

The summation of a slowly convergent series

By S. PATERSON.

The series

$$S_{2r}(x) \equiv \sum_{n=1}^{\infty} n^{2r} e^{-n^2 x} \quad (1)$$

in which r is zero or an integer is rapidly convergent if x is large but may be very slowly convergent if x is small. The object of this note is to derive an alternative series for $S_{2r}(x)$ which is rapidly convergent for small values of x .

If we let

$$\Phi_{-r}(z) \equiv \frac{2}{\sqrt{\pi}} \int_z^{\infty} \int_z^{\infty} \dots \int_z^{\infty} e^{-z} (dz)^{r+1}, \quad r \geq -1,$$

then it can readily be established¹ by induction that

$$2r \Phi_{-r}(z) = -2z \Phi_{-(r-1)}(z) + \Phi_{-(r-2)}(z) \quad (2)$$

from which it follows that

$$(2\sqrt{x})^{r-2} \Phi_{-(r-2)}(n\pi/\sqrt{x}) = \frac{d}{dx} \left\{ (2\sqrt{x})^r \Phi_{-r}(n\pi/\sqrt{x}) \right\}. \quad (3)$$

Now as a special case of Poisson's identity we have

$$S_0(x) \equiv \sum_{n=1}^{\infty} e^{-n^2 x} = -\frac{1}{2} + \frac{1}{2}\sqrt{\pi/x} + \sqrt{\pi/x} \sum_{n=1}^{\infty} e^{-n^2 \pi^2/x}$$

which may be written in the form

$$S_0(x) = \frac{1}{2}(\sqrt{\pi/x} - 1) + \frac{1}{2} \pi x^{-\frac{1}{2}} \sum_{n=1}^{\infty} \Phi_1(n\pi/\sqrt{x}). \quad (4)$$

Integrating this equation with respect to x from x to ∞ leads to the relation

$$S_{-2}(x) = \frac{\pi^2}{6} - \sqrt{\pi x} + \frac{1}{2}x - 2\pi\sqrt{x} \sum_{n=1}^{\infty} \Phi_{-1}(n\pi/\sqrt{x}) \quad (5)$$

which in turn leads to

$$S_{-4}(x) = \frac{\pi^4}{90} - \frac{\pi^2 x}{6} + \frac{2(\pi x^3)^{\frac{1}{2}}}{3} - \frac{x^2}{4} + \pi(2\sqrt{x})^3 \sum_{n=1}^{\infty} \Phi_{-3}(n\pi/\sqrt{x}). \quad (6)$$

¹ Hartree, D. R., *Mem. Proc. Manchester Lit. Phil. Soc.*, **80**, 85, (1936).

In the general case we obtain the expression

$$\begin{aligned}
 S_{-2r}(x) = & 2^{2r-1} \pi^{2r} \left\{ \frac{B_r}{(2r)!} - \frac{B_{r-1} x}{4\pi^2(2r-2)!} \right. \\
 & + \dots + (-1)^{r-1} \frac{B_1 x^{r-1}}{(4\pi^2)^{r-1} 2! (r-1)!} \left. \right\} + (-1)^r \sqrt{\pi} \frac{2^{2r-1} r! x^r}{(2r)!} \\
 & + (-1)^{r-1} \frac{x^r}{2(r)!} + (-1)^{r+1} \pi (2\sqrt{x})^{2r-1} \sum_{n=1}^{\infty} \Phi_{-(2r-1)}(n\pi/\sqrt{x}) \quad (7)
 \end{aligned}$$

in which B_r denotes the r -th Bernoulli number.

The magnitude of the contribution $\pi(2\sqrt{x})^{2r-1} \sum \Phi_{-(2r-1)}$ to the final result can be estimated by expanding $\Phi_{-(2r-1)}$ in descending powers of $\pi x^{-\frac{1}{2}}$; we then find that $\pi(2\sqrt{x})^{2r-1} \sum \Phi_{-(2r-1)}$ is of the order of $(x/\pi)^{2r-\frac{1}{2}} e^{-\pi^2/x}$ which is always negligible if $0 < x \leq 1$.

