# MEASURING THE LAYER CHARGE OF DIOCTAHEDRAL SMECTITE BY O–D VIBRATIONAL SPECTROSCOPY

ARTUR KULIGIEWICZ<sup>1</sup>, ARKADIUSZ DERKOWSKI<sup>1,\*</sup>, KATJA EMMERICH<sup>2</sup>, GEORGE E. CHRISTIDIS<sup>3</sup>, CONSTANTINOS TSIANTOS<sup>4</sup>, VASSILIS GIONIS<sup>4</sup>, AND GEORGIOS D. CHRYSSIKOS<sup>4,\*</sup>

<sup>1</sup> Institute of Geological Sciences, Polish Academy of Sciences, ul. Senacka 1, 31-002 Krakow, Poland
<sup>2</sup> Competence Center for Material Moisture (CMM) and Institute of Functional Interfaces (IFG), Karlsruhe Institute of Technology, Hermann-von-Helmholtz-Platz 1, D-76344 Eggenstein-Leopoldshafen, Germany

<sup>3</sup> School of Mineral Resources Engineering, Technical University of Crete, Chania, Greece 73100

<sup>4</sup> Theoretical and Physical Chemistry Institute, National Hellenic Research Foundation, 48 Vass. Constantinou Ave., Athens,

Greece 11635

Abstract—Layer charge (LC) is a fundamental property of smectite but its measurement remains challenging and tedious to apply on a high-throughput basis. The present study demonstrates that the position of a sharp, high-energy O–D stretching band of adsorbed D<sub>2</sub>O (vO–D, at ~2686–2700 cm<sup>-1</sup>), determined by infrared spectroscopy, correlates with LC and provides a simple method for its measurement. Twenty nine natural dioctahedral smectites and 14 reduced-charge montmorillonites with LC determined previously by different methodologies were saturated with D<sub>2</sub>O and examined by attenuated total reflectance infrared spectroscopy (ATR-IR). The samples included smectites in Mg, Ca, Na, Li, K, and Cs forms and covered the full range of the smectite LC (0.2 to 0.6 *e* per formula unit). Statistically significant correlations were found between vO–D and LC values determined with each of the two main methods of LC determination: the structural formula method (R<sup>2</sup> = 0.96,  $\sigma$  = 0.02, ~0.2 < LC < 0.6) and the alkylammonium method (R<sup>2</sup> = 0.92,  $\sigma$  = 0.01, 0.27 < LC < 0.37). These correlations were found between vO–D and LC values determined with each of the tassed on Li- and Na-saturated smectites, respectively, but other cationic forms can be employed provided that the exchangeable cations are of sufficiently high hydration enthalpy (*e.g.* Mg<sup>2+</sup> or Ca<sup>2+</sup>, but not K<sup>+</sup> or Cs<sup>+</sup>). The new method is fast, low-cost, implemented easily in laboratories equipped with ATR-FTIR, and applicable to samples as small as ~5 mg.

Key Words—ATR, D<sub>2</sub>O, Infrared Spectroscopy, Layer Charge, Smectite.

# INTRODUCTION

# Layer charge in smectite

The Clay Minerals Society (CMS) Nomenclature Committee defines smectite as a group of 2:1 phyllosilicates with layer charge (LC) between ~0.2 and 0.6 negative elementary charge per  $O_{10}(OH)_2$  formula unit (p.f.u.; Güven, 1988; Guggenheim *et al.*, 2006). Layer charge is defined as "the total negative charge deviation from an ideal, unsubstituted dioctahedral or trioctahedral composition" (CMS Nomenclature Committee, 2015). Layer charge controls many properties of smectite, such as the cation exchange capacity (CEC), water adsorption, rheological behavior, *etc.* (*e.g.* Laird, 1999; Harvey and Lagaly, 2006; Laird, 2006; Christidis *et al.*, 2006; Ferrage *et al.*, 2007; Tertre *et al.*, 2015). For these reasons, LC is a key property to determine when characterizing smectite.

Fast and accurate LC determination of smectite remains a challenge. In turn, the way LC is measured influences the specific terminology developed to describe LC in a manner that may be confusing. A

\* E-mail address of corresponding author: ndderkow@cyf-kr.edu.pl; gdchryss@eie.gr DOI: 10.1346/CCMN.2015.0630603 portion of the measured smectite charge that is caused by heterovalent, octahedral or tetrahedral substitutions (as in the aforementioned CMS definition of layer charge) is referred to as 'permanent.' Charge originating from OH groups developed at the edges of the 2:1 layer and local structural defects depends on pH and is referred to as 'variable.' The term 'total charge' is, therefore, used hereafter to describe the sum of permanent and variable charges, whereas the term 'charge' is used without specifying charge origin or location. On a purely geometrical basis, the terms 'interlayer charge,' 'basal charge,' and 'edge charge' are used when referring to a charge present in the interlayer, on basal surfaces, and on edges of 2:1 layers, respectively. Although the terms 'variable charge' and 'edge charge' refer to very similar portions of total charge, they may not be directly interchangeable (Kaufhold, 2006; Delaverhne et al., 2015).

