

Negative Differential Conductivity in AlGaN/GaN HEMTs: Real Space Charge Transfer from 2D to 3D GaN States?

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ABSTRACT

We report on non-thermal negative differential conductivity (NDC) in AlGaN/GaN HEMTs grown on sapphire substrates by low-pressure MOCVD. The sheet electron density was on the order of few times 10^{12}cm^{-2} and the Hall mobility was $1,000\text{cm}^2/\text{V}\cdot\text{s}$. The HEMTs had threshold voltage close to zero and could operate at high positive gate bias up to 3 to 3.5 Volts, with a very low gate leakage current. NDC was observed at the gate bias larger than 1.5V and at the drain biases between approximately $0.5V_g$ and V_g . We excluded the possibility of self-heating as the cause, since the NDC occurs at relatively small power levels where self-heating effects are negligible.

An explanation we provided for the NDC effect is the new mechanism of real space charge transfer from 2D to 3D GaN states, which leads to a decrease in the channel mobility at large 2D electron gas densities. The observed low leakage can be explained by an enhanced molar fraction of aluminum at the heterointerface that results in a larger conduction band discontinuity. Our model that accounts for the piezoelectric and pyroelectric effects is consistent with the observed NDC effect. The Hall mobility dependence on the gate bias and sheet carrier concentration [1] is consistent with the real space transfer mechanism.

This NDC effect in GaN/AlGaN HEMTs may find applications in high-performance digital circuits at elevated temperatures.

INTRODUCTION

Recently, there has been much progress in the research and development of AlGaN/GaN High Electron Mobility Transistors (HEMTs). However, the physics of these novel devices has not yet been fully understood. In this paper, we report on a new mechanism of the negative differential mobility in AlGaN/GaN HEMTs.

Negative differential conductance (NDC) in HEMTs was previously observed [2-3]. However, in these publications, the NDC was attributed to either gate leakage current [2] or self-heating effect [3]. In [2], it was claimed that the two dimensional electron gas was heated in the channel. This heating caused the "real-space transfer" [4] between the GaAs conducting layer and the AlGaAs barrier layer. Self-heating might also lead to NDC [3] since with the maximum source-drain current of $1\text{A}/\text{mm}$, the GaN devices grown on sapphire biased at 5V could have a temperature rise of 125K, which would significantly reduce the mobility and the maximum charge density of the channel.

In our samples, these mechanisms do not play an important role for the following reasons. Self-heating is not important since the power dissipation in our samples is only on the order of $0.5\text{W}/\text{mm}$. For an estimated impedance of $25\text{W}/\text{mm}$ [3], this power dissipation should lead to the temperature rise of 12K, which is not enough to cause NDC.

As shown below, piezoelectric and pyroelectric effects in AlGaN/GaN heterostructures hinder the real space transfer into the AlGaN layer. However, as we discuss below, real space transfer into the 3D states in GaN is a possible mechanism explaining the observed NDC.

EXPERIMENTS

The $\text{Al}_{0.2}\text{Ga}_{0.8}\text{N}/\text{GaN}$ HEMTs were grown on sapphire substrates by low-pressure metal organic chemical vapor deposition (MOCVD). The epilayer design of the samples was described in [5]. The measured 2DEG sheet density was on the order of 10^{12}cm^{-2} and the Hall mobility was $1,000\text{ cm}^2/\text{V}\cdot\text{s}$.

