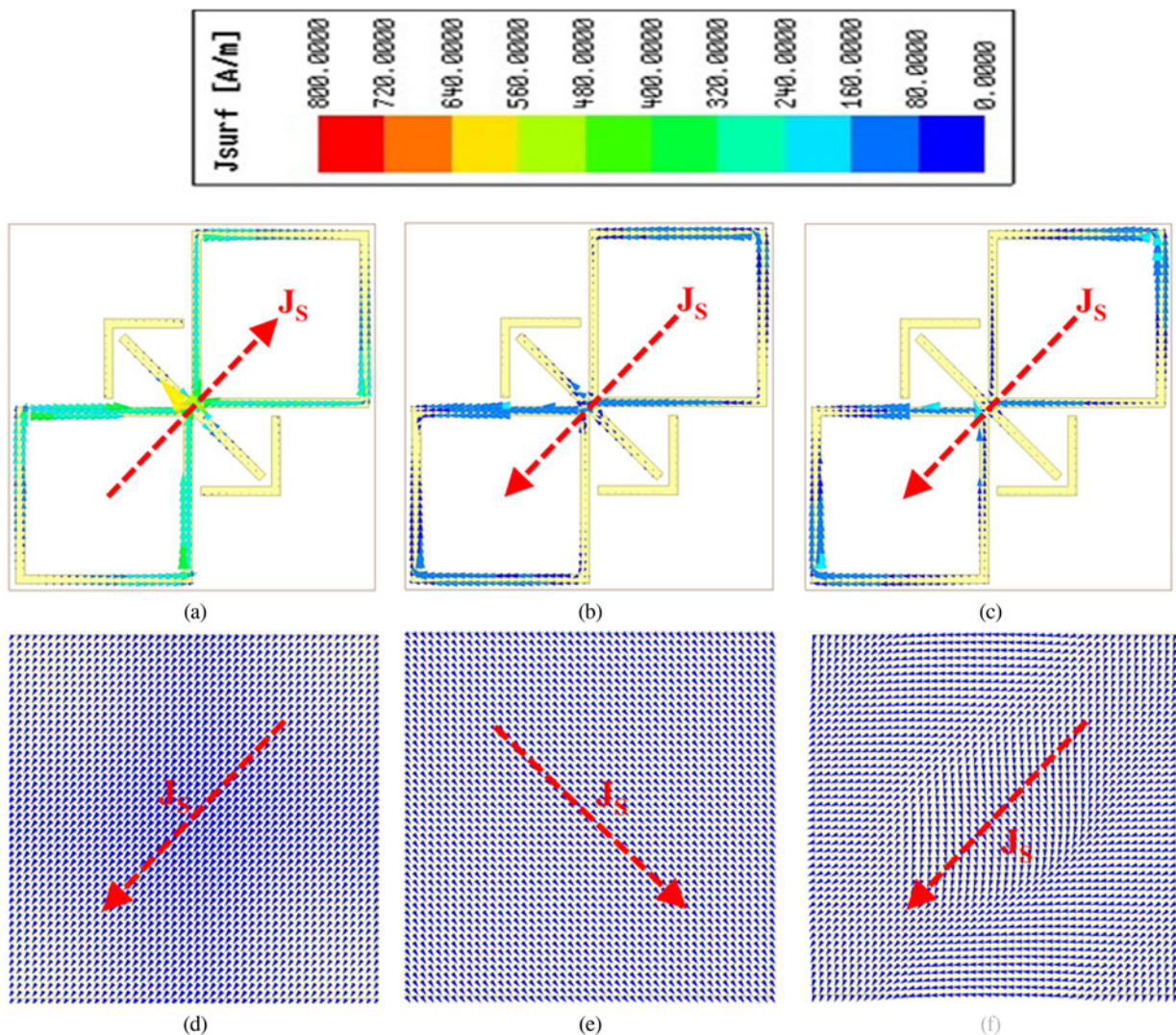


Fig. 8. Surface current distributions of top FSS at (a) 5.05 GHz, (b) 8.85 GHz, (c) 15.95 GHz, and for the ground at (d) 5.05 GHz, (e) 8.85 GHz, (f) 15.95 GHz.

