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A note on the logarithmic derivative of the gamma function

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Introduction. The object of this note is to give simpler proofs of two formulae involving the function $\psi(z)$ which I have proved elsewhere by more complicated methods.¹

The formulae are ²

$$\psi(x+1) - \log x = 2 \int_0^{\infty} \left\{ \psi(t+1) - \log t \right\} \cos 2\pi xt \, dt \quad (1)$$

for all real positive x , and

$$\begin{aligned} & \sum_{n=1}^{\infty} \left\{ \psi(1+nz) - \log nz - \frac{1}{2nz} \right\} + \frac{1}{2z} (\gamma - \log 2\pi z) \\ &= \frac{1}{z} \sum_{n=1}^{\infty} \left\{ \psi\left(1 + \frac{n}{z}\right) - \log \frac{n}{z} - \frac{z}{2n} \right\} + \frac{1}{2} \left(\gamma - \log \frac{2\pi}{z} \right) \end{aligned} \quad (2)$$

for $|\arg z| < \pi$, where γ is Euler's constant.

Proof of (1). Now ³

¹ *Journal London Math. Soc.*, 22 (1947), 14-18. $\psi(z)$ denotes $\Gamma'(z) / \Gamma(z)$.

² The first formula shows that $\psi(x+1) - \log x$ is self-reciprocal with respect to the Fourier cosine kernel $2 \cos 2\pi x$. It is strange that this result should have been overlooked, but I can find no trace of it.

Cf. B. M. Mehrotra, *Journal Indian Math. Soc. (New Series)*, 1 (1934), 209-27 for a list of self-reciprocal functions and references.

³ C. A. Stewart, *Advanced Calculus*, (London, 1940), 495 and 497.

$$\psi(x+1) = \int_0^\infty \left(\frac{e^{-t}}{t} - \frac{e^{-xt}}{e^t-1} \right) dt \quad (3)$$

$$= \frac{1}{2x} + \log x - 2 \int_0^\infty \frac{dt}{(t^2+x^2)(e^{2\pi t}-1)}. \quad (4)$$

Also

$$\int_0^\infty (e^{-t} - e^{-xt}) \frac{dt}{t} = \log x$$

by Frullani's integral,¹ and $\int_0^\infty \frac{dt}{t^2+x^2} = \frac{\pi}{2x}$.

Combining these results with (3) and (4) we have

$$\begin{aligned} \psi(x+1) - \log x &= \int_0^\infty \left(\frac{1}{t} - \frac{1}{e^t-1} \right) e^{-xt} dt \\ &= 2 \int_0^\infty \left(\frac{1}{2\pi t} - \frac{1}{e^{2\pi t}-1} \right) \frac{t dt}{t^2+x^2}. \end{aligned} \quad (5)$$

Hence

$$\begin{aligned} 2 \int_0^\infty \left\{ \psi(t+1) - \log t \right\} \cos 2\pi x t dt &= 2 \int_0^\infty \cos 2\pi x t dt \int_0^\infty \left(\frac{1}{u} - \frac{1}{e^u-1} \right) e^{-ut} du \\ &= 2 \int_0^\infty \left(\frac{1}{u} - \frac{1}{e^u-1} \right) du \int_0^\infty e^{-ut} \cos 2\pi x t dt \\ &= 2 \int_0^\infty \left(\frac{1}{u} - \frac{1}{e^u-1} \right) \frac{u du}{u^2+4\pi^2 x^2} \\ &= 2 \int_0^\infty \left(\frac{1}{2\pi t} - \frac{1}{e^{2\pi t}-1} \right) \frac{t dt}{t^2+x^2} \\ &= \psi(x+1) - \log x \end{aligned}$$

by (5). The inversion of order of integration is justified by absolute convergence for u in the range (δ, ∞) , $(\delta > 0)$ and the result follows on making δ tend to $+0$.

Proof of (2). Let $z = p/q$ where p and q are integers and $(p, q) = 1$. Consider the expression

$$\begin{aligned} S_N(p, q) &= \frac{1}{q} \sum_{n=1}^{Nq} \left\{ \psi \left(1 + \frac{np}{q} \right) - \log \frac{np}{q} - \frac{q}{2np} \right\} \\ &\quad - \frac{1}{p} \sum_{m=1}^{Np} \left\{ \psi \left(1 + \frac{mq}{p} \right) - \log \frac{mq}{p} - \frac{p}{2mq} \right\}. \end{aligned} \quad (6)$$

$$\text{Now }^2 \quad \psi(1+py) = \log p + \frac{1}{p} \sum_{r=1}^p \psi \left(y + \frac{r}{p} \right)$$

¹ C. A. Stewart, loc. cit. 457.