Figure 1(a) shows the current-voltage characteristics of the $45\text{ }\mu\text{m}$ wide AlGaN-GaN HEMT. The source-drain spacing L_{ds} was $5\text{ }\mu\text{m}$ and the gate length L_{g} was $1.7\text{ }\mu\text{m}$. The applied gate bias was as high as 3.0 V . The negative differential conductivity (NDC) was observed at the gate biases $V_{\text{g}} > 1.5\text{ V}$. The drain current increased with V_{ds} to the peak value, I_{peak} , determined by the gate bias, and then decreased until saturation current (I_{sat}) has been reached. For HEMTs with the drain-to-gate spacing of $1\text{ }\mu\text{m}$ and $2\text{ }\mu\text{m}$ the peak drain current was achieved at the source-drain voltages of approximately $V_{\text{peak}} = 1.5\text{ V}$ and $V_{\text{peak}} = 2.5\text{ V}$, respectively. The NDC was observed at the source-drain voltages $V_{\text{peak}} < V_{\text{ds}} < V_{\text{valley}}$, where V_{valley} was the drain bias when the NDC effect stopped. The drain current reached its saturation beyond V_{valley} . In Figure 1(b), we show the gate bias dependence of V_{peak} , V_{valley} and the peak-valley current ratio $R = I_{\text{peak}}/I_{\text{valley}}$. This figure demonstrates that V_{valley} increased with the gate bias with a slope of about 1, while the increase rate of V_{peak} was much smaller. On the other hand, R also increased almost linearly with V_{g} .

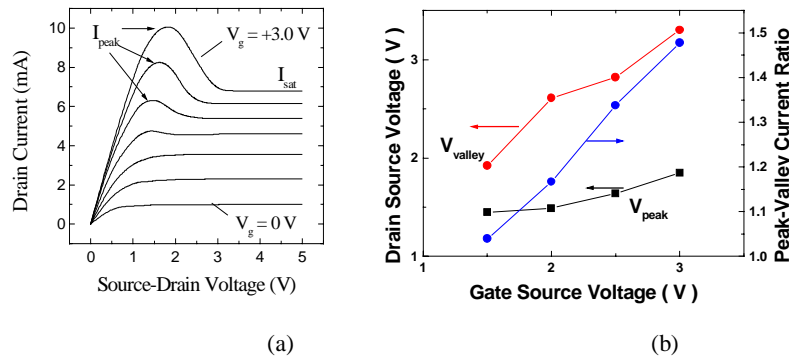


Figure 1. Drain current-voltage characteristics of the $45\text{ }\mu\text{m}$ wide AlGaN-GaN HEMT. (a) Drain current dependence on V_{d} , V_{g} varies from 0 to 3V with the step of 0.5V. (b) The V_{peak} , V_{valley} and $I_{\text{peak}}/I_{\text{sat}}$ as functions of the gate bias.

We measured NDC in the HEMTs with the source-drain spacing from $5\text{ }\mu\text{m}$ to $15\text{ }\mu\text{m}$ and the offset gate length from $2\text{ }\mu\text{m}$ to $12\text{ }\mu\text{m}$. The source-to-gate and gate-to-drain distances were approximately the same and equal to $2\text{ }\mu\text{m}$ and $1\text{ }\mu\text{m}$, respectively. Figure 2 demonstrates the normalized current drop as a function of the gate length for $95\text{ }\mu\text{m}$ wide HEMTs ($\Delta I_{\text{neg}} = I_{\text{peak}} - I_{\text{sat}}$). The obtained results clearly showed strong dependence of NDC on L_{g} . At the gate bias of $+3.0\text{ V}$, the increase in the gate length for about 6 times (source-drain spacing increases 3 times) reduces NDC by a factor of larger than 30. The saturation current for the same gate lengths decreases

linearly only by one third, from 12 mA to 8 mA. The ΔI_{neg} dependence on the gate length can be described by expression $\Delta I_{\text{neg}} \sim L_g^{-\alpha}$, where α is from 1.5 to 2.

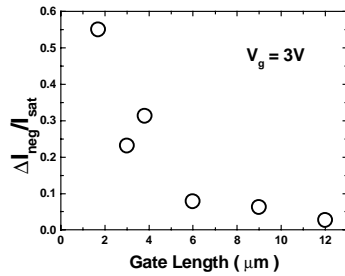


Figure 2. Gate length dependence of the current drop in NDC (normalized to the I_{sat})

DISCUSSION

The gate current versus drain voltage at different gate biases for the device in Figure 1 was measured. (See also [6]). Even for the zero source-drain voltage, the gate leakage current at high positive gate bias ($V_g = +3.0$ V) was less than 0.5 mA, whereas the drain current reduction in the NDC region for the same gate bias was more than 3mA (see Figure 1(a)). At the source-drain voltage of 1.5 V and higher (NDC region) the gate current dropped below $1\mu\text{A}$. These results indicated that the gate current of the device were negligible compared to ΔI_{neg} .