#### Layer-charge measurement methodology

Various methods have been developed to measure the charge of smectites. These involve a chemical analysis and/or CEC measurement of pure, monomineralic smectite fractions (Čícel and Komadel, 1994; Laird, 1994; Christidis, 2008); the exchange with alkyl-ammonium cations of different chain lengths (Lagaly and Weiss, 1969; Ruehlicke and Kohler, 1981; Lagaly,

1981, 1994; Olis *et al.*, 1990; Dohrmann *et al.*, 1999; Wolters *et al.*, 2009); or the saturation of smectite with particular inorganic cations solvated by polar liquids in combination with X-ray diffractometry (XRD; Christidis and Eberl, 2003). Spectroscopic methods have also been applied to a suspension of smectite in cationic dye solutions (Bujdák, 2006) or to smectites saturated with ammonium cations (Petit *et al.*, 1998).

Structural formula methods (SFM) aim to calculate the deviations from the ideal, unsubstituted dioctahedral or trioctahedral compositions, distributing analyzed elements on the basis of certain assumptions (Stevens, 1946). The SFM are biased by the presence of impurities such as quartz, opal, and feldspars. Interference from impurities can be avoided with electron microprobe techniques (Christidis and Dunham, 1993, 1997), which provide the range of LCs in smectite particles but, inevitably, add to the complexity of the measurement. Furthermore, SFM can be influenced by the presence of non-exchangeable, non-structural cations (Kaufhold *et al.*, 2011), variable charges, and local domains of different octahedral occupancy (*cf.* Wolters *et al.*, 2009).

The alkylammonium method (AAM) probes only the interlayer charge density (Lagaly, 1994). The original methodology returns an entire LC distribution based on the specific conformation of different long-chain alkyl-ammonium cations. Smectite LC is then calculated as a mean value of the LC distribution. The measurement, however, can be influenced by uncertainties in the packing of the alkylammonium cations in the interlayer and by particle-size effects (Laird *et al.*, 1989; Laird, 1994). Additionally, the AAM method is labor intensive, time consuming, and requires a relatively large quantity of sample (>1 g). Modifications of AAM proposed by Olis *et al.* (1990) and Dohrmann *et al.* (1999) aimed to reduce these drawbacks.

The comparison of charge values obtained with SFM and AAM is not straightforward because these methods are based on different assumptions, they are subject to different types of error, and they measure different portions of total charge (*e.g.* Laird, 1994; Kaufhold *et al.*, 2002; Kaufhold, 2006; Christidis, 2008; Kaufhold *et al.*, 2011). The AAM returns lower charge values than the SFM-based or calibrated methods (Laird, 1994; Kaufhold, 2006; Wolters *et al.*, 2009; Kaufhold *et al.*, 2011). Laird (1994) and Kaufhold (2006) have demonstrated experimentally that SFM- and AAM-derived charges are correlated linearly, but that linearity may also depend on the aspect ratio of the smectite crystallites (Delavernhe *et al.*, 2015).

Whereas SFM and AAM provide a value of the layer charge directly and do not require independent calibration, all other LC determination methods involve calibration against one of the two or, with certain assumptions, against the CEC. For example, the XRDbased method developed by Christidis and Eberl (2003) probes only the interlayer charge density and was calibrated against SFM charges. This method is, however, limited to smectites with charge greater than 0.39 p.f.u. The infrared spectroscopic method of Petit et al. (1998) correlated CEC with the relative integrated intensity of the  $v_4$  NH<sup>+</sup><sub>4</sub> bending envelope against the Si-O stretching envelope in a NH<sub>4</sub>-exchanged smectite. In combination with a Li<sup>+</sup> fixation pretreatment, this method can provide an additional estimate of charge location (octahedral vs. tetrahedral; Petit et al., 2006). This method can be applied to small samples (<5 mg) but is sensitive to the presence of silicate impurities. The electronic spectra of cationic dyes such as methylene blue (Bujdák, 2006; Czímrerová et al., 2006; Pentrák et al., 2012) and rhodamine 6G (Bujdák et al., 2003, 2004) proved to be very sensitive indicators of LC changes in clays, although mostly on a qualitative level.

#### Water molecules as a probe for smectite surface

Whereas some of the aforementioned probes of layer charge are extrinsic, water is a ubiquitous and intrinsic probe of the surface of clay minerals (Sposito et al., 1983; Johnston et al., 1992; Xu et al., 2000; Schoonheydt and Johnston, 2006; Johnston, 2010). A recent infrared spectroscopic study of D<sub>2</sub>O-saturated smectites (Kuligiewicz et al., 2015) reported that the position of the high-frequency O-D stretching band observed at 2680-2695 cm<sup>-1</sup> (hereafter, vO-D, corresponding to an O–H band of H<sub>2</sub>O at 3620-3635 cm<sup>-1</sup>) seemed to correlate linearly with total charge. The vibrational spectra of adsorbed water (as D<sub>2</sub>O) have, therefore, the potential to be adapted for a quick and reliable determination of layer charge in smectites. The high-frequency vO-D band, first described by Farmer and Russell (1964) and Russell et al. (1970), was assigned to the stretching modes of so-called "free O-D" ("dangling mode") groups of D<sub>2</sub>O molecules positioned in the vicinity of the siloxane surface and nearly perpendicular to the plane of basal oxygen atoms (Figure 1; Russell et al., 1970; Suquet et al., 1977). This interpretation is supported by the observations of highfrequency O-H stretching bands of H<sub>2</sub>O at the interfaces with hydrophobic phases such as silica (Jena and Hore, 2010), tetrachloromethane (Scatena et al., 2001), or at water-vapor interfaces (Sovago et al., 2009; Tian and Shen, 2009), where corresponding bands were assigned to the dangling O-H. In the case of the smectite-water interface, the bond strength and the stretching frequency of this dangling O-H (O-D) bond are probably controlled by the entire basal charge, regardless of the charge origin from the octahedral or tetrahedral sheet (Kuligiewicz et al., 2015; cf. Farmer and Russell, 1971; Suguet et al., 1977).