² C. A. Stewart, loc. cit. 504.

hence
$$\psi\left(1 + \frac{np}{q}\right) = \log p + \frac{1}{p} \sum_{r=1}^p \psi\left(\frac{np + rq}{pq}\right).$$

Substituting this result in both parts of (6) we have

$$S_N(p, q) = \frac{1}{q} \sum_{n=1}^{Nq} \left\{ \frac{1}{p} \sum_{r=1}^p \psi\left(\frac{np + rq}{pq}\right) - \log \frac{n}{q} - \frac{q}{2np} \right\} - \frac{1}{p} \sum_{m=1}^{Np} \left\{ \frac{1}{q} \sum_{s=1}^q \psi\left(\frac{mq + sp}{pq}\right) - \log \frac{m}{p} - \frac{p}{2mq} \right\}. \quad (7)$$

Now $np + rq$ ($1 \leq n \leq Nq, \quad 1 \leq r \leq p$)
 and $mq + sp$ ($1 \leq m \leq Np, \quad 1 \leq s \leq q$)

both run through the same sets of integers in the range $p + q$ to $(N + 1)pq$, inclusive. Hence the repeated sums in (7) cancel, and

$$\begin{aligned} S_N(p, q) &= -\frac{1}{q} \sum_{n=1}^{Nq} \left(\log \frac{n}{q} + \frac{q}{2np} \right) + \frac{1}{p} \sum_{m=1}^{Np} \left(\log \frac{m}{p} + \frac{p}{2mq} \right) \\ &= N \log \frac{q}{p} + \frac{1}{p} \log (Np!) - \frac{1}{q} \log (Nq!) \\ &\quad + \frac{1}{2q} \sum_{m=1}^{Np} \frac{1}{m} - \frac{1}{2p} \sum_{n=1}^{Nq} \frac{1}{n}. \end{aligned}$$

Now $\log (M!) = (M + \frac{1}{2}) \log M - M + \frac{1}{2} \log 2\pi + O(M^{-1})$,

and $\sum_{n=1}^M \frac{1}{n} = \log M + \gamma + O(M^{-1})$

as $M \rightarrow \infty$. Hence, as $N \rightarrow \infty$

$$\begin{aligned} S_N(p, q) &= N \log \frac{q}{p} + \frac{1}{p} \left\{ (Np + \frac{1}{2}) \log Np - Np + \frac{1}{2} \log 2\pi \right\} \\ &\quad - \frac{1}{q} \left\{ (Nq + \frac{1}{2}) \log Nq - Nq + \frac{1}{2} \log 2\pi \right\} \\ &\quad + \frac{1}{2q} (\gamma + \log Np) - \frac{1}{2p} (\gamma + \log Nq) + O(N^{-1}) \\ &= \frac{1}{2q} \left(\gamma - \log \frac{2\pi q}{p} \right) - \frac{1}{2p} \left(\gamma - \log \frac{2\pi p}{q} \right) + O(N^{-1}) \\ &= \frac{1}{q} \left\{ \frac{1}{2} \left(\gamma - \log \frac{2\pi}{z} \right) - \frac{1}{2z} (\gamma - \log 2\pi z) \right\} + O(N^{-1}). \end{aligned}$$

Multiplying by q and re-arranging, we have

$$\begin{aligned} & \sum_{n=1}^{Nq} \left\{ \psi(1+nz) - \log nz - \frac{1}{2nz} \right\} + \frac{1}{2z} (\gamma - \log 2\pi z) \\ &= \frac{1}{z} \sum_{n=1}^{Np} \left\{ \psi\left(1 + \frac{n}{z}\right) - \log \frac{n}{z} - \frac{1}{2n} \right\} + \frac{1}{2} \left(\gamma - \log \frac{2\pi}{z} \right) + O(N^{-1}). \end{aligned} \quad (8)$$

Further¹ if $|z| \rightarrow \infty$ in the region $|\arg z| \leq \pi - \delta$, ($\delta > 0$) then

$$\psi(1+z) - \log z - \frac{1}{2z} \sim -\frac{1}{12z^2}.$$

Hence the two series in (8) converge absolutely as $N \rightarrow \infty$, and thus (2) holds for positive rational z . Further, both the series in (2) converge absolutely for $|\arg z| < \pi$, and they define analytic functions of z in this region. Hence (2) holds by analytic continuation for $|\arg z| < \pi$.

Extensions of (2). The method of the previous section can be used to prove that

$$z \sum_{n=1}^{\infty} \left\{ \psi'(1+nz) - \frac{1}{nz} \right\} - \frac{1}{z} \sum_{n=1}^{\infty} \left\{ \psi'\left(1 + \frac{n}{z}\right) - \frac{1}{n} \right\} = \log z$$

for $|\arg z| < \pi$. For higher derivatives ($k \geq 2$) the corresponding results are

$$\sum_{n=1}^{\infty} \psi^{(k)}(1+nz) = z^{-k-1} \sum_{n=1}^{\infty} \psi^{(k)}\left(1 + \frac{n}{z}\right). \quad (9)$$

In these cases the results follow immediately on substituting²

$$\psi^{(k)}(1+ny) = (-1)^{k-1} k! \sum_{m=1}^{\infty} (m+ny)^{-k-1}$$

in each side of (9).

¹ E. T. Whittaker and G. N. Watson, *Modern Analysis* (Cambridge, 1927), 278.

² C. A. Stewart, loc. cit. 505.

Note added in proof.

Professor T. A. Brown tells me that he proved the self-reciprocal property of $\psi(1+x) - \log x$ some years ago, and that he communicated the result to the late Professor G. H. Hardy. Professor Hardy said that the result was also given in a progress report to the University of Madras by S. Ramanujan, but was not published elsewhere.

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