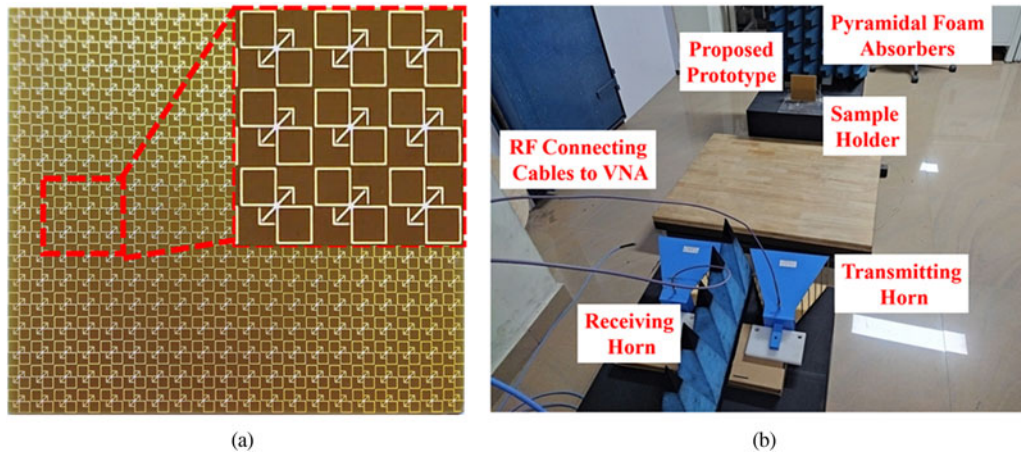


Fig. 9. (a) A 20 × 20 array of the fabricated prototype. (b) Free-space measurement setup.

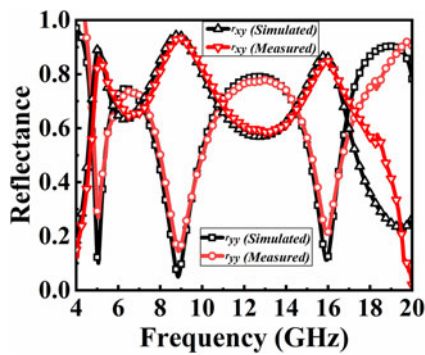


Fig. 10. Comparison of measured reflectances with the simulated results.

incident EM fields induces plasmonic resonances, which may be electric or magnetic. From Fig. 8, the induced surface currents are antiparallel for the top FSS and ground plane at 5.05 and 8.85 GHz resembling the magnetic or dielectric resonance. In contrast, the currents are parallel for the 15.95 GHz case, resulting in electric resonances [9].

### Fabrication and measurement setup

For the experimental verification, a 20 × 20 array (146 mm × 146 mm) of the proposed prototype is fabricated using traditional printed circuit board (PCB) technology, as shown in Fig. 9(a). The top FSS was printed on a thin 2.4 mm grounded FR-4

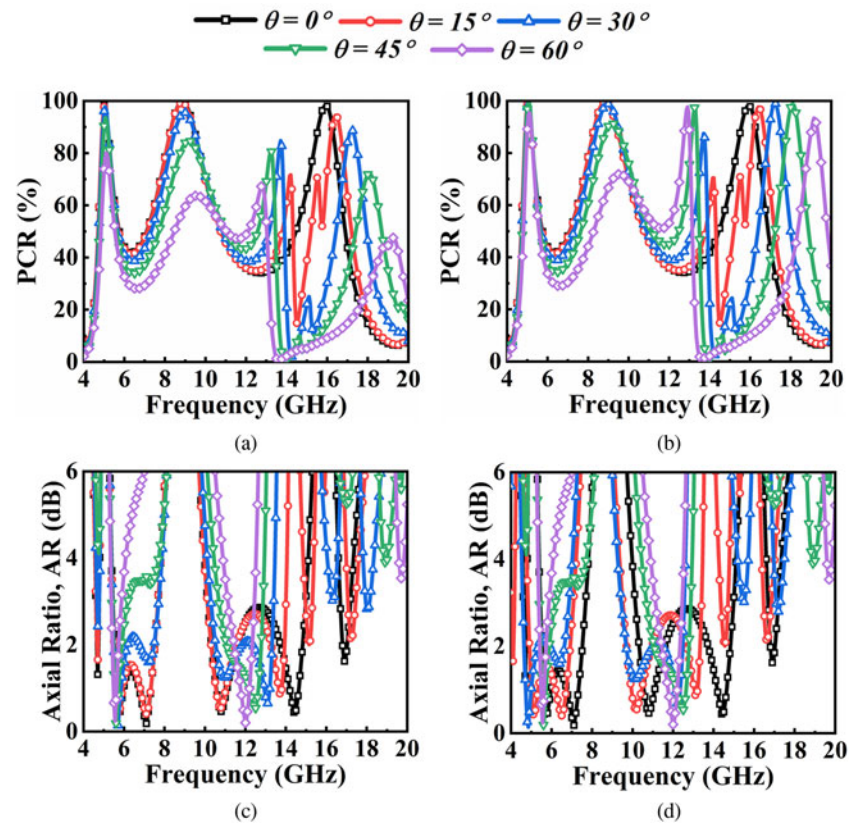


Fig. 11. Measured PCR for the (a) TE, (b) TM, and AR (dB) for the (c) TE and (d) TM modes for different elevation angles ( $\theta$ ).

