

Dark matter distribution and dynamics of dwarf spheroidal galaxies

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Abstract. I review the current status of dynamical modelling of dwarf spheroidal galaxies focusing on estimates of their dark matter content. Starting with the simplest methods using the velocity dispersion profiles I discuss the inherent issues of mass-anisotropy degeneracy and contamination by unbound stars. I then move on to methods of increasing complexity, aiming to break the degeneracy, up to recent applications of the Schwarzschild orbit superposition method. The dynamical modelling is placed in the context of possible scenarios for the formation of dwarf spheroidal galaxies, including the tidal stirring model and mergers of dwarf galaxies. The two scenarios are illustrated with examples from simulations: a comparison between the tidal evolution of dwarfs with cuspy and cored dark matter profiles and the formation of a dwarf spheroidal with prolate rotation.

Keywords. galaxies: dwarf, galaxies: evolution, galaxies: interactions, galaxies: kinematics and dynamics, galaxies: structure, Local Group, dark matter

1. Introduction

Dwarf spheroidal (dSph) galaxies are a class of faint galaxies mostly found in the Local Group, in the vicinity of the Milky Way and Andromeda. They are gas poor, supported mostly by random motions and believed to be strongly dominated by dark matter. Their unique property is that resolved stellar populations are available for study i.e. positions and radial velocities of individual stars can be measured. Over the last two decades dSph galaxies have been the subject of intensive study, with their kinematic samples increasing from few tens to few hundred and now a few thousand of stars with measured velocities for best-studied dwarfs. This allowed us to develop more and more advanced methods of modelling their dark matter content. At the same time, a few scenarios for their dynamical evolution were proposed. In the first part of this contribution I summarize the development of the mass modelling methods, from the simplest to the most advanced, while the second part is devoted to a short description of the formation scenarios.

2. Mass modelling

The simplest observable that can be constructed for dSph galaxies in order to obtain an estimate of their mass content is the line-of-sight velocity dispersion of the stars or its profile as a function of galactocentric radius. This profile can be modelled by solving the Jeans equation which relates the velocity dispersion to the underlying gravitational potential and the velocity anisotropy parameter β . The anisotropy parameter measures the amount of radial versus tangential motion in stellar orbits and can take the values between $\beta = 1$ (radial orbits) and $\beta = -\infty$ (circular orbits) with $\beta = 0$ corresponding to isotropic orbits.

Early studies have demonstrated that dSph galaxies are strongly dominated by dark matter so that their mass-to-light ratios exceed the stellar values by at least one or two orders of magnitude. It has also been shown, e.g. for the best-studied Fornax dSph galaxy (Lokas 2002), that the velocity dispersion profile can be equally well reproduced by different combinations of dark matter density profiles and anisotropy parameters. In general, shallower profiles require more radial stellar orbits and steeper profiles require more tangential orbits to fit a given data set. This is a manifestation of the so-called mass-anisotropy degeneracy inherent in the modelling of any pressure-supported systems, from dSph galaxies to galaxy clusters. For a long time, the simplest and customary way to go around the problem was to assume the stellar orbits to be isotropic in order to infer the mass distribution in dSph galaxies. Such an assumption would obviously lead to biases in the estimated density profiles if the real anisotropy was different. Simulations of galaxy formation and evolution have since demonstrated that the anisotropy indeed tends to depend on radius and departs from zero.

The mass-anisotropy degeneracy can be to some extent broken by adding constraints from a higher velocity moment, the kurtosis, that could be measured from more numerous samples of stellar kinematics which became available. Modelling the kurtosis involves solving the higher order Jeans equation. Since kurtosis is mostly sensitive to the anisotropy, in combination with the modelling of dispersion, both moments can constrain both the mass distribution and anisotropy. However, this still usually requires some assumptions concerning the functional form of orbital anisotropy, for example β to be constant with radius (Lokas *et al.* 2005; Lokas 2009).