On the other hand if we differentiate (4) repeatedly and make use of equation (3) in the form

$$(4x)^{-r-\frac{1}{2}} \Phi_{2r+1}(n\pi/\sqrt{x}) = \frac{d^r}{dx^r} \left\{ (4x)^{-\frac{1}{2}} \Phi_1(n\pi/\sqrt{x}) \right\}$$

we find that

$$\begin{aligned}
 S_{2r} = & \frac{(2r)! \sqrt{\pi}}{r! (2\sqrt{x})^{2r+1}} + \frac{(-1)^r 2\sqrt{\pi}}{(2\sqrt{x})^{2r+1}} \sum_{n=1}^{\infty} \left[\frac{d^{2r}}{dz^{2r}} \right]_{z=n\pi/\sqrt{x}} \\
 = & \frac{\sqrt{\pi}}{2^r x^{r+\frac{1}{2}}} \left\{ \frac{(2r)!}{r! 2^{r+1}} + (-1)^r \sum_{n=1}^{\infty} e^{-n^2\pi^2/2x} D_{2r}(n\pi\sqrt{2/x}) \right\} \quad (8)
 \end{aligned}$$

where D denotes Weber's parabolic cylinder function.

Since the first term in the asymptotic expansion of $D_{2r}(z)$ is $e^{-\frac{1}{2}z^2} \cdot z^{2r}$ the ratio of the series in (8) to the other term is of order $2(r!) (4\pi^2/x)^r e^{-\pi^2/x}/(2r)!$, and this is of the order 10^{-4} , or smaller, if

$$r \leq \pi^2 (1/x - 1) / \log(\pi^2/x).$$

Thus we conclude that

$$S_{2r} = \frac{\sqrt{\pi} (2r)!}{r! (4x)^{r+\frac{1}{2}}}, r \geq 1 \quad (9)$$

to better than 1 in 10,000 provided that

$x \leq$	0.7	0.62	0.52	0.45	0.38	0.32	0.28	0.25	0.22	0.20
if $r =$	1	2	3	4	5	6	7	8	9	10.

Since we have neglected the factor $2^{2r+1} (r!)/(2r)!$ which decreases as r increases, the range is actually somewhat wider for the larger values of r .

By writing y for π/\sqrt{x} in the above relations, we obtain a corresponding set of expressions for $\sum_{n=1}^{\infty} \Phi_{2r+1}(ny)$ for $y \geq 10$. For example from (8) we have

$$\sum_{n=1}^{\infty} \Phi_{2r+1}(ny) \doteq (-1)^{r+1} (2r)!/\pi^{1/2} r!$$

which is independent of y (≥ 10), and for $r = 1$ gives

$$\sum_{n=1}^{\infty} e^{-4n^2 z^2} D_2(nz) \doteq \frac{1}{2}, \text{ if } z > 14.$$

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A Generalisation of Dirichlet's Multiple Integral

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The purpose of this note is to generalise the Dirichlet-Liouville formula which expresses a certain type of multiple integral in terms of a single integral.¹ In our formula the multiple integral will involve several arbitrary functions instead of only one, and it will be expressed as a product of single integrals.

Let n be a positive integer. Let $f_1(t), f_2(t), \dots, f_n(t)$ be Lebesgue measurable functions when $0 \leq t \leq 1$. A finite sequence of n real numbers m_1, m_2, \dots, m_n is given. We write $m_{n+1} = 0$ and

$$M_r = m_1 + m_2 + \dots + m_r$$

$$X_r = x_1 + x_2 + \dots + x_r$$

¹ See, for example, G. F. Meyer, *Vorlesungen über die Theorie der bestimmten Integrale* (Leipzig, 1871), 566 *et seq.*; or E. T. Whittaker and G. N. Watson, *Modern Analysis* (4th edn., Cambridge, 1935), section 12.5; or H. Jeffreys and B. S. Jeffreys, *Methods of Mathematical Physics* (Cambridge, 1946), section 15.08; or L. J. Mordell, "Dirichlet's integrals," *Edin. Math. Notes*, No. 34 (1944), 15-17.