The aim of the present work was to develop a new method of layer-charge determination by further testing the correlation of vO-D with LC on a large number of dioctahedral smectite samples with charge determined previously by broadly accepted techniques. The



Figure 1. Schematic orientation of water molecules in the vicinity of the basal oxygen plane of a smectite 2:1 layer. The oriented dangling O–D bonds at the water-smectite interface (shaded zone) are involved in the vO–D band, as suggested by Russell *et al.* (1970), Suquet *et al.* (1977), and Kuligiewicz *et al.* (2015). Distances and atomic radii not to scale.

differences between particular charge-measurement reference methods have been discussed extensively in the literature (*e.g.* Mermut, 1994; Laird, 1994; Kaufhold *et al.*, 2002; Kaufhold, 2006; Christidis, 2008; Wolters *et al.*, 2009; Kaufhold *et al.*, 2011; Kaufhold and Dorhmann, 2013) and are beyond the scope of the present study.

### MATERIALS AND METHODS

#### Samples

A collection of 29 natural smectites in various cationic forms and fourteen reduced-charge Li-montmorillonites (RCMs) were examined. The samples studied were grouped in four sets (Table 1):

(A) Fine fractions (<1 µm and finer) thoroughly separated from six smectite reference samples from the Source Clays Repository of The Clay Minerals Society, which included montmorillonites (SAz-2, SCa-3, SWy-2), beidellites (SbCa-1, SbId-1), and a nontronite/ ferric-smectite (SWa-1). The same specimens were studied previously by Kuligiewicz et al. (2015) and were included here in additional cationic forms for comparison purposes. The total layer charge for this set was determined on XRD-pure smectites (except SbId-1, where minor kaolinite content was subtracted virtually) from the Na content of the Na-exchanged smectite fraction measured by flame photometry after conventional acid digestion (Kuligiewicz et al., 2015). The total charge estimation method applied followed the general SFM definition.

(B) Fifteen natural dioctahedral smectites (<0.2  $\mu$ m fractions, Na-exchanged) from the study of Wolters *et al.* (2009). Those authors reported the charges of the

same specimens by the AAM (Lagaly, 1994), a simplified alkylammonium method (AAOM; Olis *et al.*, 1990), and SFM (Stevens, 1946).

(C) Six montmorillonite samples (SAz-1, SWy-1, SCa-3, STx-1, Woburn, Kinney; <2  $\mu$ m size fractions) originally used in the study of Christidis and Eberl (2003) and two clay fractions (<2  $\mu$ m) of bentonites from Cyprus (FEOG and Skouriotissa, referred to hereafter as SK; Christidis, 2006). Both SFM charges and charges calculated by the method of Christidis and Eberl (2003; CEM) were reported by Christidis and Eberl (2003), Christidis (2006), and Skoubris *et al.* (2013). The spectroscopic data reported here were obtained from the Li forms of these fractions (Skoubris *et al.*, 2013).

(D) A series of 13 portions of Li-SAz-1 (<2  $\mu$ m fraction) montmorillonite subjected to increasing charge reduction by heating for 24 h at 60–180°C from the study of Skoubris *et al.* (2013). These samples were labeled Li-SAz1-xxx, where xxx denotes the heating temperature. The total charge of these samples was measured with CEM for temperatures up to 135°C (Skoubris *et al.*, 2013). An additional set of charges for the whole sample set was calculated by multiplying the SFM charge of the unheated material by the % reduction of the CEC upon heating, reported by Skoubris *et al.* (2013).

#### Mid-infrared spectroscopy

The infrared (IR) spectra of all samples were collected in the  $4000-580 \text{ cm}^{-1}$  range on two FTIR instruments (Equinox 55 by Bruker Optics, Ettlingen, Germany and Nicolet 6700 by Thermo Scientific, Waltham, Massachusetts, USA) with three different single-reflection diamond Attenuated Total Reflectance

Table 1. Positions of the sharp,	high-frequency O-D	stretching band,	vO-D, and layer	charge of the s	smectite samples used in
the present study.					