The second essential issue in the mass modelling of dSph galaxies is the question of stellar membership. In order to obtain a reliable mass estimate it is important to model a clean sample, i.e. composed only of member stars. If the sample is contaminated by tidally stripped stars or background and foreground stars from the Milky Way that do not trace the potential of the dwarf galaxy, the results will be biased. Such biased estimates will typically lead to the inference of more tangential orbits or more extended mass distributions since unbound stars tend to increase the velocity dispersion at larger galactocentric radii (Klimentowski *et al.* 2007; Lokas *et al.* 2008).

Recently, more advanced methods of mass modelling of dSph galaxies have been developed. In the Schwarzschild orbit superposition method the velocity moments are reconstructed from a combination of stellar orbits. No assumptions on the anisotropy are required and the anisotropy profile follows as a result of the application of the method. For spherical dwarfs the method works very well and is able to reproduce different anisotropy profiles as well as density profiles of simulated galaxies (Kowalczyk *et al.* 2017). For non-spherical objects application of the method with the assumption of spherical symmetry is subject to an inherent bias depending on the line of sight; in particular, the anisotropy is underestimated, and more so for the line of sight along the shortest axis of the stellar component (Kowalczyk *et al.* 2018). A recent application of the Schwarzschild method to the modelling of the Fornax dwarf used around 3300 carefully selected stars from combined catalogues and allowed to constrain both the mass and anisotropy profiles in the dwarf. The mass-to-light ratio was found to increase with radius in agreement with previous estimates and the anisotropy profile was found to be slightly decreasing with radius, but consistent with isotropy at all radii (Kowalczyk *et al.* 2019).

Future improvements in the mass modelling of dSph galaxies should include: increasing the samples of member stars to reduce sampling errors, developing better membership determination methods, using multiple stellar populations, including proper motion measurements for the stars (Strigari *et al.* 2018) and extensions to non-spherical models (Hayashi *et al.* 2018).

3. Formation scenarios

The dynamics of dSph galaxies is best elucidated by studying their formation scenarios. Although these objects could have formed in isolation, the scenarios most studied in recent years involved interactions with other systems. The two main scenarios considered were the tidal stirring due to the proximity of the Milky Way and mergers between dwarfs. In both scenarios the progenitors of dSph galaxies were disky dwarfs. The tidal stirring model (Mayer *et al.* 2001) explains the density-morphology relation observed among the Local Group dwarfs, i.e. the fact that dSph galaxies are situated close to the Milky Way or Andromeda, tend to be stripped of gas, spheroidal and non-rotating while dwarf irregular galaxies are typically found at larger distances, contain gas, are more flattened and rotate. The merger scenario was invoked to explain peculiar structural and kinematic features found in some dwarfs.

3.1. Tidal stirring

In simulations based on the tidal stirring scenario an initially disky dwarf galaxy is placed on a typical eccentric orbit around the Milky Way-like galaxy and evolved for a few gigayears. The effects of tidal stirring involve the mass loss (in both dark matter and stars), the morphological transformation of the stellar component from a disk to a spheroid and the change of rotation into random motions of the stars. The final product of tidal stirring depends strongly on the size of the orbit: it can leave the dwarf's disk almost unaffected for an extended orbit or transform it to a small perfect sphere for a tight orbit (Kazantzidis *et al.* 2011a). The evolution also critically depends on the initial dark matter distribution in the dwarf. The dwarfs with an initial cuspy dark matter profile tend to survive for many gigayears while those with a cored profile dissolve after one or two pericenter passages. A comparison between these two cases is illustrated in Figs. 1 and 2.

An interesting intermediate stage of the evolution from a disk to a spheroid involves the formation of a tidally induced bar in the dwarf (Lokas *et al.* 2014a; Gajda *et al.* 2017). The bar always forms at the first pericenter passage and becomes shorter in time but the spheroidal shape usually survives until the end of evolution which may explain the non-spherical appearance of dSph galaxies. If the gas is included in the progenitor dwarf the tidally induced stellar bars are weaker but the effect is strong only in gas-dominated dwarfs (Gajda *et al.* 2018). On the other hand, if the gas was stripped by ram pressure in the hot halo of the Milky Way then the bars would survive.

