

Lithium Depletions in Late-type Dwarfs

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Abstract. Using a self-consistent dynamic theory of non-local convection in chemically inhomogeneous stars, the lithium depletion in MS stars of masses 0.725–1.5 M_{\odot} are calculated. Both of the overshooting and microdiffusion are included in a consistent way. The comparisons of theoretical results with the observed Li depletions in open clusters show that the general characters of Li depletions can be reproduced by theory. The overshooting mixing and microdiffusion induced by gravitational setting and radiative accelerations may be two main mechanisms of Li depletion.

Keywords. Convection- star: abundance-star: late-type-stars: evolution- open clusters and associations: general

It is well known that the Li abundance is very significant for the theory of big-bang nucleosynthesis and tracking the extension of the surface convection zone in stars during the course of their evolution.

The mechanism of Li depletion is still not completely known. A number of theoretical approach has been presented so far, such as the mass loss (Weymann & Sears, 1965; Schramm *et al.*, 1990), wave-driven mixing (García López & Spruit, 1991; Montalbán & Schatzman, 1996), rotationally induced mixing (Charboneau *et al.*, 1992; Chaboyer *et al.*, 1995; Pinsonneault, *et al.*, 1999), microdiffusion and turbulent mixing (Michaud, 1986; Turcotte *et al.*, 1998), overshooting mixing (Straus *et al.*, 1976; Xiong & Deng, 2001) and so on. However, there is no any one single mechanism which can interpret all the known observed characters of Li depletion. The Li depletion, in our opinion, should be the result of several physical processes, rather than from a single mechanism. The overshooting mixing and microdiffusion induced by gravitation setting and radiation acceleration seem to be a reasonable combination.

Using our non-local convection theory in chemically inhomogeneous stars (Xiong, 1981), we calculated the Li depletions for 20 series of evolutionary models with mass of 0.725–1.5 M_{\odot} . The solar abundance ($X=0.70$, $Z=0.02$) is assumed. The overshooting mixing and microdiffusion induced by gravitational setting are included in a self-consistent way. The fundamental equations and the numerical scheme will be described in our latter paper (Xiong & Deng, 2008). Figure 1 illustrates the Li abundances as function of stellar age for the evolutionary models with different masses as marked on the curves. In Fig. 1a the microdiffusion is neglected. In Fig. 1b both of the microdiffusion and overshooting mixing are included. It can be seen that the Li abundances decrease (approximately) exponentially with age of stars. The e-folding time of Li depletion are not a monotonic function of stellar mass. It achieves the maximum at $M \approx 1.1M_{\odot}$. This is due to the fact that the Li depletion is resulted dominantly from macrodiffusion for warm stars ($M \geq 1.1M_{\odot}$ or $Te \geq 6100K$), and it is resulted dominantly from the overshooting mixing for the cool stars ($M \leq 1.0M_{\odot}$ or $Te \leq 5800K$). The sign \odot in Figs. 1 & 3d is