Set	Sample	Charge location <sup>1</sup>	Interlayer cation	vO-D (cm <sup>-1</sup> )	$1\sigma$ (cm <sup>-1</sup> )		Layer charge (p.f.u.)	•
							SFM <sup>2</sup>	
	SAz-2	0	Mg Ca Na K	2686.2 2686.0 2688.0 2688.5	0.5 0.5 0.5 0.5		0.52	
	SCa-3	0	Cs Mg Ca Na K Cs	2690.0 2685.1 2686.1 2686.2 2687.7 2689.4	0.5 0.5 0.5 0.5 0.5 0.5		0.51	
	SbCa-1	Т	Ca Na K	2685.5 2684.5 2687.2	0.5 0.5 0.5		0.50	
А	SbId-1	Т	Mg Ca Na K Cs	2688.8 2689.3 2689.3 2690.6 2691.6	0.5 0.5 0.5 0.5 0.5		0.39	
	SWy-2	O/T	Mg Ca Na K Cs	2691.8 2691.8 2692.4 2692.3 2693.2	0.5 0.5 0.5 0.5 0.5		0.36	
	SWa-1	O/T	Mg Ca Na K Cs	2691.7 2692.1 2692.9 2693.7 2691.4	0.5 0.5 0.5 0.5 1		0.31	
В	2LP 3, 7 <sup>th</sup> Mayo 4JUP 5MC 6GPC 7EMC 8UAS 14TR03 17GR02 19USA02 21D01 28SB 31BAR 3 32Volclay 33CA	0/T 0/T 7/0 0/T 0/T 0/T 0/T 0/T 0/T 0/T 0/T 0/T	Na	2692.3 2691.0 2692.1 2692.1 2690.4 2693.5 2689.0 2686.5 2692.5 2690.9 2691.2 2688.3 2690.4 2693.7 2690.9	$\begin{array}{c} 0.2 \\ 0.2 \\ 0.4 \\ 0.2 \\ 0.4 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \\ 0.2 \end{array}$	SFM <sup>3</sup> 0.37 0.42 0.37 0.40 0.45 0.40 0.48 0.47 0.38 0.39 0.42 0.47 0.41 0.39 0.42	$\begin{array}{c} AAM^4 \\ 0.28 \\ 0.30 \\ 0.29 \\ 0.28 \\ 0.33 \\ 0.28 \\ 0.34 \\ 0.37 \\ 0.29 \\ 0.30 \\ 0.29 \\ 0.30 \\ 0.29 \\ 0.36 \\ 0.32 \\ 0.27 \\ 0.30 \end{array}$	$\begin{array}{c} \text{AAOM}^5 \\ 0.29 \\ 0.31 \\ 0.30 \\ 0.28 \\ 0.34 \\ 0.29 \\ 0.33 \\ 0.38 \\ 0.30 \\ 0.31 \\ 0.31 \\ 0.39 \\ 0.32 \\ 0.27 \\ 0.30 \end{array}$
С	Li-SAz-1 Li-SWy-1 Li-Kinney Li-SK Li-FeOG Li-SCa-3 Li-STx-1 Li-Woburn	0 0/T 0 T/0 0 0 0/T 0/T	Li	2686.9 2693.7 2687.4 2689.6 2692.9 2686.9 2691.1 2692.5	0.2 0.2 0.2 0.3 0.2 0.2 0.3	$\begin{array}{c} {\rm SFM}^6 \\ 0.56 \\ 0.36 \\ 0.49 \\ 0.48 \\ 0.39 \\ 0.60 \\ 0.44 \\ 0.42 \end{array}$	$\begin{array}{c} \text{CEM}^7 \\ 0.55 \\ 0.39 \\ 0.53 \\ 0.50 \\ 0.39 \\ 0.60 \\ 0.44 \\ 0.43 \end{array}$	

447	
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Table1	(contd.)
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Set	Sample	Charge location <sup>1</sup>	Interlayer cation	vO-D (cm <sup>-1</sup> )	$1\sigma$ (cm <sup>-1</sup> )	Layer (p.:	charge f.u.)	
						SFM <sup>6</sup>	CEM <sup>7</sup>	
	Li-SAz1-60			2687.2	0.2	0.54	0.54	
	Li-SAz1-80			2687.7	0.2	0.54	0.53	
	Li-SAz1-100			2689.3	0.2	0.52	0.50	
	Li-SAz1-105			2689.3	0.2	0.49	0.46	
	Li-SAz1-110			2690.6	0.2	0.47	0.44	
	Li-SAz1-115			2691.3	0.2	0.44	0.44	
D	Li-SAz1-120	Ο	Li	2694.0	0.3	0.39	0.42	
	Li-SAz1-125			2694.4	0.3	0.37	0.40	
	Li-SAz1-130			2695.2	0.4	0.31	0.39	
	Li-SAz1-135			2696.4	0.5	0.26	0.39	
	Li-SAz1-140			2697.3	0.5	0.24	N.A. <sup>8</sup>	
	Li-SAz1-160			2697.6	0.6	0.22	N.A. <sup>8</sup>	
	Li-SAz1-180			2699.0	0.7	0.19	N.A. <sup>8</sup>	

<sup>1</sup> O: Octahedral, T: tetrahedral

<sup>2</sup> Total charge from Na content of Na-exchanged pure smectites (Kuligiewicz et al., 2015; SFM equivalent)

<sup>3</sup> Total charge by structural formula method (SFM) of Stevens (1946)

<sup>4</sup> AAM – alkylammonium method (Lagaly, 1994)

<sup>5</sup> AAOM – simplified alkylammonium method (Olis *et al.*, 1990)

<sup>6</sup> Total charge recalculated from SFM data of unheated sample and CEC reduction of RCM (SFM equivalent);

 $^{7}$  CEM – XRD method of Christidis and Eberl (2006).

