

A CLASS OF CRITICAL KIRCHHOFF PROBLEM ON THE HYPERBOLIC SPACE \mathbb{H}^n

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Abstract. We investigate questions on the existence of nontrivial solution for a class of the critical Kirchhoff-type problems in Hyperbolic space. By the use of the stereographic projection the problem becomes a singular problem on the boundary of the open ball $B_1(0) \subset \mathbb{R}^n$. Combining a version of the Hardy inequality, due to Brezis–Marcus, with the mountain pass theorem due to Ambrosetti–Rabinowitz are used to obtain the nontrivial solution. One of the difficulties is to find a range where the Palais Smale converges, because our equation involves a nonlocal term coming from the Kirchhoff term.

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1. Introduction. In this paper, we investigate questions on the existence of nontrivial solution for the following Kirchhoff-type equation

$$-\left(a + b \int_{\mathbf{B}^3} |\nabla_{\mathbf{B}^3} u|^2 dV_{\mathbf{B}^3}\right) \Delta_{\mathbf{B}^3} u = \lambda |u|^{q-2} u + |u|^4 u \quad \text{in } u \in H^1(\mathbf{B}^3) \quad (1.1)$$

in Hyperbolic space \mathbf{B}^3 , where a , b , and λ are positive constants, $4 < q < 6$, $H^1(\mathbf{B}^3)$ is the usual Sobolev space on the disc model of the Hyperbolic space \mathbf{B}^3 , and $\Delta_{\mathbf{B}^3}$ denotes the Laplace Beltrami operator on \mathbf{B}^3 . For the hyperbolic space \mathcal{H}^n , we make use of the stereographic projection $E: \mathcal{H}^n \rightarrow \mathbb{R}^n$, where each point $P' \in \mathcal{H}^n$ is projected to $P \in \mathbb{R}^n$, where P is the intersection of the straight line connecting P' and the point $(0, \dots, 0, -1)$. More exactly, we have explicitly the projection operator $G: \mathbb{R}^n \rightarrow \mathcal{H}^n$ and $G^{-1}: \mathcal{H}^n \rightarrow \mathbb{R}^n$ given by

$$G(x) = (x.p(x), (1 + |x|^2)p/2) \quad \text{and} \quad G^{-1}(y) = \frac{1}{y_{n+1}} y, \quad x, y \in \mathbb{R}^n,$$

$$\text{where } p(x) = \frac{2}{1 - |x|^2}.$$

This projection takes \mathcal{H}^n onto the open ball $B_1(0) \subset \mathbb{R}^n$, and we denote by $D \subset B_1(0)$ the stereographic projection of $D' \subset \mathcal{H}^n$. See more details in the excellent books [54, 58].

We will consider the metric

$$ds = p(x)|dx|, \quad \text{where } p(x) = \frac{2}{1 - |x|^2},$$

and we denote by \mathbf{B}^n the ball $B_1(0)$ endowed with the above metric.

The gradient, the Dirichlet integral, and the Laplace–Beltrami operator corresponding to this metric are

$$\begin{aligned} \nabla_{\mathbf{B}^n} u &= \frac{\nabla u}{p^2}, & Du &= \int_D |\nabla_{\mathbf{B}^n} u|^2 dV_{\mathbf{B}^n} = \int_D |\nabla u|^2 p^{n-2} dx, \\ \Delta_{\mathbf{B}^n} u &= p^{-n} \operatorname{div}(p^{n-2} \nabla u). \end{aligned}$$

See [13, 37, 57].

Elliptic problems, in an euclidean space, involving Sobolev's critical exponent, that is, when the non-linearity behaves as a polynomial function of degree $2^* = \frac{2N}{N-2}$, $N \geq 3$, were studied in a pioneering and remarkable article due to Brezis and Nirenberg [19]. In that paper, the lack of compactness was overcome by analyzing the critical level set of the functional associated with the problem. From this work, several authors have been working on the theme trying to extend or complement existing results in several directions, in that sense, we would like to mention some articles, and we apologize for not mentioning all the authors. For such problems modeled in a bounded domain, we cite [3, 5, 10, 14, 16, 20, 23, 24, 29, 32, 56, 62], while we mention [33, 51, 50, 63] for problems in unbounded domains. For the problems involving the p -Laplacian operator or more general degenerate operators, the following works have treated these subjects, [21, 38–40, 51]. The authors in [8, 9, 12] and also in [11, 13, 15, 22, 36, 37, 41, 49, 57] have treated some critical problem in a sphere and in a hyperbolic space, respectively. We cite [15, 41] for the related problems in the cases linear or supercritical. See references therein, as well as the book [61] for additional remarks and results. On the other hand, our equation in an euclidean space is related to a stationary Kirchhoff equation [45], namely,

$$u_{tt} - M \left(\int_{\Omega} |\nabla_x u|^2 dx \right) \Delta_x u = f(x, t), \quad (x, t) \in \Omega \times \mathbb{R},$$

where Ω is bounded domain, $M(s) = a + bs$, $a, b > 0$, and f is a suitable function, which is an extension of the classical D'Alembert's wave equation, since in this case, the model considers the effects of the changes in the length of the strings during the vibrations. See [46]. The main difficulty is because the term containing M in the equation makes this equation nonlocal, that is, the equation does not satisfy a pointwise identity any longer.