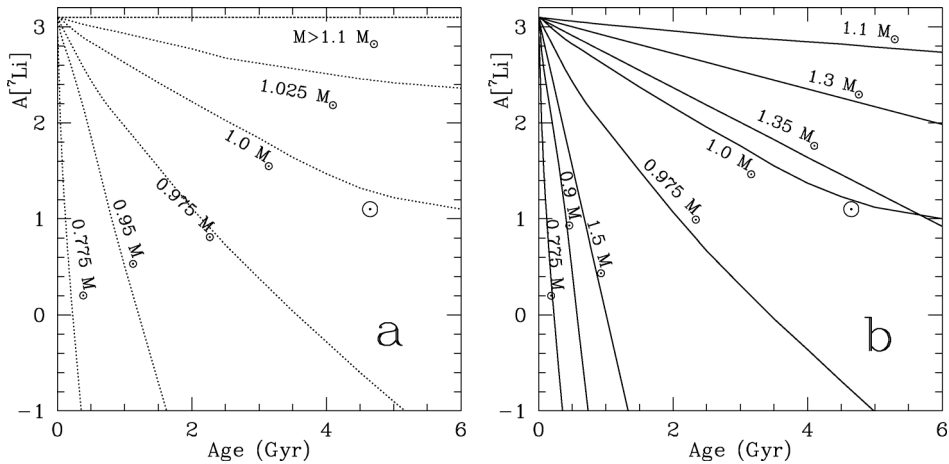


Figure 1. The Li abundances as functions of star's age for the different masses marked on the curves. (a) the microdiffusion is neglected, (b) both of the microdiffusion and overshooting mixing are included.

the observed location of the Sun. It is necessary to indicate that the depth of our solar model agrees with the requirement of helioseismology.

Figs. 2a – d show the lithium distribution in stellar interior at different ages. The short dashed lines mark the bottom boundary of convectively unstable zone. The sharply drop at bottom is due to violent burn of lithium. It can be seen that the convective overshooting (the shadowy regions) penetrates deeply into the convectively stable zone and the full mixing zone is extended greatly. For the warm stars with mass greater than about $1.1 M_{\odot}$ (Figs. 2a & b), the convection zone is too shallow. The gravitational setting is not enough to bring the surface lithium to the deep burn region. It drive the surface lithium to the deeper interior of stars and the surface lithium will be stored in the lower radiative region under the overshooting zone, so the surface lithium abundance decreases. The convection zone becomes deeper with decrease of stellar temperature. overshooting brings the surface lithium to the deep burn region, therefore, the surface lithium abundance decreases (Figs 2c & d). We can know from the above analysis that the mechanisms of lithium depletion are different for the warm and cool stars. For the warm stars the microdiffusion is domanical, however the overshooting mixing becomes domanical for the cool stars.

Figure 3 give The Li abundance vs. the effect temperature for the clusters with different ages, and the corresponding theoretical isochne lines of Li abundance are also drawn here. It can be seen from Fig. 3 that within the observational uncertain the theoretical results reproduced roughly the general profiles of Li depletion except the following points: (1) The Li depletions seem to be underestimated for stars arround $Te \approx 5900K$. It is possible there are depletion mechanisms ignored in this work; (2) In Fig. 3d the Li abundances of five stars on the right of theoretical isochne line seem to be too high. However, two of them are only given the upper limits of their Li abundance. The another two stars, named as I-2a and I-2b in the Table 1 of Jones *et al.* (1999), are the two components of a double-lined spectroscopic binary. Their Li abundance are quite uncertain. The samples of lower temperature star are too few for giving a confident conclusion.

It is the conclusion that the overshooting mixing and microdiffusions induced by gravitation setting and radiative acceleration may be two main mechanisms of Li depletion. Not only the convective overshooting interpret the lithium depletion of cool stars but