Obvious candidates for tidally stirred dwarfs in the vicinity of the Milky Way include the Sagittarius and Carina dwarfs, although most of dSph galaxies in the Local Group show signs of tidal extensions. In the case of the Sagittarius dwarf it has been demonstrated that its present structural and kinematic properties can be reproduced by a model in which the initially disky progenitor has just passed the second pericenter on its orbit around the Milky Way (Lokas *et al.* 2010). In this model the stellar component still preserves its bar-like shape (slightly inclined with respect to the orbit) and shows very little rotation. The recent compilation of data for the Carina dSph galaxy, showing its elongated shape and significant rotation, can also be reproduced by a late-stage tidally stirred disky dwarf (Fabrizio *et al.* 2016).

3.2. Mergers of dwarfs

Some dwarfs show features that cannot be explained by tidal stirring, including e.g. stellar shells in Fornax. In addition, there are distant dSph galaxies in the Local Group that probably did not have enough time to strongly interact with the Milky Way or

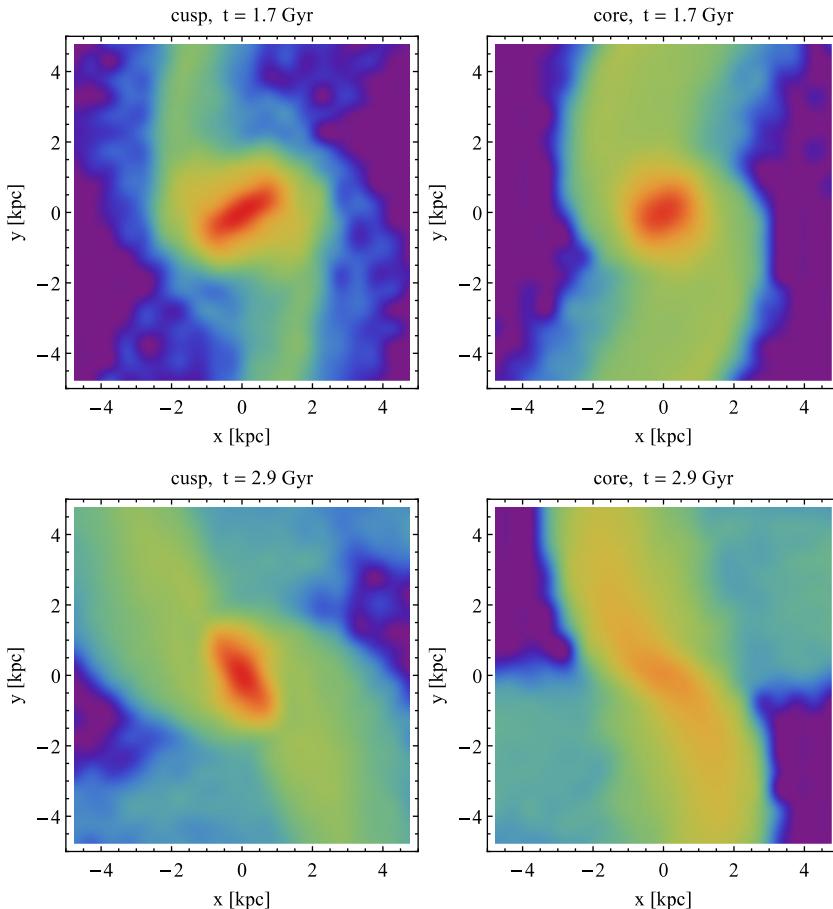


Figure 1. Effect of tidal stirring on two initially disk-like dwarf galaxies evolving around the Milky Way-like galaxy. Both dwarfs were placed on a prograde orbit with apo- to pericenter ratio equal to 100 kpc/20 kpc. The initial structural parameters of the dwarfs differed only by their initial dark matter profile. The dwarf with the initial cuspy (r^{-1}) density profile (left panels) forms a bar at the first pericenter passage and survives for 10 Gyr. The dwarf with a cored ($r^{-0.2}$) profile (right panels) dissolves after the second pericenter passage (Lokas 2016). The stellar components of both dwarfs are shown in projection onto the orbital plane after the first (upper row) and second (lower row) pericenter.