<sup>8</sup> The method could not be applied

Available charges were compiled from Kuligiewicz et al. (2015), set A; Wolters et al. (2006), set B; Christidis and Eberl (2003), Christidis (2006), and Skoubris et al. (2013), set C; Skoubris et al. (2013), set D.

(ATR) accessories, employed with their pressure tips removed (DuraSampl IR II by SensIR Technologies, presently Smiths Detection, Edgewood, Maryland, USA; Golden Gate by Thermo Scientific, Waltham, Massachusetts, USA; and MIRacle by PIKE Technologies, Madison, Wisconsin, USA). The wavenumber accuracy of both spectrometers is better than  $0.05 \text{ cm}^{-1}$  (typically <0.03 cm<sup>-1</sup>). All spectra were measured at room temperature by averaging 100 scans at 4 cm<sup>-1</sup> resolution, ( $\Delta v = 2$  cm<sup>-1</sup>). Detailed sample handling and data processing parameters were given by Kuligiewicz et al. (2015). Briefly, a few drops of sonicated smectite suspensions in D<sub>2</sub>O (99.9 atom% D, Sigma-Aldrich Chemie GmbH, Steinheim, Germany) were deposited on the ATR crystal, covered with a custom-made cup as in Bukas et al. (2013), and purged with N<sub>2</sub> of purity better than 99.999 vol.% ( $<5 \text{ ppm H}_2\text{O}$ ) until a dry ~10 µm thick film was created in good contact with the ATR element. Subsequently, the cap was flushed with N<sub>2</sub> gas enriched with D<sub>2</sub>O vapor to 60-80% (±2%) relative humidity (RH) using a custommade gas blending setup equipped with a calibrated hygrometer. The sample was first equilibrated at ~60% RH ( $D_2O$ ) for at least 15 min in order to remove traces of  $H_2O$  or HDO. Subsequently, 5–10 spectra were collected in the 60 to 80% RH (D<sub>2</sub>O) range. This range was chosen because the vO-D value is stable between 60 and 80% RH (figure 7 of Kuligiewicz et al., 2015). Absorbance (ATR) maxima were determined for each

spectrum from the 2<sup>nd</sup> derivative minima calculated with a Savitzky-Golay algorithm (13 point smoothing) using the *Opus*<sup>(m)</sup> software by Bruker (Ettlingen, Germany). The reported values of vO–D (Table 1) represent the average from 5–10 spectra with errors given as one standard deviation (1 $\sigma$ ).

#### RESULTS

The spectra of D<sub>2</sub>O-saturated samples display structural O-H stretching bands in the region 3750 to  $3500 \text{ cm}^{-1}$  (Figure 2). O–D stretching and bending bands are at  $2800-2200 \text{ cm}^{-1}$ , and  $\sim 1200 \text{ cm}^{-1}$ , respectively, originating from the adsorbed D<sub>2</sub>O (Farmer and Russell, 1964; Russell et al., 1970; Kuligiewicz et al., 2015). A trace HDO bending mode could be observed at ~1450  $\text{cm}^{-1}$  and was employed to test the progress of the H/D exchange (Figure 2). The O–D stretching envelope included the high-frequency sharp band (Figure 2, inset) with vO-D in the  $2685-2700 \text{ cm}^{-1}$  range (Table 1). Because this work investigates the systematics of vO-D following Kuligiewicz et al. (2015), the vibrational features of the aluminosilicate network (e.g. Madejová et al., 2002; Madejová, 2003; Zviagina et al., 2004; Gates, 2005 and references therein) are not discussed here.

Samples from set A were examined against charges derived from the quantitative analysis of Na in the Na forms (Table 1, Figure 3). This is an extension of a



Figure 2. Examples of the IR spectra of  $D_2O$ -saturated dioctahedral smectites in the  $3800-1150 \text{ cm}^{-1}$  range. Three representative samples from set B are shown. The inset illustrates the determination of vO-D from the  $2^{nd}$  derivative spectra.

similar correlation by Kuligiewicz *et al.* (2015) with additional cationic forms. Values of vO–D in the available Mg, Ca, and Na forms differed by  $2\sigma$  or less (Table 1, Figure 3), and, therefore, could be fitted with a single linear regression equation with  $R^2 = 0.92$ . In contrast, the corresponding Cs forms exhibited a narrower distribution of vO–D values with a very weak dependence on charge, whereas the K forms occupy intermediate positions between Cs and the other cation forms (Figure 3). The values of vO–D for different cation forms converge toward the low-charge limit.