The above equation has been received special attention after the work by Lions [46], where a non-linear functional analysis approach was proposed. Up to our knowledge, Ma and Rivera [48] were the pioneers to study this problem by variational methods, more exactly, by using the minimization method. In [1], was employed the mountain pass theorem, while in [53] a topologic argument was used, more specifically, the Yang index and critical groups, and in [44] is studied the equation by using the minimization arguments and fountain theorem. We would like to cite [26, 35, 60] for more multiplicity results. For the Kirchhoff equation involving critical exponents we refer to [2, 34, 42, 43, 47] and references therein. See [4, 7, 28, 30] and [25, 27, 55] for some related results.

Returning to our subject, restricted to $n = 3$, if u is a solution of equation (1.1), putting $v := p^{\frac{1}{2}}u$, then v satisfies the following problem:

$$\begin{cases} (a + b\|v\|^2) (-\Delta v + (3/4)p^2v) = \lambda p^\alpha |v|^{q-2}v + |v|^4v, & \text{in } B_1(0) \\ v = 0, & \text{on } \partial B_1(0), \end{cases} \tag{1.2}$$

where $\alpha = (6 - q)/2$ and $\|v\|^2 = \int_{B_1(0)} (|\nabla v|^2 + (3/4)p^2v^2)$.

From now on, we will consider $\Omega := B_1(0)$. We will denote by $H_{0,r}^1(\Omega)$ the subspace of $H_0^1(\Omega)$ of the radial functions which is endowed with the norm given by

$$\|v\|^2 = \int_{\Omega} (|\nabla v|^2 + (3/4)p^2v^2).$$

Since the euclidean sphere with center at the origin $0 \in \mathbb{R}^N$ is also a hyperbolic sphere with center at the origin $0 \in \mathbf{B}^n$, $H_{0,r}^1(\Omega)$ can also be seen as the subspace of $H_0^1(\Omega)$ consisting of the hyperbolic radial functions. See this characterization as well as others remarks in [13, Appendix]. We observe that in [13, Theorem 3.1], $H_{0,r}^1(\Omega)$ is embedded compactly in $L^q(\Omega)$ for $2 < q < 2^*$. Note also that here $2^* = 6$.

We have the following functional $J : H_{0,r}^1(\Omega) \rightarrow \mathbb{R}$ associated with problem (1.2)

$$J(v) = \frac{a}{2}\|v\|^2 + \frac{b}{4}\|v\|^4 - \frac{\lambda}{q} \int_{\Omega} p^\alpha |v|^q - \frac{1}{6} \int_{\Omega} |v|^6, \tag{1.3}$$

whose Gateaux derivative is given by

$$J'(v)w = (a + b\|v\|^2) \int_{\Omega} \left(\nabla v \nabla w + \frac{3}{4}p^2vw \right) - \lambda \int_{\Omega} p^\alpha |v|^{q-2}vw - \int_{\Omega} |v|^4vw. \tag{1.4}$$

Now, we present our main result.

THEOREM 1.1. *Suppose $4 < q < 6$. Then, for every $\lambda > 0$ the problem (1.1) has a nontrivial solution $u \in H^1(\mathbf{B}^3)$.*

This result of existence of nontrivial solution, in the hyperbolic space, extends results presented in [22, 43].

2. Proof of the main result. The proof is made by applying the mountain pass theorem (See Willem [61] for a reference). To this end, we have the following

LEMMA 2.1. *(Mountain pass geometry).*

- (a) *There exist $\beta > 0$ and $\rho > 0$ such that $J(v) \geq \beta$ when $\|v\| = \rho$.*
- (b) *There exists an element $e \in H_{0,r}^1(\Omega)$ with $\|e\| > \rho$ such that $J(e) < 0$.*

Proof. For item (a), we observe that by [17, 18], there exists a constant $C > 0$, such that

$$\int_{\Omega} p^\alpha v^q \leq C \left(\int_{\Omega} |\nabla v|^2 \right)^{\frac{q}{2}} \leq C \left[\int_{\Omega} (|\nabla v|^2 + (3/4)p^2v^2) \right]^{\frac{q}{2}}.$$

Thus,

$$J(u) \geq \frac{a}{2} \|v\|^2 + \frac{b}{4} \|v\|^4 - \frac{C\lambda}{q} \left[\int_{\Omega} (|\nabla v|^2 + (3/4)p^2 v^2) \right]^{\frac{q}{2}} - \frac{1}{6} \int_{\Omega} |v|^6,$$

and by the Sobolev continuous embedding, there is a constant $\tilde{C} > 0$, verifying

$$J(u) \geq \frac{a}{2} \|v\|^2 + \frac{b}{4} \|v\|^4 - \frac{C\lambda}{q} \|v\|^q - \frac{\tilde{C}}{6} \|v\|^6 \geq \beta,$$

where the conclusion follows by making $\|v\| = \rho$ sufficiently small.

