

# Dispersion in DLA metallicities and deuterium abundances

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**Abstract.** Recent chemical abundance measurements of damped Lyman-alpha absorbers (DLAs) revealed a large intrinsic scatter in their metallicities. We discuss a semi-analytic model that was specifically designed to study this scatter by tracing the chemical evolution of the interstellar matter in small regions of the Universe with different mean density, from over- to underdense regions. It is shown that different histories of structure formation in these regions are reflected in the chemical properties of the proto-galaxies. We also address deuterium abundance measurements, which constitute a complementary probe of the star formation and infall histories.

**Keywords.** ISM: abundances, galaxies: abundances, galaxies: evolution, galaxies: ISM

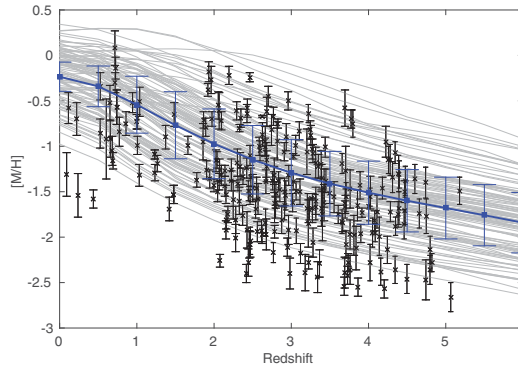
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## 1. Introduction

Numerous studies of damped Ly $\alpha$  absorbers (DLAs) conducted over the past several years have provided extensive samples of high-redshift systems (e.g. Prochaska *et al.* 2005). High resolution observations of nearly 250 of these systems established a statistically significant decrease of DLA metallicity with increasing redshift (e.g. Rafelski *et al.* 2012) and a large intrinsic dispersion of  $\sim 0.5$  dex out to  $z \sim 5$ . We study this dispersion in the context of cosmological structure formation with an efficient semi-analytic model and explore the evolution of DLA metallicities and deuterium abundances.

## 2. Cosmological chemical evolution model

We take  $V_{tot} = 10^6$  (Mpc/h)<sup>3</sup> as the total comoving volume in our calculation and divide it into 1000 smaller *regions*  $\Delta V_i = 10^3$  (Mpc/h)<sup>3</sup>. The volume  $V_{tot}$  is populated with dark matter (DM) halos according to the Sheth-Tormen mass function (MF) at  $z = 0$  accounting for large-scale clustering effects: regions with large-scale overdensities form halos more easily. We use the results of Barkana & Loeb (2004), for the MF bias caused by clustering. We then build a merger tree for each halo using the algorithm in the GALFORM model (Parkinson *et al.* 2008) and follow its evolution backwards in time up to  $z_f = 15$ . We use  $M_{min} = 10^8 M_\odot$  as the minimal halo mass able to form stars. Our next step is to calculate the mean mass fraction in collapsed structures,  $f_{coll}$  and the mean escape velocity in each region  $i$ . Those regions that host a group or a cluster at low redshift have a higher concentration of structures already present at higher redshift, whereas present day voids are relatively empty at early times. This small-scale



**Figure 1.** Metallicity abundance evolution in 100 regions, each with a volume of  $\Delta V_i = 10^3 \text{ Mpc}^3/h^3$  (thin grey lines) and their mean (thick blue line) for  $M_{min} = 10^8 M_\odot/h$ . The dispersion among the realizations and their mean is represented by the blue points with large error bars. Black crosses represent data from Rafelski *et al.* (2012).

inhomogeneity creates the observed dispersion in the different observables. The resulting  $f_{coll}$  and escape velocity for each region serve as inputs for the chemical evolution code developed in Daigne *et al.* (2004; 2006), which follows the exchange of mass between the gas within and outside of collapsed structures, the star formation rate (SFR) at each redshift and the rate of metal production in stars. Our model is calibrated so as to match the cosmic SFR from Behroozi *et al.* (2013). Further details on our model can be found in Dvorkin *et al.* (2015; 2016).

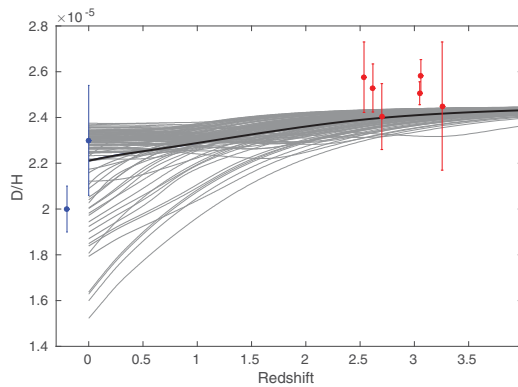
### 3. Results

In Figure 1 we show the evolution of the metal abundance and compare it to observations of DLAs from Rafelski *et al.* (2012). The results shown here include an estimate of the contribution from the mass-metallicity relation using the dispersion in stellar masses within each region (see Dvorkin *et al.* 2015 for details). It can be seen that our model successfully reproduces the observed dispersion with the exception of very metal-poor systems.

Deuterium is created during big bang nucleosynthesis (BBN) and is destroyed in stars. Since deuterium is not created after BBN its abundance constitutes a sensitive probe of the star formation and gas infall history of any given galaxy. The mean evolution of deuterium abundance in our model assuming the primordial value of  $10^5 (D/H)_p = 2.45$  (Coc *et al.* 2015) is shown by the thick black line in Figure 2. The thin grey lines show the evolution in 100 individual regions in our model. It can be seen that even in the most extreme cases the deuterium abundance reduces to about  $\sim 1/3$  of its primordial value, reflecting the overall inefficiency of star formation. We also find that in some cases the evolution of deuterium abundance is non-monotonic due the complex interplay between star formation, which depletes deuterium, and infall of primordial gas which raises its abundance (see Dvorkin *et al.* 2016).

### 4. Discussion

We have shown that the dispersion in the fraction of collapsed structures, escape velocity and SFR between different regions in the Universe contributes to the dispersion in the metallicity-redshift relation of the DLAs. In addition, we calculated the dispersion



**Figure 2.** The evolution of deuterium abundance in each region (thin grey lines) and the mean (thick black line). Measurements of  $z \sim 3$  DLAs from Cooke *et al.* (2014) and Riemer-Sørensen *et al.* (2015), local ISM measurements from Linsky *et al.* (2006) and Prodanović *et al.* (2010) (one of the points at  $z = 0$  is shifted for clarity).

in deuterium abundances expected from different structure formation histories. More observational and theoretical work is needed to improve our understanding of the relation between deuterium and metal abundances, which can provide important constraints on galaxy formation models.

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