

# COMPACT DOUBLES: TESTING THE LENSING HYPOTHESIS

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**Abstract.** If a compact double (CD) is caused by lensing, the orientation and/or flux ratio of its components can substantially change in  $\approx 1$  year.

## 1. Compact doubles as gravitationally lensed images

Major surveys using VLBA are underway to find lensed radio sources with separations  $\approx 1 - 100$  milliarcseconds (e.g., Patnaik et al. 1995; Wilkinson 1995). Further, it has been suggested that some of the known CDs (Phillips & Mutel 1982) may in fact be such lensed images (e.g., Ostriker 1995). The required lensing agents of  $\sim 10^6 - 10^7 M_{\odot}$  may be remnants of a generation of pregalactic stars or dwarf galaxies. Since the individual components of CDs are  $\approx 1$  mas, the lensed object is likely to be the radio core of a quasar. Such cores are identified with the base of a relativistic jet pointed roughly towards us (Blandford & Konigl 1979), which is itself often resolved by VLBI into one or more bright emission knots apparently separating at superluminal velocities,  $v \sim 5 - 10c$ , from a nucleus near the jet's origin (e.g., Vermeulen & Cohen 1994). We suggest that such superluminal motion of radio knot(s) in the lensed source can lead to significant structural variations in the small-separation images, even on time scales of  $\approx 1$  yr. VLBI monitoring can thus help in distinguishing any (*milli-*) lensed CD images from genuine CDs.

## 2. Temporal effects in the lensing scenario

For evaluating the temporal changes we concentrate on just the superluminal components (knots); they would dominate the core structure except in the strongly self-absorbed spectral regime. Consider a superluminal knot being multiply imaged by a point-like lens of mass  $M$ . The angular positions

TABLE 1. Evolution of PA ( $\Phi$ ) and flux-ratio ( $R$ ) for two ejection directions ( $\psi$ ).

t (yr)	$\psi = 60^\circ$		$\psi = 120^\circ$		t (yr)	$\psi = 60^\circ$		$\psi = 120^\circ$	
	$\Phi$ (deg)	R	$\Phi$ (deg)	R		$\Phi$ (deg)	R	$\Phi$ (deg)	R
0	0	1.5	0	1.5	3	37	2.4	79	1.7
1	19	1.7	30	1.4	4	41	2.8	90	2.0
2	30	2.0	60	1.5	5	44	3.4	97	2.5

of the two images,  $\alpha$ , measured from the lens (which is the origin of our coordinate system) are related to the source position  $\beta$  by :  $\alpha^2 - \beta\alpha - \theta_L^2 = 0$ . Here  $\theta_L = 5(D_{LS}/D_L)^{1/2}(D_S/1000\text{Mpc})^{-1/2}(M/3 \times 10^6 M_\odot)^{1/2}$  mas is the Einstein ring radius, and  $D_S, D_L$  and  $D_{LS}$  are the source, lens and lens-source angular-diameter distances. For a  $3 \times 10^6 M_\odot$  lens located roughly midway between us and a source  $\sim 1000$  Mpc away the image separation  $\Delta\alpha \sim 2\theta_L \sim 10$  mas, which is characteristic of CDs (Carvalho 1985). Another characteristic is the flux ratio of the components,  $R \sim 1$  to 2 (note that  $R \simeq 1.5$  if  $\beta \simeq \theta_L/5$ ).

Now, suppose an emission knot initially located on the  $x$ -axis at  $x = x_A$ , moves with  $v = \gamma c$  at an angle  $\psi$  to the  $x$ -axis. After a time  $t$  :  $\beta^2(t) = x_A^2 + 2lx_A \cos \psi + l^2$ , where  $l = \dot{\theta}t = \gamma ct/D_S$  with  $\dot{\theta} = 0.5(\gamma/8)(D_S/1000\text{Mpc})^{-1}$  mas/yr. The flux ratio is then:  $R(t) = (2 + u(t)^{1/2} + u(t)^{-1/2})/(u^{1/2} + u^{-1/2} - 2)$ , where  $u = [1 + 4\theta_L^2/\beta^2]$ . The position angle (or PA)  $\Phi(t)$ , of the line joining the two images relative to the  $x$ -axis varies as :  $\tan \Phi(t) = (l \sin \psi)/(x_A + l \cos \psi)$ . The evolution of these quantities is given in Table 1, assuming characteristic values  $\theta_L = 5$  mas,  $x_A = \theta_L/5 = 1$  mas and  $\dot{\theta} = 0.5$  mas yr $^{-1}$  (see above). Clearly, substantial variations can occur in the PA and  $R$  of the CD, if indeed it results from gravitational milli-lensing of a quasar core. Hence, by imaging a CD with VLBI arrays even  $\approx 1$  year apart one can verify if it is merely an illusion caused due to milli-lensing.

## References

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