Now for item (b), take $0 < v \in H_{0,r}^1(\Omega)$ and $0 < t$. Note that

$$J(tv) = \frac{at^2}{2} \|v\|^2 + \frac{bt^4}{4} \|v\|^4 - \frac{\lambda t^q}{q} \int_{\Omega} p^\alpha |v|^q - \frac{t^6}{6} \int_{\Omega} |v|^6.$$

Therefore, $J(tv) \rightarrow -\infty$ as $t \rightarrow +\infty$. Consequently, J satisfies the geometry of the mountain pass theorem. □

Lemma 2.1 and Ekeland’s Variational Principle [6] allow us to use the general minimax principle, see [61, Theorem 2.9], which gives us a Palais-Smale sequence, $(v_k) \subset H_{0,r}^1(\Omega)$, at the level c , i.e.,

$$J(v_k) \rightarrow c \text{ and } \|J'(v_k)\|_{H_{0,r}^1(\Omega)^*} \rightarrow 0, \text{ as } k \rightarrow \infty, \tag{2.5}$$

where

$$c = \inf_{\gamma \in \Gamma} \sup_{t \in [0,1]} J(\gamma(t)),$$

and $\Gamma = \{ \gamma \in C([0, 1], H_{0,r}^1(\Omega)); \gamma(0) = 0, J(\gamma(1)) < 0 \}$.

LEMMA 2.2. *The sequence $(v_k) \subset H_{0,r}^1(\Omega)$ defined above is bounded.*

Proof. Since (v_k) is a Palais-Smale sequence at the level c ,

$$c + 1 + \|v_k\| \geq J(v_k) - \frac{1}{q} J'(v_k)v_k.$$

Thus, by Sobolev continuous embedding, there is a constant $\tilde{C}' > 0$, such that

$$c + 1 + \|v_k\| \geq \left(\frac{a}{2} - \frac{a}{q} \right) \|v_k\|^2 + \left(\frac{b}{4} - \frac{b}{q} \right) \|v_k\|^4 + \tilde{C}' \left(\frac{1}{q} - \frac{1}{6} \right) \|v_k\|^6.$$

Therefore, the sequence is bounded. □

LEMMA 2.3. *We have $c < \frac{1}{4}abS^3 + \frac{1}{24}b^3S^6 + \frac{1}{24}(b^2S^4 + 4aS)^{3/2}$, where*

$$S := \inf_{u \in H_{0,r}^1(\Omega)} \frac{\int_{\Omega} |\nabla u|^2}{\left(\int_{\Omega} u^6 \right)^{1/3}}.$$

Proof. In this proof, we will follow some of the arguments made in [19], see also, for instance, [43, 42, 50]. First, we observe that it suffices to show that there exists a $v_0 \in H_{0,r}^1(\Omega)$, $v_0 \neq 0$ such that

$$\sup_{t \geq 0} J(tv_0) < \frac{1}{4}abS^3 + \frac{1}{24}b^3S^6 + \frac{1}{24}(b^2S^4 + 4aS)^{3/2}. \tag{2.6}$$

Indeed, observing that $J(tv_0) \rightarrow -\infty$ as $t \rightarrow \infty$, there exists $R > 0$ such that $J(Rv_0) < 0$. Now, we write $u_1 := Rv_0$, and from Lemma 2.1, we have

$$0 < \beta \leq c = \inf_{\gamma \in \Gamma} \max_{\tau \in [0,1]} J(\gamma(\tau)) \leq \sup_{t \geq 0} J(tv_0) < \frac{1}{4}abS^3 + \frac{1}{24}b^3S^6 + \frac{1}{24}(b^2S^4 + 4aS)^{3/2}.$$

Therefore, we are going to prove the existence of a function v_0 such that (2.6) holds.

Let $0 < R < \frac{1}{2}$ be fixed, and let $\varphi \in C_0^\infty(\Omega)$ be a cut-off function with support at B_{2R} , such that φ is identically 1 on B_R and $0 \leq \varphi \leq 1$ on B_{2R} , where B_r denotes the ball in \mathbb{R}^3 with center at the origin and radius r .