Andromeda (e.g. Cetus and Tucana). Besides, in tidally stirred dwarfs the remnant rotation is always around the short axis of the stellar component while some dwarfs show rotation around the long axis (prolate rotation), e.g. Andromeda II (Ho *et al.* 2012) and Phoenix (Kacharov *et al.* 2017). Such features point toward a different scenario for the formation of at least some dSph galaxies, such as mergers between initially disk-like dwarfs (Kazantzidis *et al.* 2011b). This possibility is supported by constrained simulations of the Local Group, aiming to reproduce the properties of both big galaxies, the Milky Way and Andromeda, as well as their satellite population. In particular, it has been demonstrated that in such environments a significant fraction of subhaloes experienced a strong interaction with another subhalo (Klimentowski *et al.* 2010). Since the typical velocities of single subhaloes around the Milky Way are rather too large for such interactions to be significant, mergers usually occur between subhaloes that were accreted in pairs, whose relative velocities are low.

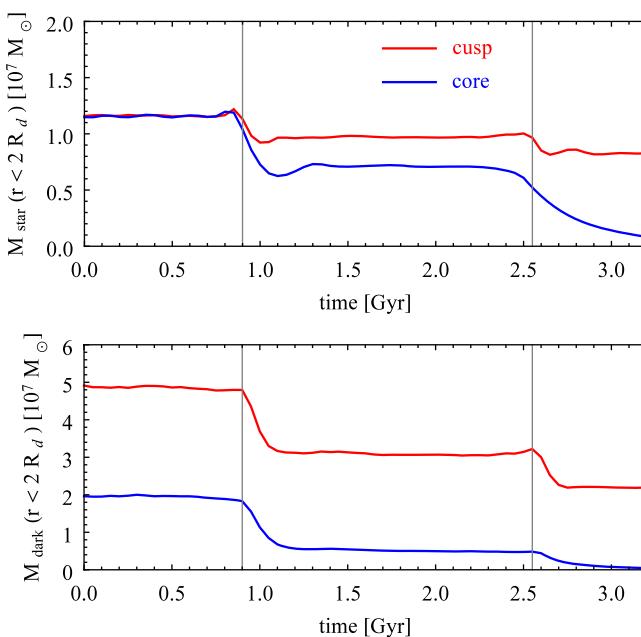


Figure 2. Evolution of the mass contained within two disk scale-lengths $2R_d = 0.82$ kpc for stars (upper panel) and dark matter (lower panel) for the dwarfs with initially cuspy (red upper line) and cored (blue lower line) dark matter profile. The vertical gray lines indicate pericenter passages on an orbit with apo- to pericenter of 100 kpc/20 kpc. For the cored dwarf both masses approach zero after the second pericenter passage signifying the disruption of the dwarf.

The merger scenario has been proposed as a way to explain the origin of prolate rotation in Andromeda II (Lokas *et al.* 2014b). In this scenario, the progenitors are two disk-like dwarfs of equal mass and similar structural parameters approaching each other on a radial orbit. Their disks are initially inclined at the right angle but both have significant components of their angular momenta aligned with the merger axis. These parts of angular momenta are preserved after the merger in the form of prolate rotation. The remnant dwarf rotates around its long axis while the rotation around the short axis is close to zero. Three different stages of such a merger are illustrated in Fig. 3. The model reproduces well the non-spherical shape of the stellar component and if the gas is included (Fouquet *et al.* 2017) also the presence and distribution of different stellar populations, with the younger one formed during the merger. In order to lose the remaining gas, the dwarf galaxy must have additionally experienced some interaction with its host, Andromeda, whose hot gaseous halo would strip the dwarf's gas via ram pressure. It is worth emphasizing that tidal stirring is unable to induce significant stable rotation around the longest axis. Even if it occurs, it is always much smaller than the velocity dispersion, while in Andromeda II these two quantities are comparable. We have also shown that prolate rotation can be obtained from less idealized initial conditions, i.e. from dwarfs merging on non-radial orbits and with different inclinations of the disks (Ebrouvá & Lokas 2015).