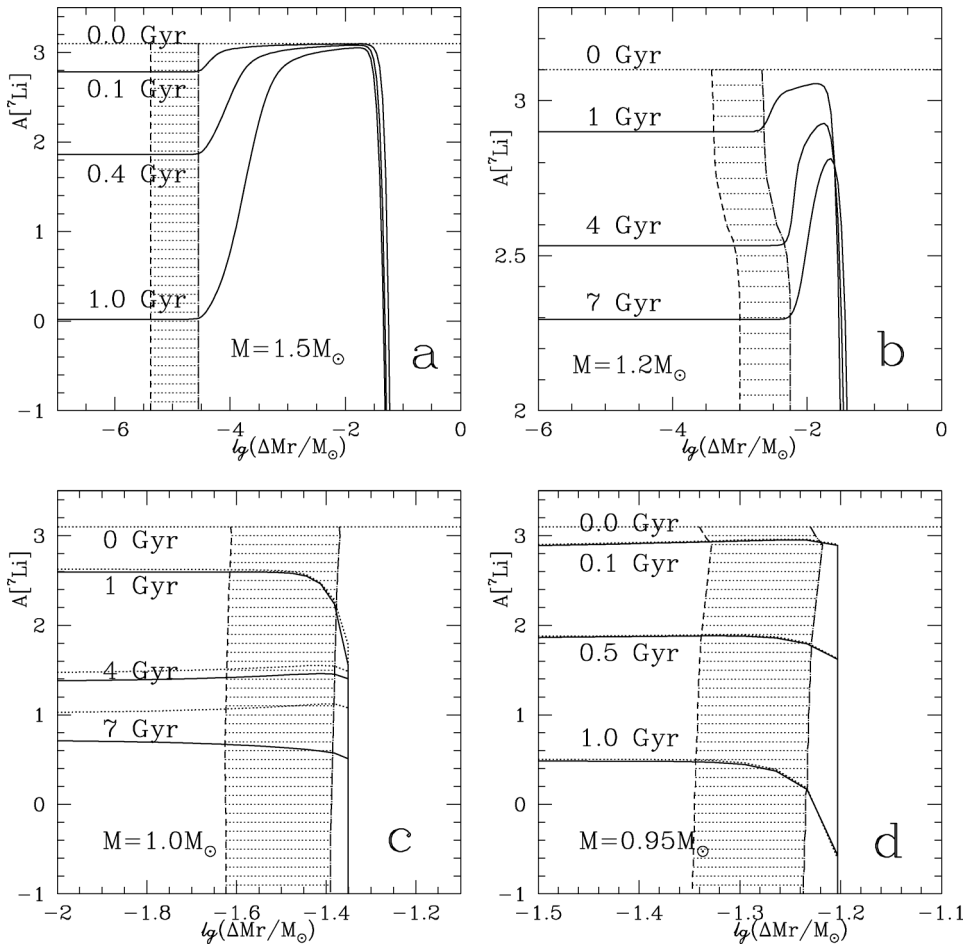


Figure 2. The lithium distribution in stellar interior at different ages, (a) $M = 1.5M_{\odot}$ (b) $M = 1.2M_{\odot}$ (c) $M = 1.0M_{\odot}$ (d) $M = 0.95M_{\odot}$

also it increase the mass of the surface full mixing region, therefore the too quick depletion of lithium in warm stars predicted by single microdiffusion theory is removed automatically. It is not necessary to induce a phenomenological turbulent diffusion such as induced by Proffitt & Michaut (1991) and Richard, Michaud & Richer (2005). In other words, convective overshooting is just the so called turbulet diffusion needed by them.