Given $\varepsilon > 0$, we set $\psi_\varepsilon(x) := \varphi(x)\omega_\varepsilon(x)$, where

$$\omega_\varepsilon(x) = (3\varepsilon)^{\frac{1}{4}} \frac{1}{(\varepsilon + |x|^2)^{\frac{1}{2}}},$$

and ω_ε satisfies (see[59])

$$\int_{\mathbb{R}^3} |\nabla \omega_\varepsilon|^2 = \int_{\mathbb{R}^3} |\omega_\varepsilon|^6 = S^{1/2}. \tag{2.7}$$

From the definition of ω_ε , it can be shown that

$$\int_{B_R} |\nabla \omega_\varepsilon|^2 \leq \int_{B_R} |\omega_\varepsilon|^6, \tag{2.8}$$

and

$$\int_{B_1 - B_R} |\nabla \psi_\varepsilon|^2 = O(\varepsilon^{\frac{1}{2}}) \text{ as } \varepsilon \rightarrow 0. \tag{2.9}$$

Now, we define

$$v_\varepsilon := \frac{\psi_\varepsilon}{\left(\int_{B_{2R}} \psi_\varepsilon^6\right)^{1/6}}$$

and also $X_\varepsilon := \int_{B_1} |\nabla v_\varepsilon|^2$. Then, we have

$$\begin{aligned} X_\varepsilon &= \int_{\Omega} |\nabla v_\varepsilon|^2 = \int_{\Omega} \frac{|\nabla \psi_\varepsilon|^2}{B^2} \\ &= \int_{B_R} \frac{|\nabla \psi_\varepsilon|^2}{B^2} + \int_{B_{2R} - B_R} \frac{|\nabla \psi_\varepsilon|^2}{B^2}, \end{aligned}$$

where $B := \left(\int_{B_{2R}} \psi_\varepsilon^6\right)^{1/6}$. Then, since, $\varphi \equiv 1$ and consequently $\nabla \varphi \equiv 0$ on B_R , we have

$$\begin{aligned} X_\varepsilon &= \frac{1}{B^2} \int_{B_R} |\nabla \psi_\varepsilon|^2 + \frac{1}{B^2} \int_{B_{2R} - B_R} |\nabla \psi_\varepsilon|^2 \\ &= \frac{1}{B^2} \int_{B_R} |\nabla \omega_\varepsilon|^2 + \int_{B_{2R} - B_R} |\nabla \psi_\varepsilon|^2. \end{aligned}$$

By equations (2.8) and (2.9) and considering $\delta = \frac{1}{2}$ we obtain

$$\begin{aligned} X_\varepsilon &\leq \frac{1}{B^2} \int_{B_R} |\omega_\varepsilon|^6 + O(\varepsilon^\delta) = \frac{\int_{B_R} |\omega_\varepsilon|^6}{\left(\int_{B_R} |\omega_\varepsilon|^6 + \int_{B_{2R}-B_R} |\psi_\varepsilon|^6\right)^{1/3}} + O(\varepsilon^\delta) \\ &\leq \frac{\int_{B_R} |\omega_\varepsilon|^6}{\left(\int_{B_R} |\omega_\varepsilon|^6\right)^{1/3}} + O(\varepsilon^\delta) \leq \left(\int_{B_R} |\omega_\varepsilon|^6\right)^{2/3} + O(\varepsilon^\delta) \\ &\leq \left(\int_{\mathbb{R}^3} |\omega_\varepsilon|^6\right)^{2/3} + O(\varepsilon^\delta) = S + O(\varepsilon^\delta) \end{aligned}$$

Therefore, we have

$$X_\varepsilon \leq S + O(\varepsilon^\delta). \tag{2.10}$$

On the other hand, we have

$$\lim_{t \rightarrow +\infty} J(tv_\varepsilon) = -\infty, \forall \varepsilon > 0.$$

This implies that there exists $t_\varepsilon > 0$ such that $\sup_{t \geq 0} J(tv_\varepsilon) = J(t_\varepsilon v_\varepsilon)$. Now, we will prove an estimate for this t_ε .

$$J'(tv_\varepsilon)|_{t=t_\varepsilon} = 0.$$

Thus,

$$at_\varepsilon \|v_\varepsilon\|^2 + bt_\varepsilon^3 \|v_\varepsilon\|^4 - \lambda t_\varepsilon^{q-1} \int_\Omega p^\alpha |v_\varepsilon|^q - t_\varepsilon^5 \int_\Omega |v_\varepsilon|^6 = 0,$$

which implies

$$a \|v_\varepsilon\|^2 + bt_\varepsilon^2 \|v_\varepsilon\|^4 - \lambda t_\varepsilon^{q-2} \int_\Omega p^\alpha |v_\varepsilon|^q - t_\varepsilon^4 \int_\Omega |v_\varepsilon|^6 = 0,$$

Since $\int_\Omega |v_\varepsilon|^6 = 1$, we have

$$-a \|v_\varepsilon\|^2 - bt_\varepsilon^2 \|v_\varepsilon\|^4 + t_\varepsilon^4 \leq 0.$$

Hence

$$0 \leq t_\varepsilon^2 \leq \frac{b \|v_\varepsilon\|^4 + [(b \|v_\varepsilon\|^4)^2 + 4a \|v_\varepsilon\|^2]^{1/2}}{2} := t_0.$$