The scenario has been placed in the cosmological context although for more massive galaxies due to resolution limitations. Almost 60 galaxies with prolate rotation have been identified in the Illustris simulation and most of them were confirmed to be the result of a major merger (Ebrouvá & Lokas 2017). In most of them a very tight correlation was found between the time of merger and the time of transition to prolate rotation and prolate shape. In massive galaxies prolate rotation tends to occur only in gas-poor objects.

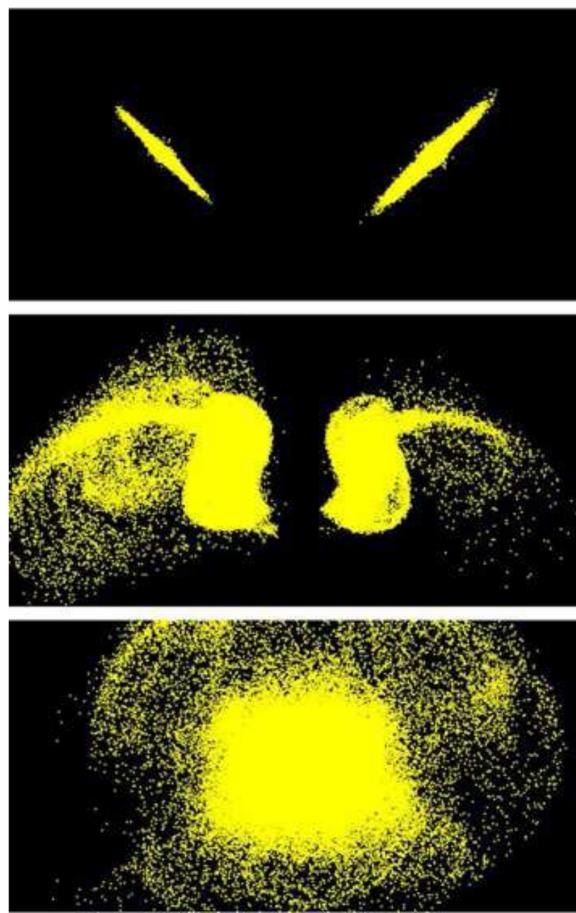


Figure 3. Formation of a dSph galaxy similar to Andromeda II by a merger of two initially disky dwarfs (Lokas *et al.* 2014b). The three panels show the stellar components in three stages of the merger: before, with the two galaxies approaching (upper panel), after the first encounter (middle panel) and at the end of evolution, after the merger (lower panel). The image in each panel has the size of 40 kpc \times 20 kpc.

Although originally believed to be relatively simple, near-spherical, old-population, pressure-supported systems, dSph galaxies of the Local Group only now start to reveal their complexity. Andromeda II has been recently shown to host at least two stellar populations with different kinematics. While the older population shows prolate rotation, the younger one rotates around the minor axis (del Pino *et al.* 2017b). Even larger degree of complexity has been discovered in the well-studied Fornax dwarf. The multiple stellar populations of Fornax display different distributions and kinematics with the youngest stars centrally concentrated and showing a distinct disky shape. Tracing the kinematics of different stellar populations as a function of age leads to the conclusion that also Fornax probably experienced a merger about 8 Gyr ago (del Pino *et al.* 2017a).

4. Conclusions

DSpH galaxies of the Local Group remain an interesting subject of investigation. Determining the dark matter distribution in dSph galaxies remains challenging but promising pathways for improvement exist. In the era of large surveys we can expect the kinematic samples for the dwarfs to increase dramatically allowing for the modelling methods to reach the next level of refinement. These should include in the near future

non-spherical orbit-based modelling and attempts to constrain the inner slope of the dark matter distribution which remained elusive so far. The measurements of proper motions for a significant number of bright stars in dSphs should put additional constraints on the models, both in terms of dark matter distribution and anisotropy of stellar orbits. These studies could also benefit from the use of multiple stellar populations that at least in principle should be able to constrain dark matter distribution at different scales.