As we mentioned before, the device had a maximum drain current on the order of 0.1~0.2A/mm, which would only raise the temperature by 10~20K at a thermal impedance of 25 $\text{K}\cdot\text{mm}/\text{W}$ [3]. Also, the self-heating NDC led to a monotonous decrease of the drain current in the saturation regime, whereas the NDC reported here took place in a limited drain voltage range (see Figure 1(a)). Therefore, we concluded that both the gate current and self-heating were not the mechanism leading to the NDC in our samples.

The explanation we provided for the NDC was the mobility and sheet carrier density decrease due to the real space transfer from the 2D to 3D GaN states. First, we compared the mobility of the two-dimensional and the three-dimensional electron gas. The results are shown in Figure 3. Even at room temperature, the mobility of 2D-electron gas is around 50% larger than the mobility of 3D electrons. This difference is much larger at cryogenic temperatures.

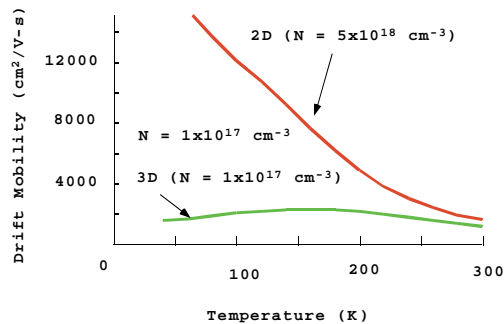


Figure 3. Calculated 2D and 3D electron mobility in GaN at different temperatures

In Figure 4, we compared the band diagram of HEMTs in different material systems – GaN-based and GaAs-based. Due to the piezoelectric and pyroelectric effects, the conduction band profile in the AlGaIn layer does not have a minimum that is needed for the real space transfer from GaN into AlGaIn. Also, in contrast to AlGaAs/GaAs system, the energy difference between the Fermi level in the 2D gas and the bottom of the conduction band in the GaN is relatively small due to the residual doping in GaN. This small difference made the electron transition into the GaN bulk more important.

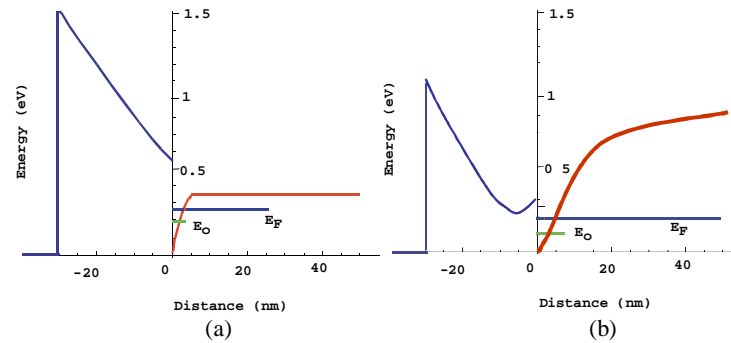


Figure 4. Comparison between the band diagram of GaN based HEMTs and GaAs based HEMTs. (a) Calculated band diagram of GaN HEMTs at the doping level of 10^{17}cm^{-3} (b) Qualitative band diagram of GaAs HEMTs

Hence, we believe that this transfer of 2D electrons into 3D GaN states reduces the channel mobility and the 2D-electron sheet concentration and leads to the current drop in the NDC region. The Monte-Carlo simulation results for AlGaAs/GaAs HEMTs reported in [7] also pointed out to the possibility of such a mechanism (see Figure 5). The dash-dotted line in Figure 5 shows the electron concentration in the quantum well and the full lines show the bulk GaAs electron concentration. This figure demonstrates that the electrons inside the quantum well experience the transition into the GaAs buffer along the channel. However, due to the low barrier height of the GaAs transistors (Figure 4(b)), the gate current becomes dominant at higher gate biases, which makes it difficult to observe the NDC related to such a transfer in GaAs devices at room temperature.