Since the function $t \mapsto \frac{a}{2} t^2 \|v_\varepsilon\|^2 + \frac{b}{4} t^4 \|v_\varepsilon\|^4 - \frac{t^6}{6}$ is increasing on $[0, t_0]$, denoting $C_1 = a \|v_\varepsilon\|^2$ and $C_2 = b \|v_\varepsilon\|^4$, we have

$$\begin{aligned}
 J(t_\varepsilon v_\varepsilon) &\leq \frac{at_0}{2} \|v_\varepsilon\|^2 + \frac{bt_0^2}{4} \|v_\varepsilon\|^4 - \frac{\lambda t_\varepsilon^q}{q} \int_\Omega p^\alpha v_\varepsilon^q - \frac{t_0^3}{6} \\
 &\leq \frac{t_0 C_1}{2} + \frac{t_0^2 C_2}{4} - \frac{\lambda t_\varepsilon^q}{q} \int_\Omega p^\alpha v_\varepsilon^q - \frac{t_0^3}{6} \\
 &\leq \frac{1}{2} \left[\frac{C_2 + (C_2^2 + 4C_1)^{1/2}}{2} \right] C_1 + \frac{1}{4} \left[\frac{C_2 + (C_2^2 + 4C_1)^{1/2}}{2} \right]^2 C_2 \\
 &\quad - \frac{1}{6} \left[\frac{C_2 + (C_2^2 + 4C_1)^{1/2}}{2} \right]^3 - \frac{\lambda t_\varepsilon^q}{q} \int_\Omega p^\alpha v_\varepsilon^q \\
 &\leq \frac{C_1 C_2}{4} + \frac{C_2^3}{24} + \frac{1}{24} (C_2^2 + 4C_1)^{3/2} - \frac{\lambda t_\varepsilon^q}{q} \int_\Omega p^\alpha v_\varepsilon^q.
 \end{aligned}$$

Considering $A = 3/4 \int_\Omega p^2 v_\varepsilon^2$, by definition of the norm, and the inequality (2.10), we obtain

$$\begin{aligned}
 J(t_\varepsilon v_\varepsilon) &\leq \frac{ab}{4} (X_\varepsilon + A)^3 + \frac{b^3}{24} (X_\varepsilon + A)^6 + \frac{1}{24} [b^2 (X_\varepsilon + 4)^4 + 4a(X_\varepsilon + A)]^{3/2} \\
 &\quad - \frac{\lambda t_\varepsilon^q}{q} \int_\Omega p^\alpha v_\varepsilon^q \\
 &\leq \frac{ab}{4} (S + O(\varepsilon^{1/2}) + A)^3 + \frac{b^3}{24} (S + O(\varepsilon^{1/2}) + A)^6 \\
 &\quad + \frac{1}{24} [b^2 (S + O(\varepsilon^{1/2}) + A)^4 + 4a(S + O(\varepsilon^{1/2}) + A)]^{3/2} - \frac{\lambda t_\varepsilon^q}{q} \int_\Omega p^\alpha v_\varepsilon^q.
 \end{aligned}$$

By using several times the standard inequality (see, e.g., [50, p. 778])

$$(a + b)^\beta \leq a^\beta + \beta(a + b)^{\beta-1}b, \quad \forall \beta \geq 1, \forall a, b > 0,$$

we infer that

$$J(t_\varepsilon v_\varepsilon) \leq \frac{abS^3}{4} + \frac{b^3S^6}{24} + \frac{1}{24} (b^2S^4 + 4aS)^{3/2} + O(\varepsilon^{1/2}) + \int_{B_{2R}} \left(\frac{3C}{4} p^2 v_\varepsilon^2 - \lambda C_\varepsilon p^\alpha v_\varepsilon^q \right), \tag{2.11}$$

for some constant $C > 0$, where $C_\varepsilon = \frac{t_\varepsilon^q}{q}$.

At this point, we can assume that there exists a positive constant C_0 such that $C_\varepsilon \geq C_0 > 0, \forall \varepsilon > 0$. If that was not the case, we could find a sequence $\varepsilon_k \rightarrow 0$ as $k \rightarrow \infty$, such that $t_{\varepsilon_k} \rightarrow 0$ as $k \rightarrow \infty$, since $C_\varepsilon \geq 0$. Now, passing to a subsequence, if necessary, which still denoted by ε_k , we have $t_{\varepsilon_k} v_{\varepsilon_k} \rightarrow 0$ as $k \rightarrow \infty$.

Therefore,

$$0 < c \leq \sup_{t \geq 0} J(tv_{\varepsilon_k}) = J(t_{\varepsilon_k} v_{\varepsilon_k}) = J(0) = 0,$$

which is a contradiction.

Observing that $\int_{B_{2R}} p^2 v_\varepsilon^2 < \infty$, we claim

$$\text{Claim: } \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon^{1/2}} \int_{B_{2R}} \left(\frac{3C}{4} p^2 v_\varepsilon^2 - C_\varepsilon \lambda p^\alpha v_\varepsilon^q \right) = -\infty.$$

Assuming the Claim is proved, from equation (2.11), we have

$$J(t_\varepsilon v_\varepsilon) < \frac{abS^3}{4} + \frac{b^3S^6}{24} + \frac{1}{24}(b^2S^4 + 4aS)^{3/2},$$

for some $\varepsilon > 0$ sufficiently small, and the proof is complete.