Many dwarfs of the Local Group show signs of interactions that shaped them. The two dominant scenarios of such interactions are the tidal stirring in the vicinity of the Milky Way and Andromeda, and mergers between dwarfs. For example, tidal stirring can account for the non-spherical shape and remnant rotation in Carina and a merger can explain the prolate rotation in Andromeda II. More advanced models created in close relation to recent observational findings are needed to reproduce the properties of these quite complicated systems in more detail.

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References

- del Pino, A., Aparicio, A., Hidalgo, S. L., & Łokas, E. L. 2017a, *MNRAS*, 465, 3708
 del Pino, A., Łokas, E. L., Hidalgo, S. L., & Fouquet, S. 2017b, *MNRAS*, 469, 4999
 Ebrová, I. & Łokas, E. L. 2015, *ApJ*, 813, 10
 Ebrová, I. & Łokas, E. L. 2017, *ApJ*, 850, 144
 Fabrizio, M., Bono, G., Nonino, M., et al. 2016, *ApJ*, 830, 126
 Fouquet, S., Łokas, E. L., del Pino, A., & Ebrová, I. 2017, *MNRAS*, 464, 2717
 Gajda, G., Łokas, E. L., & Athanassoula, E. 2017, *ApJ*, 842, 56
 Gajda, G., Łokas, E. L., & Athanassoula, E. 2018, *ApJ*, 868, 100
 Hayashi, K., Fabrizio, M., Łokas, E. L., Bono, G., Monelli, M., Dall’Ora, M., & Stetson, P. B., 2018, *MNRAS*, 481, 250
 Ho, N., Geha, M., Muñoz, R. R., et al. 2012, *ApJ*, 758, 124
 Kacharov, N., Battaglia, G., Rejkuba, M., et al. 2017, *MNRAS*, 466, 2006
 Kazantzidis, S., Łokas, E. L., Callegari, S., Mayer, L., & Moustakas, L. A. 2011, *ApJ*, 726, 98
 Kazantzidis, S., Łokas, E. L., Mayer, L., Knebe, A., & Klimentowski, J. 2011, *ApJ*, 740, L24
 Klimentowski, J., Łokas, E. L., Kazantzidis, S., Prada, F., Mayer, L., & Mamon, G. A. 2007, *MNRAS*, 378, 353
 Klimentowski, J., Łokas, E. L., Knebe, A., Gottlöber, S., Martinez-Vaquero, L. A., Yepes, G., & Hoffman, Y. 2010, *MNRAS*, 402, 1899
 Kowalczyk, K., Łokas, E. L., & Valluri, M. 2017, *MNRAS*, 470, 3959
 Kowalczyk, K., Łokas, E. L., & Valluri, M. 2018, *MNRAS*, 476, 2918
 Kowalczyk, K., del Pino, A., Łokas, E. L., & Valluri, M. 2019, *MNRAS*, 482, 5241
 Łokas, E. L. 2002, *MNRAS*, 333, 697
 Łokas, E. L. 2009, *MNRAS*, 394, L102
 Łokas, E. L. 2016, *Galaxies*, 4, 74
 Łokas, E. L., Mamon, G. A., & Prada, F. 2005, *MNRAS*, 363, 918
 Łokas, E. L., Klimentowski, J., Kazantzidis, S., & Mayer, L. 2008, *MNRAS*, 390, 625
 Łokas, E. L., Kazantzidis, S., Majewski, S. R., Law, D. R., Mayer, L., & Frinchaboy, P. M. 2010, *ApJ*, 725, 1516
 Łokas, E. L., Athanassoula, E., Debattista, V. P., Valluri, M., del Pino, A., Semczuk, M., Gajda, G., & Kowalczyk, K. 2014a, *MNRAS*, 445, 1339
 Łokas, E. L., Ebrová, I., del Pino, A., & Semczuk, M. 2014b, *MNRAS*, 445, L6
 Mayer, L., Governato, F., Colpi, M., Moore, B., Quinn, T., Wadsley, J., Stadel, J., & Lake, G. 2001, *ApJ*, 559, 754
 Strigari, L. E., Frenk, C. S., & White, S. D. M. 2018, *ApJ*, 860, 56