The dependence of vO–D on SFM-based layer charge for sets A–C (Table 1, excluding the K and Cs forms of set A) illustrated that linear correlations could be established for all sets, but the quality of the fit is variable ( $R^2 = 0.70-0.92$ ) and depends on the method of charge determination (Figure 4). Strong correlations were found between vO–D and AAM ( $R^2 = 0.92$ ) and AAOM ( $R^2 = 0.84$ ) charges, available for set B (Figure 5). A comparison of  $R^2$  coefficients for different trendlines available for sets A–C revealed that the SFM-based correlation for set B was considerably weaker than either the corresponding AAM-based correlations or the SFM-based correlations for the other sets (Figures 4, 5).

Finally, set D, based on the charge-reduced Li-SAz-1 montmorillonite, exhibited the largest shifts of vO–D toward  $\sim$ 2700 cm<sup>-1</sup> upon increase of the temperature of the charge-reduction treatment (Figure 6). Above  $\sim$ 120°C,

this shift was accompanied by a pronounced decrease in intensity of the sharp D<sub>2</sub>O band. As a result of decreasing D<sub>2</sub>O content, vO–D could not be measured for samples heated above 180°C (Figure 6). The correlation between vO-D and the CEC-derived total charges based on Skoubris et al. (2013) for set D was excellent (Figure 7,  $R^2 = 0.98$ ). The  $R^2$  for set D was greater than for all other sample sets, mostly due to the greatly expanded range of charges (and frequencies) probed down to ~0.2 p.f.u. (and up to ~2700 cm<sup>-1</sup>). Notably, the trend line of set D (originating from the same parent sample) could not be distinguished from that of set C (independent samples), suggesting that a common regression of vO-D against SFM-based charges can be applied for both Li-exchanged sample sets, regardless of their origin and structure ( $R^2 =$ 0.96; Figure 7). For the same sample set D, CEM charges were in good agreement with the CEC-derived SFM data but deviated from the SFM trend line as the layer charge approached the lower limit of CEM (~0.4 p.f.u.; Figure 7).

# DISCUSSION

## Charge prediction

The experimental results implied that strong correlations could be obtained between vO–D and either type of charge reference data (SFM or AAM). Some variability in the intercept and slope of the different SFM-based regressions was attributed to the different corrections or assumptions involved in the calculation of



Figure 3. Relationship between vO-D and total charge determined for set A in different cationic forms (Na, Ca, and Cs from Kuligiewicz *et al.*, 2015).

structural formulae in the reference methods and, to a lesser extent, to the different interlayer cations (Figure 4). Validation procedures can be employed to spot outliers in the literature data: *e.g.* removing Kinney montmorillonite from set C using SFM charges, would improve  $\mathbb{R}^2$  from 0.89 (Figure 4) to 0.97.

Any of the correlations between vO-D and SFM or AAM charges can be converted into an equation for layer-charge prediction (Figures 4, 5, 7). The two strongest linear correlations for the prediction of SFMor AAM-based charges were derived from sets C + D and B, respectively (Table 2, Figure 8). In both

Table 2. Linear regression parameters for the prediction of layer charge on the basis of vO–D. LC = p + q(vO-D - 2686)

	I I I	/
	Set B Na form (AAM)	Sets C and D Li form (SFM)
<b>p</b> <b>q</b> R <sup>2</sup> Number of samples Charge range (p.f.u.) 1σ		$\begin{array}{c} 0.60{\pm}0.01\\ -0.030{\pm}0.001\\ 0.96\\ 21\\ {\sim}0.2{-}0.6\\ 0.02 \end{array}$



Figure 4. Correlation between vO–D and SFM-based total layer charge for sample sets A–C. Only the Mg, Ca, and Na forms are presented for Set A. CEM charges for set C were included for comparison.

equations, the independent variable was written as (vO-D - 2686) in order to increase the variance of the frequency datasets. The offset value of 2686 cm<sup>-1</sup> was chosen on the basis of the Li-exchanged sets C and D as the closest integer wavenumber to the high-charge vO-D limit of a smectite definition: 0.6 p.f.u.

The calculated regression equations enabled the prediction of charge as deduced from either SFM or AAM in the  $\sim 0.2-0.6$  and 0.27-0.37 p.f.u. ranges, with standard deviations of 0.02 and 0.01 p.f.u., respectively (Table 2). These predictions were based specifically on the ATR infrared spectra of D<sub>2</sub>O-saturated samples and the determination of the position of the sharp highfrequency mode of  $D_2O$  from  $2^{nd}$  derivative analysis. The combination of spectral resolution (4  $\text{cm}^{-1}$ ), data spacing ( $\Delta v = 2 \text{ cm}^{-1}$ ), and Savitzky-Golay filter parameters for the calculation of 2<sup>nd</sup> derivatives (13 point smoothing) used in the present study were optimized previously by Kuligiewicz et al. (2015) to balance between resolution and signal-to-noise ratio requirements. A change in these parameters is feasible, but could affect the values of the regression coefficients (Table 2) and, therefore, may require recalibration or calibration transfer.