Now, we are going to prove the Claim. For this, it is sufficient to show that

$$\lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon^{1/2}} \left(\int_{B_R} \left(\frac{3C}{4} p^2 \omega_\varepsilon^2 - C_\varepsilon \lambda p^\alpha \omega_\varepsilon^q \right) \right) = -\infty \tag{2.12}$$

and

$$\int_{B_{2R}-B_R} \left(\frac{3C}{4} p^2 v_\varepsilon^2 - C_\varepsilon \lambda p^\alpha v_\varepsilon^q \right) = O(\varepsilon^{1/2}). \tag{2.13}$$

First, we will consider

$$\begin{aligned} J_\varepsilon &= \frac{1}{\varepsilon^{1/2}} \int_{B_R} \left(\frac{3C}{4} p^2 \omega_\varepsilon^2 - C_\varepsilon \lambda p^\alpha \omega_\varepsilon^q \right) \\ &= \frac{3C}{4\varepsilon^{1/2}} \int_{B_R} \left(\frac{2}{1-|x|^2} \right)^2 \frac{(3\varepsilon)^{1/2}}{(\varepsilon+|x|^2)} - \frac{\lambda C_\varepsilon}{\varepsilon^{1/2}} \int_{B_R} \left(\frac{2}{1-|x|^2} \right)^\alpha \frac{(3\varepsilon)^{q/4}}{(\varepsilon+|x|^2)^{q/2}} \\ &= \tilde{C} \int_{B_R} \left(\frac{2}{1-|x|^2} \right)^2 \frac{1}{(\varepsilon+|x|^2)} - \lambda \tilde{C}_\varepsilon \varepsilon^{\frac{(q-2)}{4}} \int_{B_R} \left(\frac{2}{1-|x|^2} \right)^\alpha \frac{1}{(\varepsilon+|x|^2)^{q/2}} \\ &= J_1 - J_2, \text{ for some constant } \tilde{C} > 0. \end{aligned} \tag{2.14}$$

We observe that on B_R ,

$$2 < \frac{2}{1-|x|^2} \leq \frac{2}{1-R^2}. \tag{2.15}$$

Therefore, making the change of variables $x = \varepsilon^{1/2}y$ and then using the polar coordinates, we obtain, for some constant $\tilde{C} > 0$,

$$\begin{aligned} J_1 &\leq \frac{4\tilde{C}}{(1-R^2)^2} \int_{B_R} \frac{1}{(\varepsilon+|x|^2)} = \frac{4\tilde{C}}{(1-R^2)^2} \int_{B_{R\varepsilon^{-1/2}}} \frac{\varepsilon^{3/2}}{(\varepsilon+\varepsilon|y|^2)} \\ &= \frac{4\tilde{C}}{(1-R^2)^2} \omega \varepsilon^{1/2} \int_0^{R\varepsilon^{-1/2}} \frac{r^2}{(1+r^2)} dr. \end{aligned} \tag{2.16}$$

Now, for J_2 , we have, considering again equation (2.15), the change of variables $x = \varepsilon^{1/2}y$ and then using the polar coordinates that, we get for some constant $\tilde{C}_\varepsilon > 0$,

$$\begin{aligned} J_2 &\geq \lambda \tilde{C}_\varepsilon \varepsilon^{\frac{(q-2)}{4}} \int_{B_R} \left(\frac{2}{1-|x|^2} \right)^\alpha \frac{1}{(\varepsilon+|x|^2)^{q/2}} \\ &\geq \lambda \tilde{C}_\varepsilon \varepsilon^{\frac{(q-2)}{4}} 2^\alpha \int_{B_{R\varepsilon^{-1/2}}} \frac{\varepsilon^{3/2}}{(\varepsilon+\varepsilon|y|^2)^{q/2}} \\ &= \lambda \tilde{C}_\varepsilon 2^\alpha \varepsilon^{\frac{(q-2)}{4}} \frac{\varepsilon^{3/2}}{\varepsilon^{q/2}} \int_{B_{R\varepsilon^{-1/2}}} \frac{1}{(1+|y|^2)^{q/2}} \\ &= \lambda \tilde{C}_\varepsilon 2^\alpha \omega \varepsilon^{-\frac{q}{4}+1} \int_0^{R\varepsilon^{-1/2}} \frac{r^2}{(1+r^2)^{q/2}} dr. \end{aligned} \tag{2.17}$$

Thus, combining equations (2.14), (2.16), and (3.24), we obtain

$$J_\varepsilon \leq \frac{4\tilde{C}}{(1-R^2)^2} \omega \varepsilon^{1/2} \int_0^{R\varepsilon^{-1/2}} \frac{r^2}{(1+r^2)} dr - \lambda \tilde{C}_\varepsilon 2^\alpha w \varepsilon^{-\frac{q}{4}+1} \int_0^{R\varepsilon^{-1/2}} \frac{r^2}{(1+r^2)^{q/2}} dr.$$

Observing that

$$\int_0^{R\varepsilon^{-1/2}} \frac{r^2}{1+r^2} dr = R\varepsilon^{-1/2} - \tan^{-1}(R\varepsilon^{-1/2})$$

and $\lim_{\varepsilon \rightarrow 0^+} \varepsilon^{-\frac{q}{4}+1} = \infty$ as $4 < q$, we conclude that equation (2.12) holds.