Acknowledgements

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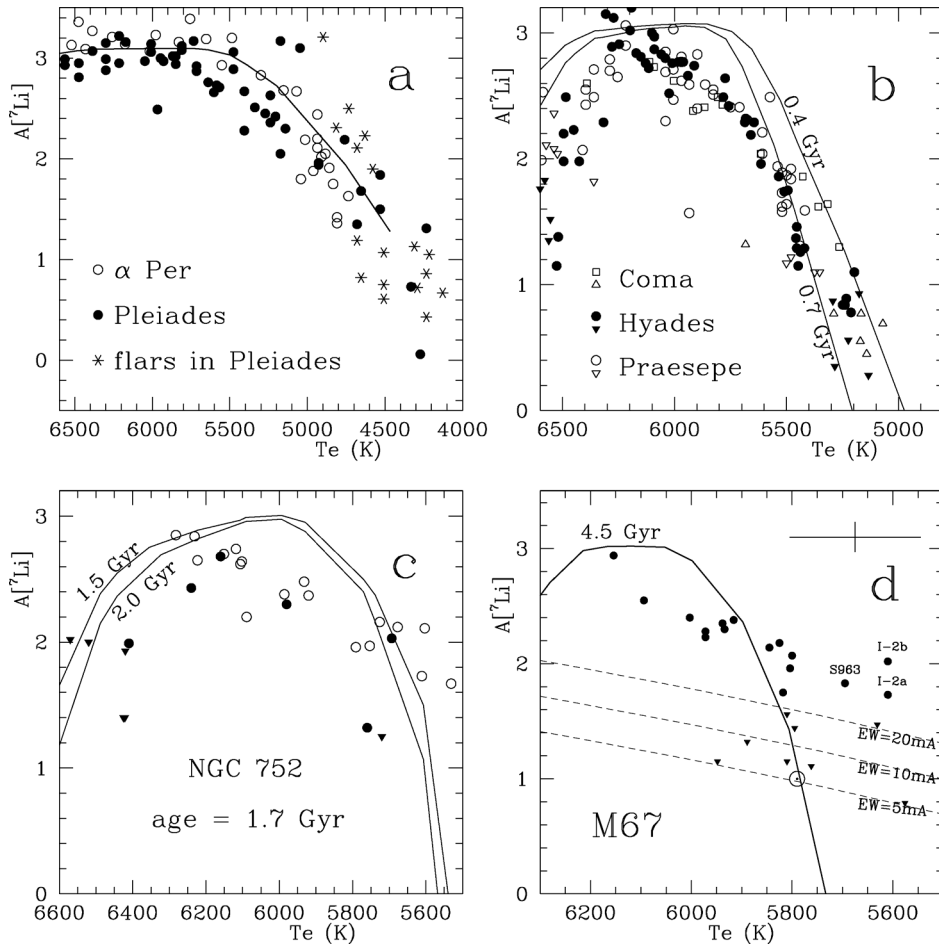


Figure 3. The Li abundances versus effective temperature of stars for (a) α Persei, Pleiades, (b) Coma Berenices, Hyades Praesepe, (c) NGC 752, and (d) M67 clusters. The triangles and inverse triangles give the upper limits for Li abundance. The estimated uncertainties of Li abundance and effect temperature are indicated in the upper-right corner of the plots.

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Discussion

ASPLUND: Have you investigated how much ${}^7\text{Li}$ depletion your convective overshooting model predict for metal-poor turn-off stars?

XIONG: We have only considered the depletion of Lithium for solar type stars. We understand that Lithium depletion in metal-poor stars is extremely important in exploring

the initial cosmological Lithium abundance and nucleosynthesis of Big-Bang. Due to the lack of reliable evolutionary models of metal-poor stars, such work can only be done in a future time when such models become available.

LANGER: Are there any free parameters in your model for overshooting? and: Are your models (for main sequence stars) in strict thermal equilibrium?

XIONG: 1). Indeed, there is also a free parameter in our statistical theory of non-local turbulent convection, which is similar to that of MLT representing the characteristic length (wave number) of the turbulent spectrum. However, such a parameter is not at all “free”. Instead, it has to be regulated by observational constraints from stellar evolution, depletion of Lithium in stars, helioseismology and stellar pulsation, etc. As we have shown in our previous work, the theory stands against all these challenges.

2). The models computed so far are all located on the main sequence, therefore the assumption of perfect thermal equilibrium is reasonable.

KUPKA: Have you also made any computations for the Li abundance in hotter (main sequence) stars than the ones you have shown here? I mean early F type and late A type, and compared to some data? This could be useful to test the model for microscopic diffusion you have used in your computations in the OV zone of cool main sequence stars too!

XIONG: We have done Lithium depletion models only for late F type – K type stars so far. The late A type – early F type stars are not modelled, mainly because we could not good monochromatic opacity. For microdiffusion, only gravitational settling has been taken into account. Radiative acceleration has been excluded. Once monochromatic opacity is available, radiative acceleration can be included into microdiffusion easily. In that case, the combination of microdiffusion and overshooting mixing can be well applied to the studies of Li depletion and A_p , F_m phenomena in A-F type stars. Turbulent diffusion is very likely not needed, pretty the same way as in the present work that overshooting can naturally reproduce the effect of turbulent diffusion