The regressions for the prediction of SFM and AAM charges (Table 2) were based on the spectra of homoionic Li- and Na-exchanged smectites, respectively. This choice was dictated by the availability of sample sets with independent and self-consistent reference charges. The spectroscopic prediction of SFM and AAM charges of an unknown smectite would strictly imply the preparation of the same cationic forms as in the relevant regression. Nevertheless, the study of set A has provided evidence that vO–D is nearly identical in the Na, Ca, and Mg forms of the same smectite (Figure 3; *cf.* Kuligiewicz *et al.*, 2015). This result implies that the choice of interlayer cation has a small effect on the prediction of the charge from vO–D, as long as its hydration enthalpy is high enough to keep the interlayers filled with D<sub>2</sub>O and allow for undisrupted D<sub>2</sub>O-basal surface interactions.

# Influence of interlayer cation and smectite charge location

The negligible effect that the common interlayer cations (Na<sup>+</sup>, Ca<sup>2+</sup>, and Mg<sup>2+</sup>) have on the value of vO–D is well explained by the assignment of this band to the dangling O–D bonds of interlayer D<sub>2</sub>O. In a hydrated smectite, these bonds point away from the strongly hydrogen-bonded D<sub>2</sub>O in the interlayer and towards the weakly charged basal surfaces (Figure 1; *cf.* Russell *et al.*, 1970; Suquet *et al.*, 1977). As such, their bond-length and associated vibrational energy are most sensitive to the density of the basal charge and not to the bonding of the hydrated interlayer cations. This is in contrast to the position of the bending mode of interlayer



Figure 5. Correlation between vO–D and smectite charges measured with different methods by Wolters *et al.* (2009) for set B. SFM total charges were calculated according to Stevens (1946), AAM and AAOM interlayer charges were calculated after Lagaly (1994) and Olis *et al.* (1990), respectively. Each trend line was derived using 15 samples (*cf.* table 6 of Wolters *et al.*, 2009).



Figure 6. Second derivative of vO-D infrared absorbance for set D. The numbers indicate the heating temperature of Li-SAz-1.



Figure 7. The vO-Dvs. total charge calculated for sets C and D with SFM (set D: based on CEC reduction). CEM charges for set C are included for comparison. The numbers indicate the heating temperature of SAz-1.

H<sub>2</sub>O (D<sub>2</sub>O) at ~1630 (~1200) cm<sup>-1</sup> and the broad stretching envelope at ~3400 (~2500) cm<sup>-1</sup>, which are cation-dependent (Farmer and Russell, 1966; Johnston *et al.*, 1992; Bishop *et al.*, 1994; Madejová *et al.*, 2002; Kuligiewicz *et al.*, 2015).

The K- and Cs-exchanged smectites are less suitable for LC prediction, especially towards the high-charge limit where the direct interaction of K<sup>+</sup> and Cs<sup>+</sup> with the siloxane surfaces occurs (Eberl et al., 1986; Onodera et al., 1998; Skiba, 2013). The greater the basal charge, the greater the attraction of Cs<sup>+</sup> (or K<sup>+</sup>) ions to charged basal oxygen atoms of the tetrahedral sheet. Binding a cation to the basal surface leads to local charge neutralization and, in extreme cases, to the collapse of the interlayer (cf. Christidis and Eberl, 2003; Ferrage et al., 2005, 2007; Skiba, 2013). In this case, a high-charge portion of the charged surface will not be available to  $H_2O$  ( $D_2O$ ) and, thus, it would not be probed by vO-D. Hence, the high-charge Cs-saturated smectites would behave as smectites with lower charge, leading to the shift of vO-D to higher wavenumbers (Figure 3). In agreement with the experimental data (Table 1, set A; Figure 3), an increase of vO-D from Mg-, Ca-, or Na- to K- and

Cs-exchanged smectite is expected to be most pronounced at the high-charge end because charge heterogeneity tends to increase with charge (Christidis and Eberl, 2003). Such a mechanism explains the similarity of the vO–D values in low-charge smectites regardless of the cationic form (Figure 3). Having intermediate hydration enthalpy, K<sup>+</sup> is bonded directly to fewer highcharge domains in the first shell of the siloxane surface than Cs<sup>+</sup>, leaving more surface available to the D<sub>2</sub>O probe.

Perhaps counter-intuitively, the vO–D values were found to be independent of charge location (tetrahedral *vs.* octahedral; Table 1), which implies that charges predicted on the basis of vO–D would also be unaffected by charge location. The sample sets A–C included dioctahedral smectites with a variable fraction of tetrahedral charge, including almost pure beidellite and montmorillonite end-members. In contrast, charge in set D was controlled by an octahedral-charge neutralization mechanism that involves Li<sup>+</sup>-migration and fixation. None of these factors introduced a measurable bias on vO–D. Samples 5MC, 2LP, and 4JUP (set B) had identical AAM charges and vO–D (0.28±0.01 p.f.u. and



Figure 8. Regressions for the prediction of AAM-based charges (set B, Na-exchanged) and SFM-based charges (sets C+D, Li-exchanged). Both trend lines intercept at the charge value of  $\sim$ 0.2 p.f.u. Dotted lines present 68% prediction bands. IC – interlayer charge; TC – total charge.