Now, we will prove equation (2.13). First, we observe that we can find and fix an $\varepsilon > 0$ sufficiently small such that $O(\varepsilon^\delta) + \varepsilon^\delta I_\varepsilon < 0$. As in [19], we obtain

$$\int_{B_{2R}} |\psi_\varepsilon|^6 = 3^{3/2} \int_{\mathbb{R}^3} \frac{dx}{(1+|x|^2)^3} + O(\varepsilon^{3/2}). \tag{2.18}$$

From equation (2.18), we obtain

$$\frac{1}{\varepsilon^{1/2}} \int_{B_{2R}-B_R} \left(\frac{3C}{4} p^2 v_\varepsilon^2 - \lambda C_\varepsilon p^\alpha v_\varepsilon^q \right) \leq \frac{C'}{\varepsilon^{1/2}} \int_{B_{2R}-B_R} p^2 \varphi^2 \omega_\varepsilon^2.$$

We define $\Theta = B_{2R} - B_R$. Since $R \leq |x| \leq 2R$, we have

$$\frac{2}{1-R^2} \leq p(x) \leq \frac{2}{1-4R^2},$$

therefore,

$$I_1 := \frac{C'}{\varepsilon^{1/2}} \int_{\Theta} p^2 \varphi^2 \omega_\varepsilon^2 \leq \frac{4C'}{\varepsilon^{1/2}(1-4R^2)^2} \int_{\Theta} \varphi^2 \frac{\varepsilon^{1/2}}{(\varepsilon+|x|^2)}.$$

Making the change of variables $x = \varepsilon^{1/2}y$ and later changing to polar coordinates we obtain

$$\begin{aligned} I_1 &\leq \frac{4C'}{(1-4R^2)^2} \int_{\Theta'} \varphi^2(\varepsilon^{1/2}y) \frac{\varepsilon^{3/2}}{(\varepsilon+\varepsilon|y|^2)} \\ &\leq \frac{4C' \omega \varepsilon^{1/2}}{(1-4R^2)^2} \int_{R\varepsilon^{-1/2}}^{2R\varepsilon^{-1/2}} \frac{r^2}{(1+r^2)} dr, \end{aligned}$$

where $\Theta' = B_{2R\varepsilon^{-1/2}} - B_{R\varepsilon^{-1/2}}$.

By the Mean Value Theorem for integrals, there exists $r_0 \in [R\varepsilon^{-1/2}, 2R\varepsilon^{-1/2}]$ such that

$$\begin{aligned} I_1 &\leq \left[\frac{4C' \omega \varepsilon^{1/2}}{(1-4R^2)^2} \right] \frac{r_0^2}{(1+r_0^2)} (2R\varepsilon^{-1/2} - R\varepsilon^{-1/2}) = \left[\frac{4C' \omega \varepsilon^{1/2}}{(1-4R^2)^2} \right] \frac{Rr_0^2 \varepsilon^{-1/2}}{(1+r_0^2)} \\ &\leq \left[\frac{4C' \omega \varepsilon^{1/2}}{(1-4R^2)^2} \right] \frac{R(2R\varepsilon^{-1/2})^2 \varepsilon^{-1/2}}{(1+(R\varepsilon^{-1/2})^2)} = \left[\frac{4C' \omega}{(1-4R^2)^2} \right] \frac{2^2 R^3 \varepsilon^{-1}}{(1+\frac{R^2}{\varepsilon})}. \end{aligned}$$

Therefore,

$$I_1 \leq \frac{C(R)}{\varepsilon + R^2}.$$

Since $0 < \varepsilon \leq 1$, then

$$\frac{1}{1 + R^2} \leq \frac{1}{\varepsilon + R^2} \leq \frac{1}{R^2}.$$

Therefore,

$$I_1 \leq \frac{C(R)}{R^2}.$$

□

3. Proof of the theorem. Taking the sequence $\{v_n\}$ given by equation (2.5), by Lemma 2.2 this sequence $\{v_n\}$ is bounded in $H_{0,r}^1(\Omega)$. So that, we can assume, passing to a subsequence, that $v_n \rightharpoonup v$, weakly in $H_{0,r}^1(\Omega)$, as $n \rightarrow \infty$ and

$$J'(v_n)w = o(1), \forall w \in H_{0,r}^1(\Omega). \tag{3.19}$$

Now, note that

$$|J'(v_n)w - J'(v)w| \rightarrow 0, \tag{3.20}$$

as $n \rightarrow \infty$, for all $w \in C_{c,rad}^\infty(\Omega)$. From this, it follows that $J'(v)w = 0$, for all $w \in C_{c,rad}^\infty(\Omega)$. By density we conclude that

$$J'(v)w = 0, \forall w \in H_{0,r}^1(\Omega), \tag{3.21}$$

and v is a critical point of the functional J restricted to the space $H_{0,r}^1(\Omega)$.