**Table 2.** Performance comparison of the proposed design with other recently reported state-of-the-art polarizers

Ref	Substrate used	Periodicity (in $\lambda_L^2$ )	Thickness (in $\lambda_L$ )	Volume (in $\lambda_L^3/1000$ )	Polarization conversion	Freq range (GHz)	FBW (%)	Angular stability	Orthogonality in circular polarization
[9]	Teflon	$0.219 \times 0.219$	0.082	3.932	Linear-cross	8.2–23	94.8	–	–
[11]	Duroid 4730	$0.176 \times 0.176$	0.044	1.362	Linear-cross	4.4–5.30 9.45–13.60	18.56 36.01	–	–
[20]	FR-4	$0.175 \times 0.175$	0.040	1.225	Linear-cross Linear-circular	8–11 7.5–7.7 11.5–11.9	31.58 2.63 3.42	–	–
[21]	FR-4	$0.091 \times 0.091$	0.091	0.753	Linear-cross Linear-circular	17.4–18.9 9.1–16.5 20–25.4	8.26 57.81 23.79	30°	Yes
[22]	FR-4	$0.265 \times 0.265$	0.066	4.628	Linear-cross	12.4–27.96	77	30°	–
[23]	RO4003C	$0.233 \times 0.233$	0.071	3.850	Linear-cross	7–16	78.2	–	–
[24]	F4B-2	$0.210 \times 0.210$	0.066	2.910	Linear-cross	9.04–20.83	78.9	–	–
[25]	RT Duroid	$0.175 \times 0.175$	0.078	2.400	Linear-cross	7.5–24.5	106	Up to 30°	–
[16]	PET	$0.188 \times 0.188$	0.080	2.820	Linear-cross	11.3–20	56.5	Up to 40°	–
[26]	FRB	$0.294 \times 0.294$	0.078	6.711	Linear-circular	8.81–14.34	40.9	–	–
[27]	FR-4	$0.230 \times 0.230$	0.069	3.650	Linear-cross	6.91–14.31	69.7	–	–
[28]	F4B-2	$0.182 \times 0.182$	0.084	2.780	Linear-cross	12.4–27.96	98.9	Up to 40°	–
[29]	F4B	$0.313 \times 0.313$	0.094	9.170	Linear-cross	9.38–13.36 14.84–20.36	35 31.36	Up to 45°	–
[30]	Arlon	$0.244 \times 0.034$	0.086	0.639	Linear-circular	1.9–2.3 7.9–8.3	19.04 4.93	Up to 25°	–
[31]	Neltec	$0.032 \times 0.426$	0.069	0.922	Linear-circular	13.5–22.7 26–30.9	50.83 17.22	Up to 15°	Yes
[32]	Taconic RF-35	$0.152 \times 0.312$	0.059	2.820	Linear-circular	11.7–12.5 17.3–18.1	6.61 4.51	Up to 30°	Yes
<b>This work</b>	<b>FR-4</b>	<b><math>0.112 \times 0.112</math></b>	<b>0.037</b>	<b>0.464</b>	<b>Linear-circular</b>	<b>4.61–4.77</b>	<b>0.16</b>	<b>45°</b>	<b>Yes</b>
						<b>5.44–7.68</b>	<b>34.15</b>		
						<b>10.13–14.98</b>	<b>38.63</b>		
					<b>16.68–17.19</b>	<b>3.01</b>			
					<b>Linear-cross</b>	<b>4.92–5.18</b>	<b>5.15</b>		
					<b>8.32–9.41</b>	<b>1.09</b>			
						<b>15.56</b>	<b>16.25</b>		

\* $\lambda_L$  is the free-space wavelength corresponding to the band's lowest frequency.

substrate. The free-space measurement technique is adopted for the reflection properties extraction of the sample [12], as depicted in Fig. 9(b). Two horn antennas (800 MHz–18 GHz) for the transmission and reception were connected to Keysight Power Network Analyzer (PNA) N5224B. The calibration of the PNA is done with the perfect electric conductor (PEC) sheet having the same dimension as the prototype. The measured reflectance (co-pol and cross-pol) is compared with the simulated one, and good matching is observed (see Fig. 10). The measured PCR and AR (dB) are depicted in Fig. 10.

For the oblique TE incidence, the transmitting and receiving antennas were moved away from the normal of the designed metasurface. The incident angle is varied from 0° to 60° in steps of 15°, and the reflectance measurement was done for both co- and cross-pol components. Both the antennas were rotated by 90° for the TM case. The measured PCR and AR (dB) for TE and TM incidences are depicted in Fig. 11. The proposed polarizer's performance is compared to Table 2's other recently reported state-of-the-art designs. The authors strongly believe that this design has the potential which can be suitable for real-time applications due to its compact size and multi-conversion performance with circular rotation switching in consecutive bands.

## Conclusion

The novel figure-of-eight squared loop linked with arrow type FSS-based multiband linear-cross and linear-circular reflective type polarizer is proposed in this article. The prototype is fabricated, and a good understanding is observed with the simulated results. Surface plasmonic resonances are the reason behind these conversions. The design novelties lie in its size, multiband conversion (linear-cross and linear-circular) with switching circular rotations sense in the consecutive band, and considerable stable performance for oblique incidence, making this polarizer a better candidate for real-time needs.

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