

High-Q photonic chip-based temporal phase plates for electron microscopy

Armin Feist¹, Arslan Sajid Raja², Jan-Wilke Henke¹, Junqiu Liu², Germaine Arend¹, Guan hao Huang², Fee Jasmin Kappert¹, Rui Ning Wang², Jiahe Pan², Ofer Kfir¹, Tobias Kippenberg² and Claus Ropers³

¹4th Physical Institute – Solids and Nanostructures, University of Göttingen, Göttingen, Germany, United States, ²Institute of Physics, Swiss Federal Institute of Technology Lausanne (EPFL), Lausanne, Switzerland, United States, ³4th Physical Institute – Solids and Nanostructures, University of Göttingen, Göttingen, Germany, Goettingen, Niedersachsen, Germany

The active optical shaping of free-electron beams enables a broad range of applications, from efficient free-space acceleration [1] and attosecond bunching of electrons [2-4] to the implementation of laser-driven phase plates [5,6] and beam splitters [7].

Recent experiments on inelastic electron-light scattering (IELS) [8,9] with dielectric microcavities such as microspheres [10] or photonic crystal cavities [11] have demonstrated the potential of using resonant field enhancement for increasing the interaction between electrons and photons. However, bringing IELS to state-of-the-art continuous-beam electron microscopes is challenging, due to the optical transition probabilities normally being far below unity, thus requiring the use of femtosecond high-intensity laser pulses.

In this work, we demonstrate IELS on a CW-pumped Si₃N₄ microresonator with a Q-factor of >10⁶, achieving an unprecedented high coupling to a continuous electron beam. The interaction strength between the electrons and the evanescent cavity field is spatially mapped and we discuss the potential application as a versatile temporal phase plate for electrons.

In an ultrafast transmission electron microscope (UTEM) [12], a continuous electron beam passes by a CW pumped Si₃N₄ microresonator (see Fig. 1.a)), which is installed in a custom-designed TEM sample holder. The air-cladded microresonator chips (Fig. 1.b)) are fabricated using the photonic Damascene process [13] and exhibit a linewidth of ~70 MHz (Fig. 1.c)), as well as a free spectral range of ~1 THz for the quasi-TM fundamental mode. The electron beam interacts with the whispering gallery mode, confined in the fiber-coupled photonic chip.

When the CW laser is tuned to a resonance of the cavity, the initially narrow energy distribution is significantly broadened up to about 50 eV, with a shape typical for multi-order absorption and emission of photons (Fig. 1.a,d)). We retrieve a spatial map of the electron-light coupling parameter g , which exhibits a pronounced modulation along the structure (Fig. 1.e)), due to Ramsey-type interference [14] between the two crossings of the electrons and the resonator ring.

In conclusion, we demonstrate strong inelastic electron scattering off a CW optical field, enabled by high-Q microresonators. Hereby, new opportunities arise in electron microscopy, from optical beam structuring with state-of-the-art integrated photonics, to the exploration of free-electron cavity quantum optics [15].

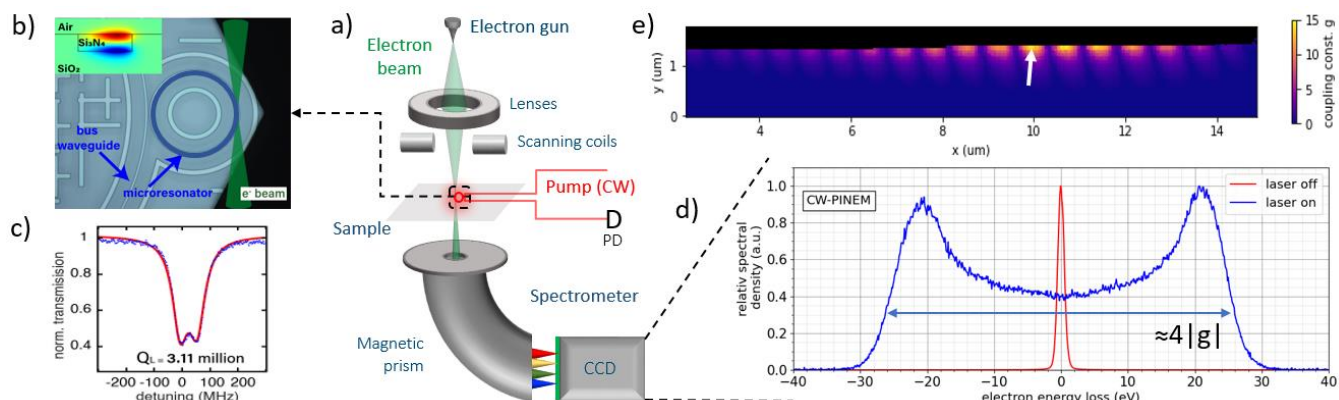


Figure 1. Figure 1. a) Inside of a transmission electron microscope (TEM), an electron beam passes close to a CW-pumped integrated Si₃N₄ microresonator in a loop geometry (b, SEM image: side view of the chip, Inset: Electric field component parallel to the electron beam.) c) Optical transmission spectrum of a typical microresonator resonance exhibiting a linewidth of ~70 MHz (Q-factor of $3.11 \cdot 10^6$). d) Electron energy distribution of the transmitted beam recorded with a magnetic-prism spectrometer. Tuning the CW laser to a cavity resonance, a spectral broadening up to ~50 eV is observed (probing position indicated by the white arrow in e)). e) Spatial map of the electron-light coupling constant g , obtained from scanning the electron beam across the resonator.

References

- [1] R. J. England *et al.*, *Rev. Mod. Phys.* **86**, 1337–1389 (2014).
- [2] K. E. Priebe *et al.*, *Nat. Photonics* **11**, 793 (2017).
- [3] Y. Morimoto and P. Baum, *Nat. Phys.* **14**, 252 (2018).
- [4] M. Kozáket *et al.*, *Phys. Rev. Lett.* **120**, 103203 (2018).
- [5] O. Schwartz *et al.*, *Nat. Methods* **16**, 1016–1020 (2019).
- [6] G. M. Vanacore *et al.*, *Nat. Mater.* **18**, 573–579 (2019).
- [7] A. Feist *et al.*, *Phys. Rev. Research* **2**, 043227 (2020).
- [8] B. Barwick, D. J. Flannigan, and A. H. Zewail, *Nature* **462**, 902 (2009).
- [9] F. J. García de Abajo, *et al.*, *Rev. Mod. Phys.* **82**, 0034–6861 (2010).
- [10] O. Kfir *et al.*, *Nature* **582**, 46–49 (2020).
- [11] K. Wang *et al.*, *Nature* **582**, 50–54 (2020).
- [12] A. Feist *et al.*, *Ultramicroscopy* **176**, 63 (2017).
- [13] J. Liu *et al.*, *Optica* **5**, 1347–1353 (2018).
- [14] K. E. Echternkamp *et al.*, *Nature Physics* **12**, 1000–1004 (2016)
- [15] O. Kfir *et al.*, *Phys. Rev. Lett.* **123**, 10 (2019).