Now, we will follow the ideas of [14, 21, 31] (see also [52]). Since $H_{0,r}^1(\Omega)$ is a closed subspace of $H_0^1(\Omega)$, we can write

$$H_0^1(\Omega) = H_{0,r}^1(\Omega) \oplus H_{0,r}^1(\Omega)^\perp,$$

where \cdot^\perp denotes the orthogonal complement of the space. Therefore, for each $w \in H_0^1(\Omega)$, there exist $\vartheta \in H_{0,r}^1(\Omega)$ and $\vartheta^\perp \in H_{0,r}^1(\Omega)^\perp$ such that

$$w = \vartheta + \vartheta^\perp. \tag{3.22}$$

As $H_{0,r}^1(\Omega)$ is a Hilbert space and $J'(v) \in H_{0,r}^1(\Omega)^*$, from the Riesz Representation Theorem there exists $z \in H_{0,r}^1(\Omega)$ such that

$$J'(v)w = \int_\Omega \nabla z \cdot \nabla w, \quad \text{for all } w \in H_{0,r}^1(\Omega).$$

Thus, as $z \in H_{0,r}^1(\Omega)$ and $\vartheta^\perp \in H_{0,r}^1(\Omega)^\perp$, we have

$$J'(v)\vartheta^\perp = 0. \tag{3.23}$$

From equations (3.21)–(3.23), for each $w \in H_0^1(\Omega)$, we obtain

$$J'(v)w = J'(v)\vartheta + I'(v)\vartheta^\perp = 0.$$

This allows us to conclude that v is a critical point of the functional J in $H_0^1(\Omega)$ and consequently v is a weak solution for the problem (1.1).

If $v \neq 0$ we are done.

Suppose now that $v \equiv 0$. Considering $v_n \rightarrow 0$, as $n \rightarrow \infty$, we have

$$J'(v_n)v_n = a\|v_n\|^2 + b\|v_n\|^4 - \lambda \int_{\Omega} p^\alpha |v_n|^q - \int_{\Omega} |v_n|^6 = o_n(1). \tag{3.24}$$

But

$$\lambda \int_{\Omega} p^\alpha |v_n|^q \rightarrow 0, \text{ as } n \rightarrow \infty. \tag{3.25}$$

Let $L_1 > 0, L_2 > 0$ be such that

$$a\|v_n\|^2 \rightarrow L_1 \text{ and } b\|v_n\|^4 \rightarrow L_2, \text{ as } n \rightarrow \infty. \tag{3.26}$$

By equations (3.24)–(3.26)

$$\int_{\Omega} |v_n|^6 \rightarrow L_1 + L_2, \text{ as } n \rightarrow \infty. \tag{3.27}$$

But

$$S \left(\int_{\Omega} v_n^6 \right)^{1/3} \leq \int_{\Omega} |\nabla v_n|^2, \tag{3.28}$$

which implies

$$aS \left(\int_{\Omega} v_n^6 \right)^{1/3} \leq a \int_{\Omega} |\nabla v_n|^2 \leq a \int_{\Omega} (|\nabla v_n|^2 + (3/4)p^2 v_n^2) = a\|v_n\|^2, \tag{3.29}$$

and

$$bS^2 \left(\int_{\Omega} v_n^6 \right)^{2/3} \leq b \left[\int_{\Omega} |\nabla v_n|^2 \right]^2 \leq b \left[\int_{\Omega} (|\nabla v_n|^2 + (3/4)p^2 v_n^2) \right]^2 = b\|v_n\|^4. \tag{3.30}$$

Thus, by equations (3.26), (3.27), (3.29), and (3.30)

$$L_1 \geq aS (L_1 + L_2)^{1/3} \text{ and } L_2 \geq bS^2 (L_1 + L_2)^{2/3}. \tag{3.31}$$

On the other hand, $J(v_n) = c + o(1)$. So

$$c = \frac{L_1}{2} + \frac{L_2}{4} - \frac{1}{6}(L_1 + L_2) = \frac{L_1}{3} + \frac{L_2}{12}. \tag{3.32}$$

By equation (3.31), we have

$$(L_1 + L_2)^{1/3} \geq \frac{bs^2 + (b^2s^4 + 4as)^{1/2}}{2}. \tag{3.33}$$

Hence by equations (3.31)–(3.33)

$$\begin{aligned} c &\geq \frac{1}{3}L_1 + \frac{1}{12}L_2 \geq \frac{1}{3}aS(L_1 + L_2)^{1/3} + \frac{1}{12}bS^2 [(L_1 + L_2)^{1/3}]^2 \\ &\geq \frac{1}{4}abS^3 + \frac{1}{24}b^3S^6 + \frac{1}{24}(b^2S^4 + 4aS)^{3/2}, \end{aligned}$$

which is a contradiction with Lemma 2.3. Therefore, we conclude that $v \neq 0$.

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