2692.2 $\pm$ 0.1 cm<sup>-1</sup>, Table 1) despite the very different fractions of beidellitic charge (13, 24, and 62%, respectively; Wolters *et al.* (2009)). Similarly, Li-Tx-1 (set C) and the Li-SAz-1-115 (set D) had identical SFM charges and vO-D (0.44 p.f.u. and 2691.2 $\pm$ 0.1 cm<sup>-1</sup>, Table 1) despite the obviously different origin of their layer charges.

# Other factors influencing layer-charge measurement by the O-D method

The charge predicted from vO–D is essentially the part of the total charge that is spread onto the basal surfaces of the water-accessible interlayers of smectite. In this sense, the proposed spectroscopic method is an interlayer-sensitive technique, conceptually similar to AAM or CEM, but based on the isotopic H/D exchange of H<sub>2</sub>O, a ubiquitous and intrinsic constituent of smectite. Smectites are heterogeneous with sample-specific layercharge distribution that can be probed in detail by AAM (Lagaly, 1994) and, in a coarser manner, by CEM (Christidis and Eberl, 2003). Individual vO–D contributions from layers of different charge could not be resolved spectroscopically. Hence, the O–D method yields a charge averaged over all water-accessible interlayers of all particles in the sample. This average is not necessarily identical to the average layer charge. Collapsed layers, should they exist (*cf.* Ferrage *et al.*, 2005, 2007), would not be probed by D<sub>2</sub>O. Similarly, the vO–D of interstratified illite-smectite (Cetin and Huff, 1995; Środoń *et al.*, 2009) or kaolinite-smectite (Cuadros and Dudek, 2006; Dudek *et al.*, 2006) minerals would only reflect the charge of the smectitic component.

Both AAM- and SFM-based charges are highly correlated to vO–D; therefore, these two types of layer charge are linearly related to each other. Similar conclusions were reached by Laird (1994) and Kaufhold (2006). Interestingly, the two types of charge measurement converge at the low-charge limit of LC (Figure 8): a value of vO–D equal to 2699 cm<sup>-1</sup> would correspond to essentially identical SFM- and AAM-type charges ( $0.21\pm0.02$  and  $0.19\pm0.01$ , respectively), which coincides with the low-charge limit in the definition of

smectite (*cf.* Laird, 1994). Detailed discussion of reasons for discrepancies between AAM and SFM is, however, beyond the scope of the present work.

## CONCLUSIONS

The linear relationship between the position of the dangling O–D stretching mode in the FTIR spectra of D<sub>2</sub>O-saturated smectites and their layer charge (Kuligiewicz *et al.*, 2015) was confirmed on various sample sets with charges determined by different established techniques. Statistically significant linear correlations were obtained between vO–D and both SFM and AAM charges. These linear regressions can be used for predicting the charge of dioctahedral smectites with high accuracy ( $\leq 0.02$  p.f.u.) and over the entire range of smectite layer charge.

The charge-predicting equations were based on Liand Na-exchanged smectites (SFM and AAM charges, respectively), which added a cation exchange step prior to the spectroscopic measurement. On the other hand, Ca<sup>2+</sup> and Mg<sup>2+</sup>, in addition to Na<sup>+</sup>, are the most common cations in sedimentary environments and the most predominant exchangeable cations on clay surfaces (e.g. Odom, 1984; Dohrmann et al., 2012). The present results suggest that interlayer cations of relatively high hydration enthalpy such as Na<sup>+</sup>, Ca<sup>2+</sup>, and Mg<sup>2+</sup> do not bias the vO-D values (Figure 3). On such a basis, a precise estimate of layer charge can probably be obtained directly from smectites in their natural cationic forms. More accurate charge determinations would require samples in the same cationic form as those used to establish the prediction equation.

The stretching energy of the dangling O-D bond of adsorbed  $D_2O$  probes the basal/interlayer charge, *i.e.* the portion of charge referred to as permanent charge or layer charge. The proposed method and prediction equations provide an average charge of the smectite layers that are hydrated and accessible to H/D exchange. The prediction equations are based on band positions rather than intensities and, therefore, do not require normalization. The method is insensitive to the presence of anhydrous impurities. The measurement is fast and requires small amounts of sample (~5 mg), as well as relatively common instrumentation. As such, the proposed O-D method is an appealing low-cost alternative for determining the layer charge of dioctahedral smectite in high-throughput applications, such as the mapping of bentonite deposits (e.g. Christidis, 2001; Christidis et al., 2006; Christidis and Huff, 2009). On a more fundamental research front, this can be the techniqueof-choice for determining the charge of expandable layers in mixed-layered clay minerals such as illitesmectite, chlorite-smectite, and kaolinite-smectite.

In its present form, the method was tested on natural and reduced-charge dioctahedral smectites. Further effort is needed to extend this approach to additional sample sets, such as synthetic smectites, vermiculites, and trioctahedral smectites. Preliminary results on a synthetic high-charge saponite (Kuligiewicz *et al.*, 2015) and beidellites (Kuligiewicz *et al.*, unpublished data), suggest that an approach similar to that proposed can be developed for trioctahedral and synthetic phases.

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