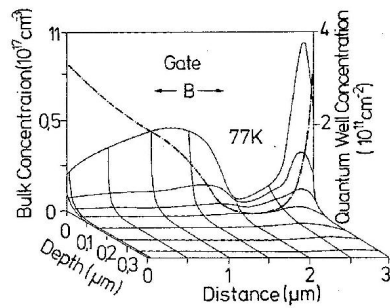


Figure 5. Monte Carlo simulation results on electron concentration in GaAs based HEMTs. Dashed dotted line, concentrations in a quantum well; full line, concentrations in bulk GaAs. After [7]

In order to illustrate this new mechanism of NDC, we calculated the current-voltage characteristics of our devices for different values of the electron mobility using the AIM-Spice [8], which helped to understand the device operation. In Figure 6(a), a mobility of $750 \text{ cm}^2/\text{V}\cdot\text{s}$ was used and reasonable agreement was reached in the low drain bias region. The mobility dropped to $400 \text{ cm}^2/\text{V}\cdot\text{s}$ in Figure 6(b), which give reasonable fit in the saturation region. The maximum sheet charge density for these two regions was also different. (See Table 1) These simulation results confirm that a mobility drop at the higher gate bias might be the cause of the NDC effect, as illustrated in Figure 3.

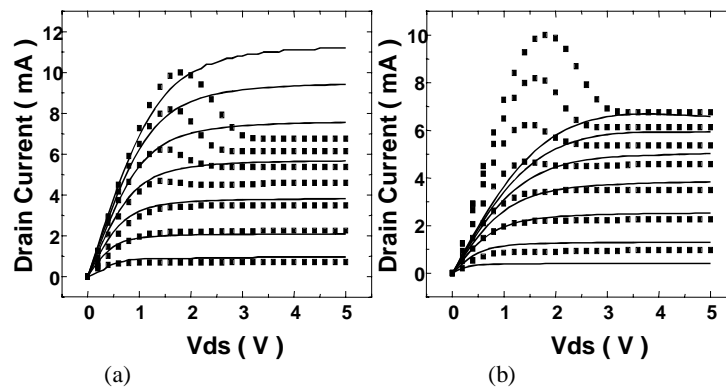


Figure 6. AIM-Spice model of the above device. The model parameters are in Table 1. The symbols are the experimental results and the lines are from simulation.

CONCLUSIONS

In this paper, we report the non-thermal negative differential conductance in AlGaIn/GaN HEMTs. With the help of device characterization and modeling, we excluded the possibility of self-heating and real space transfer into the AlGaIn layer current as the causes of the NDC. We linked

the NDC to the mobility reduction caused by the transfer from 2D to 3D states in the channel. This mechanism was supported by the device simulation and band diagram calculations.

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Table 1. AIM-Spice parameters for Figure 6.

Parameters	2D channel	3D channel	Unit
W_g (Gate width)	45	45	μm
L_g (Gate length)	1.7	1.7	μm
D_i (Distance to buffer layer charge)	30	30	nm
λ (Output conductance parameter)	0	0	V^{-1}
μ (Low-field mobility)	0.075	0.04	$\text{m}^2 \text{V}^{-1} \text{s}^{-1}$
n_{max} (Sheet charge density)	$1 \cdot 10^{17}$	$4 \cdot 10^{16}$	m^{-2}
Φ_b (Heterojunction barrier height)	1.1	1.1	V
R_s (Drain series resistance)	30	30	Ω
R_d (Source series resistance)	30	30	Ω
v_s (Saturation velocity)	40000	40000	m/s
M (Knee voltage parameter)	2	2	-
V_T (Zero-bias threshold voltage)	-0.6	-0.6	V

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