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The massive Proca field

The massive vector field is the model which describes massive vector bosons, such as the W and Z particles of the electro-weak theory.

22.1 Action and field equations

The action for the Proca field is

$$S = \int (dx) \left\{ \frac{1}{4} F^{\mu\nu} F_{\mu\nu} + \frac{1}{2} m^2 A^\mu A_\mu - J^\mu A_\mu \right\}, \quad (22.1)$$

where

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu. \quad (22.2)$$

The variation of the action gives

$$\delta A = \int (dx) \left[-\partial^\nu F_{\mu\nu} + m^2 A_\mu - J_\mu \right] \delta A_\mu + \int d\sigma^\mu F_{\mu\nu} \delta A^\nu. \quad (22.3)$$

This yields the field equation

$$-\partial^\nu F_{\mu\nu} + m^2 A_\mu = J_\mu, \quad (22.4)$$

also writable as

$$-\square A_\mu - \partial_\mu (\partial^\nu A_\nu) + m^2 A_\mu = J_\mu, \quad (22.5)$$

and associated continuity conditions identical to those of the Maxwell field. The conjugate momentum ($d\sigma^\mu = d\sigma^0$) is

$$\Pi_i = F_{0i}. \quad (22.6)$$

If the surface σ is taken to separate two regions of space rather than time, one has the continuity conditions for the Proca field in a vacuum:

$$\begin{aligned}\Delta F_{i0} &= 0 \\ \Delta F_{ij} &= 0,\end{aligned}\tag{22.7}$$

and we have assumed that δA_μ is a continuous function. Taking the $n + 1$ divergence of eqn. (22.4), we obtain

$$\partial^\mu A_\mu = 0.\tag{22.8}$$

Here we have used the anti-symmetry of $F_{\mu\nu}$ and the assumption that the source is conserved, $\partial^\mu J_\nu = 0$. Thus the field equations, in the form of eqn. (22.5), become

$$(-\square + m^2)A_\mu = J_\mu\tag{22.9}$$

$$\partial^\mu A_\mu = 0.\tag{22.10}$$

In contrast to the electromagnetic field, this has both transverse and longitudinal components.